Giant Resonance Studies
Perspectives

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- **Introduction**

  Giant resonances: Fundamental modes of nuclear excitation

- **Importance of studying compression modes in nuclei**

  ISGMR and ISGDR excitation

- **Experimental studies**

  Incompressibility of nuclear matter

  Asymmetry term

- **Conclusions and outlook**

In the following:

IS = Iso-Scalar

IV = Iso-Vector

S = Spin

G = Giant

M = Monopole

D = Dipole

Q = Quadrupole

O = Octupole

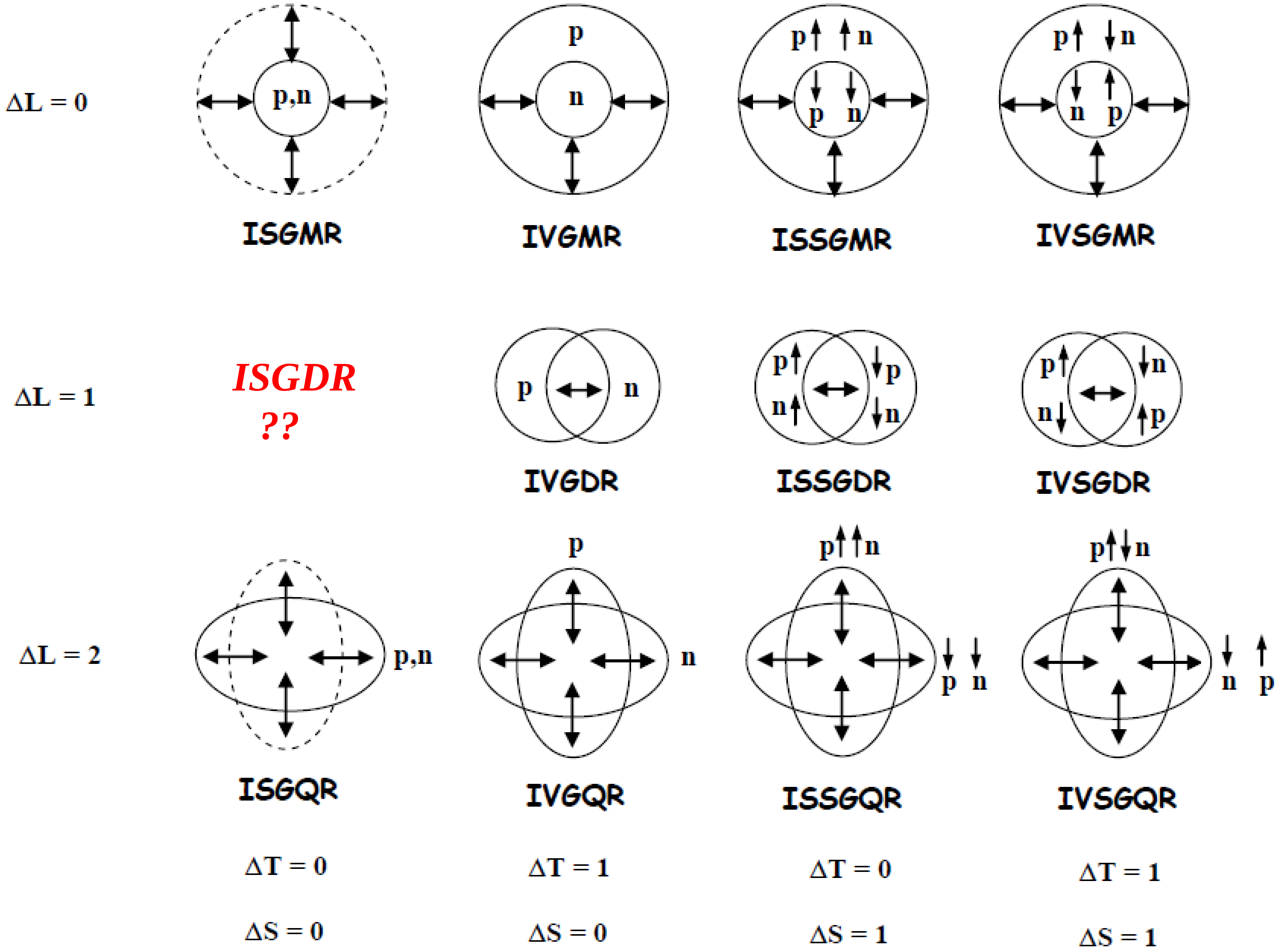
E.g. ISGMR = Isoscalar giant monopole resonance

ISGDR = Isoscalar giant dipole resonance

IVGDR = Isovector giant dipole resonance

IVSGMR= Isovector spin giant monopole resonance

IVSGDR = Isovector spin giant dipole resonance


ΔL = 0
p,n
ISGMR
p
n
IVGMR
p n
p n
ISSGMR
p n
n p
IVSGMR
ΔL = 1
ISGDR
??
p
n
IVGDR
p
n
p
n
ISSGDR
p
n
n
p
IVSGDR
ΔL = 2
p,n
ISGQR
p
n
IVGQR
p n
p n
ISSGQR
p n
n p
IVSGQR
ΔT = 0
ΔS = 0
ΔT = 1
ΔS = 0
ΔT = 0
ΔS = 1
ΔT = 1
ΔS = 1

Microscopic picture: GRs are coherent (1p-1h) excitations induced by single-particle operators.

- **Excitation energy depends on**

  ***i*) multipole *L* (*L*ħω, since radial operator ∝ r^L; except for ISGMR and ISGDR, 2ħω & 3ħω, respectively),**

  ***ii*) strength of effective interaction and**

  ***iii*) collectivity.**
- **Exhaust appreciable % of EWSR**
- **Acquire a width due to coupling to continuum and to underlying 2p-2h configurations.**

Microscopic structure of ISGMR & ISGDR

Transition operators:

$$O^{L=0} = \sum_i r_i^0 Y_0^0 + \frac{1}{2}\sum_i r_i^2 Y_0^0 + \ ...$$

Constant Overtone

2ħω excitation

$$O^{L=1} = \sum_i r_i^1 Y_0^1 + \frac{1}{2}\sum_i r_i^3 Y_0^1 + \ ...$$

Spurious c.o.m. motion Overtone

3ħω excitation (overtone of c.o.m. motion)

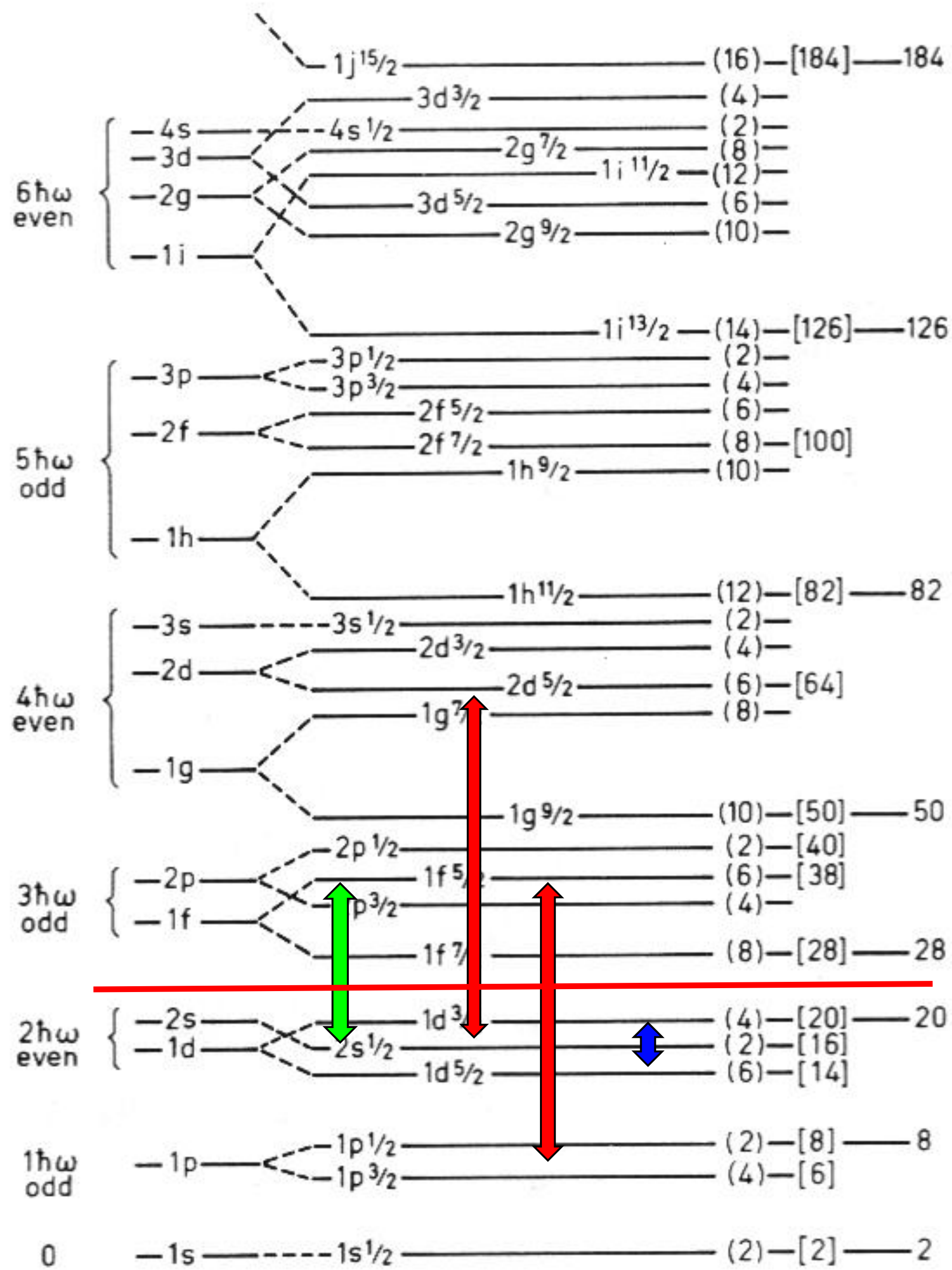
Nucleus → **Many-body system with a finite size**

Vibrations → **Multipole expansion with r, Y_{lm}, τ, σ**

	ΔS=0, ΔT=0	ΔS=0, ΔT=1	ΔS=0, ΔT=1	ΔS=1, ΔT=1	ΔS=1, ΔT=1
L=0: Monopole	**ISGMR** r^2Y_0	**IAS** τY_0	**IVGMR** $\tau r^2 Y_0$	**GTR** $\tau \sigma Y_0$	**IVSGMR** $\tau \sigma r^2 Y_0$
L=1: Dipole	**ISGDR** r^3Y_1		**IVGDR** $\tau r Y_1$		**IVSGDR** $\tau\sigma r Y_1$
L=2: Quadrupole	**ISGQR** r^2Y_2		**IVGQR** $\tau r^2 Y_2$		**IVSGQR** $\tau\sigma r^2 Y_2$
L=3: Octupole	**LEOR, HEOR** r^3Y_3				

ΔN = 1 E1 (IVGDR)

ΔN = 2 E2 (ISGQR)
&
ΔN = 0 E0 (ISGMR)



Decay of giant resonances

- **Width of resonance**

  $\Gamma, \Gamma^{\uparrow}, \ \Gamma^{\downarrow} \ (\Gamma^{\downarrow\uparrow}, \Gamma^{\downarrow\downarrow})$

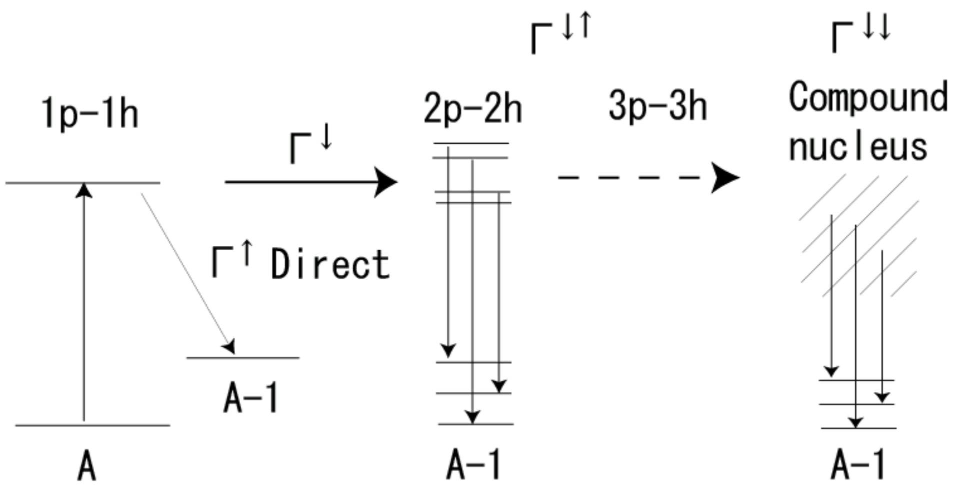
  - **$\Gamma^{\uparrow}$: direct or escape width**
  - **$\Gamma^{\downarrow}$: spreading width**

    $\Gamma^{\downarrow\uparrow}$: pre-equilibrium, $\Gamma^{\downarrow\downarrow}$: compound

- **Decay measurements**

  ⇒ Direct reflection of damping processes

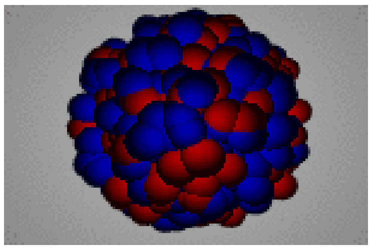
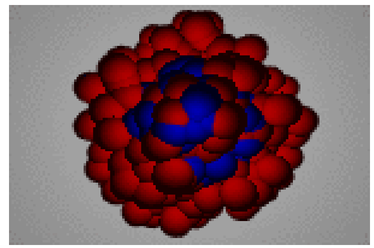
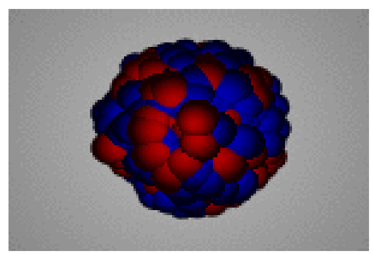
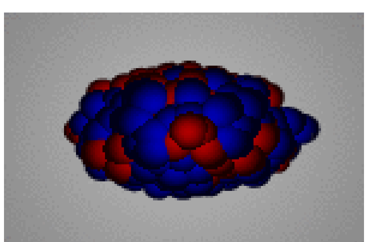
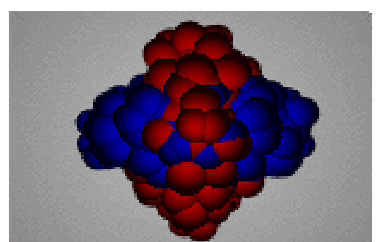
  Allows detailed comparison with theoretical calculations



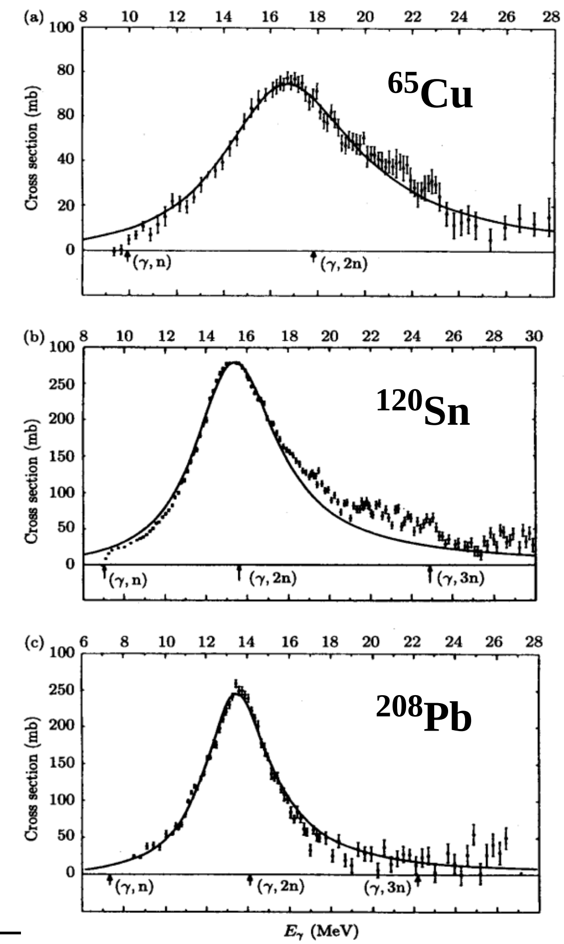
The collective response of the nucleus
Giant Resonances

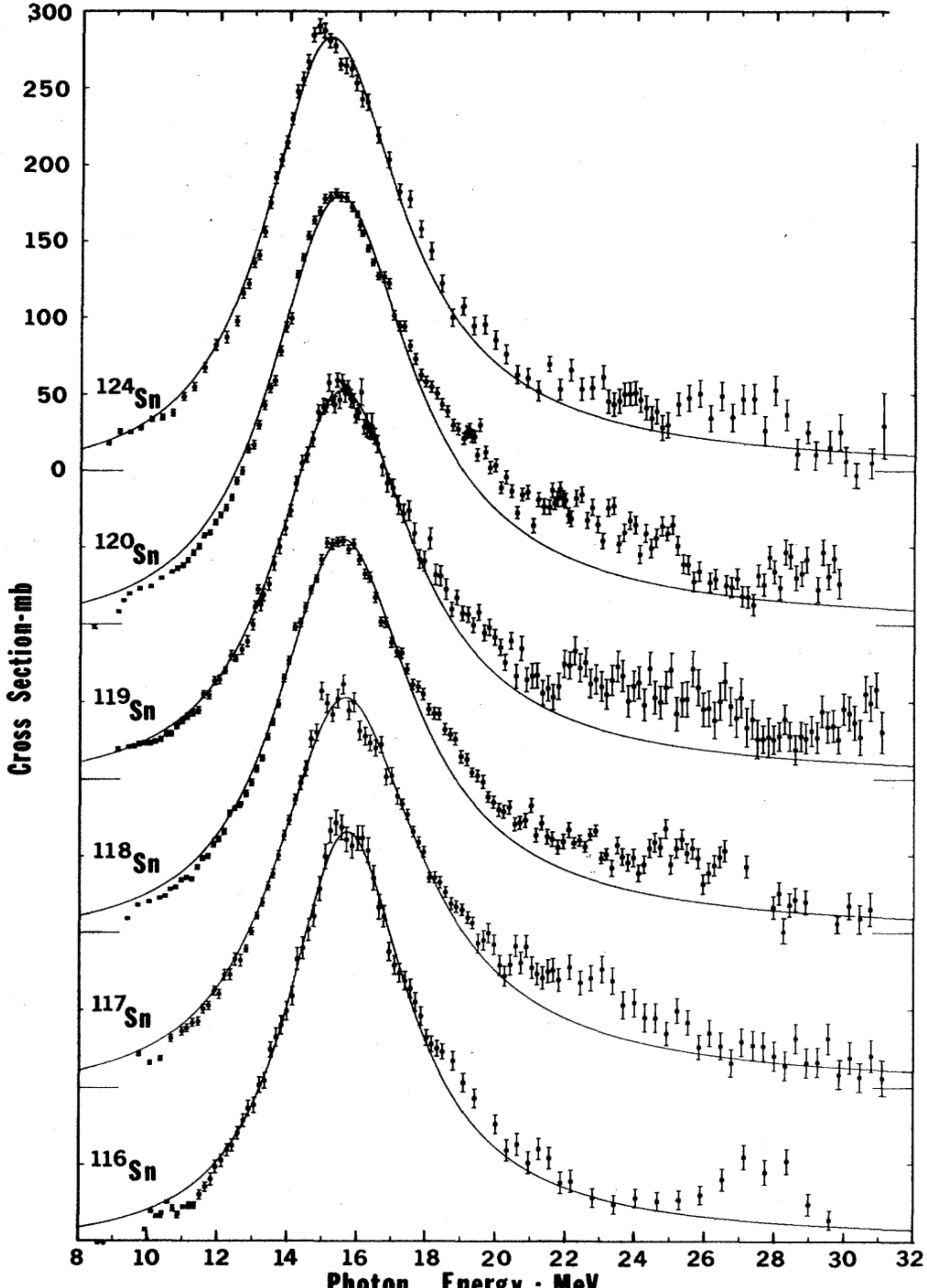
Electric giant resonances

Photo-neutron cross sections

	Isoscalar	Isovector
Monopole (GMR)		
Dipole (GDR)		
Quadrupole (GQR)		

Berman and Fulz, Rev. Mod. Phys. 47 (1975) 47



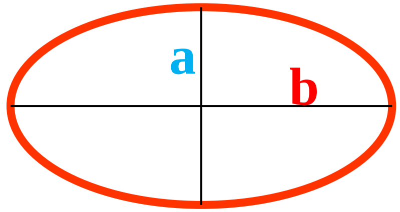


Measurement of the giant dipole resonance with mono-energetic photons

B.L. Berman and S.C. Fultz

Rev. Mod. Phys. 47 (1975) 713

Nucleus	Centroid (MeV)	Width (MeV)
^{116}Sn	15.68	4.19
^{117}Sn	15.66	5.02
^{118}Sn	15.59	4.77
^{119}Sn	15.53	4.81
^{120}Sn	15.40	4.89
^{124}Sn	15.19	4.81



Quadrupole deformation:

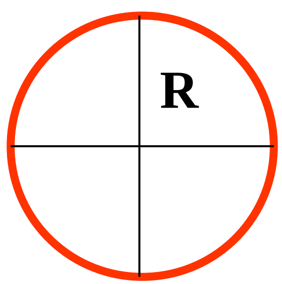
$\beta_2 = 0.275$

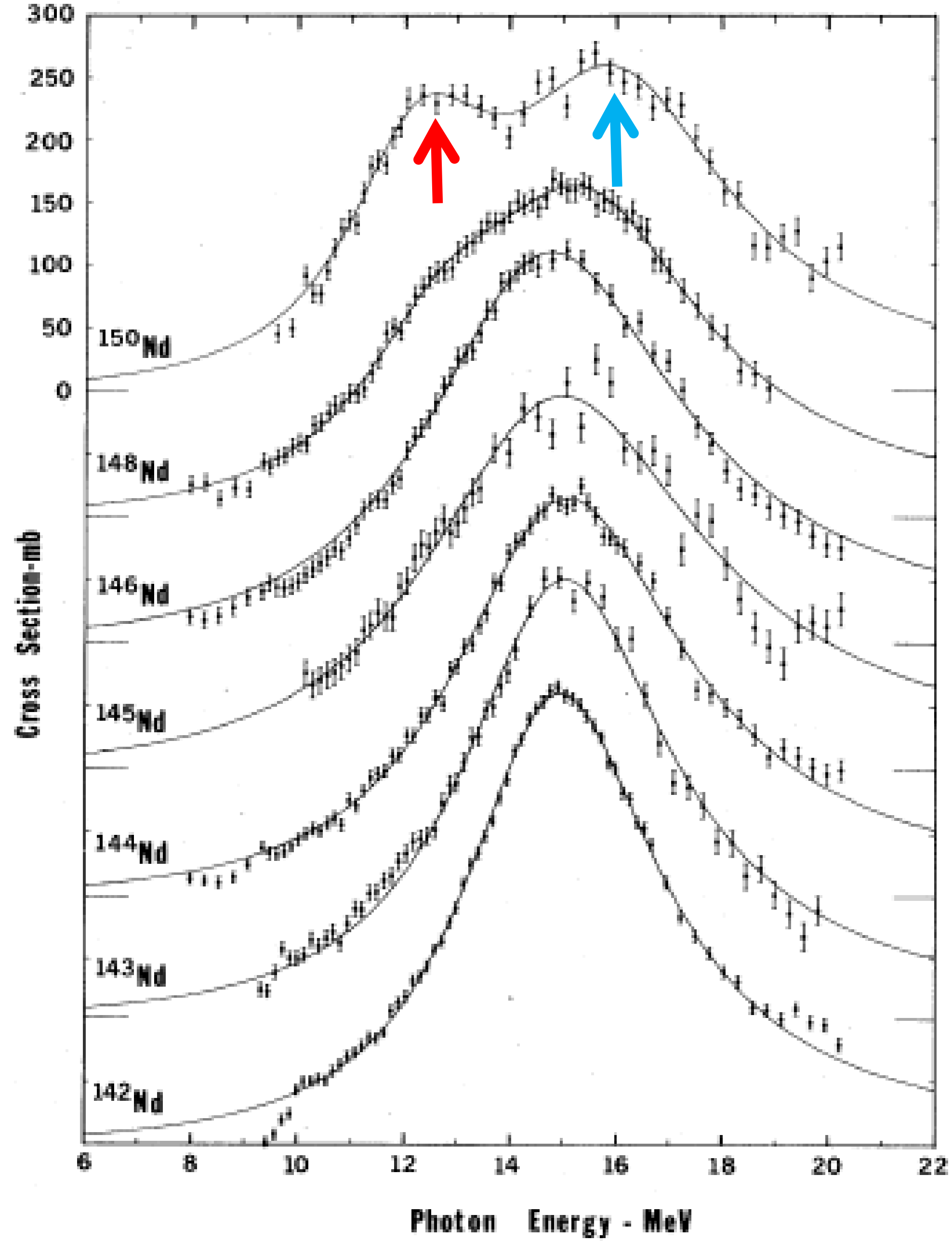
Excitation energies:

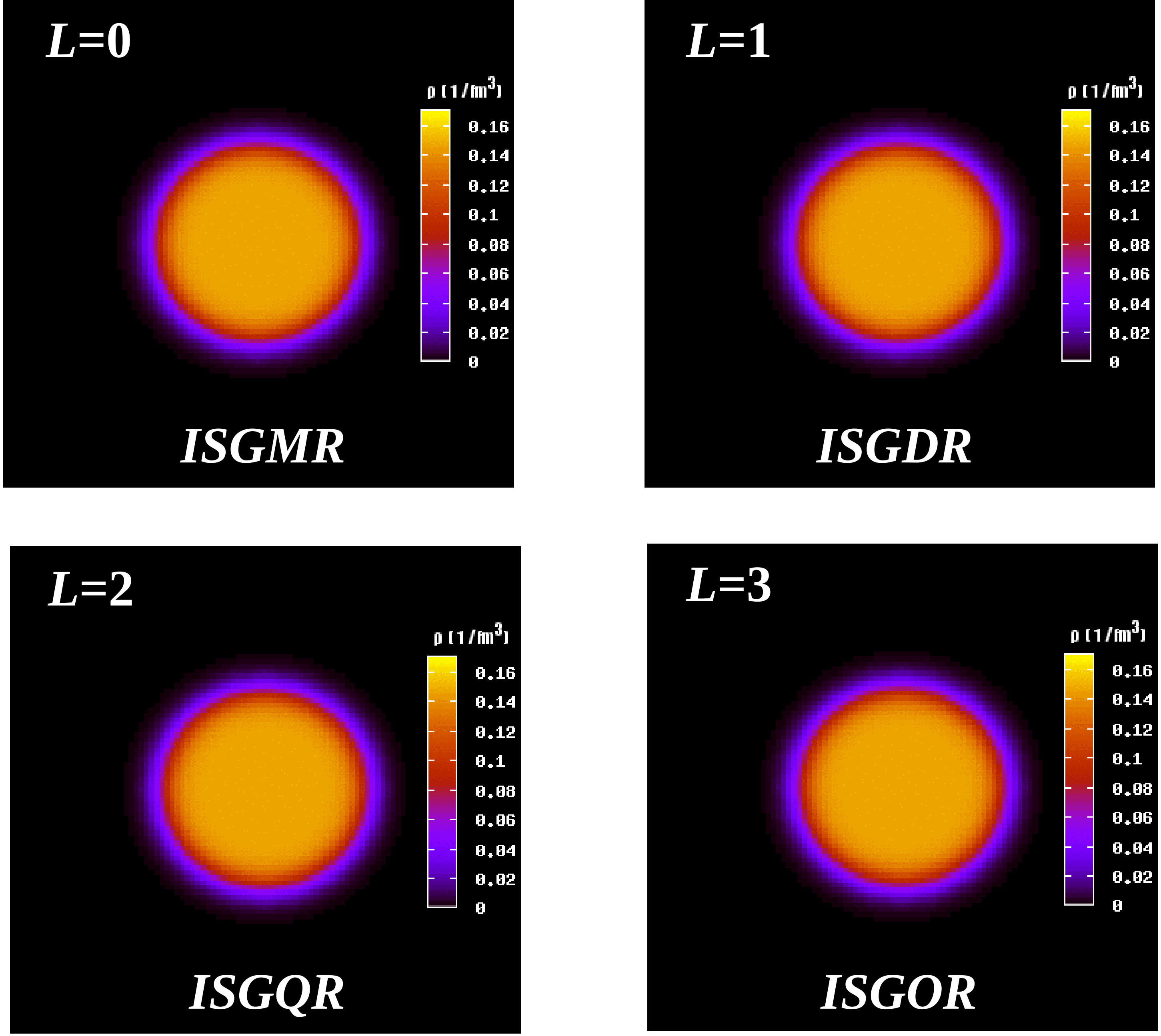
$E_2/E_1 = 0.911\eta + 0.089$

Where $\eta = b/a$

$s_1/s_2 = 1/2$






L=0
ρ (1/fm³)
0.16
0.14
0.12
0.1
0.08
0.06
0.04
0.02
0
ISGMR
L=1
ρ (1/fm³)
0.16
0.14
0.12
0.1
0.08
0.06
0.04
0.02
0
ISGDR
L=2
ρ (1/fm³)
0.16
0.14
0.12
0.1
0.08
0.06
0.04
0.02
0
ISGQR
L=3
ρ (1/fm³)
0.16
0.14
0.12
0.1
0.08
0.06
0.04
0.02
0
ISGOR

M. Itoh

In fluid mechanics, **compressibility** is a measure of the relative volume change of a fluid as a response to a pressure change.

$$\beta = -\frac{1}{V}\frac{\partial V}{\partial P}$$

where P is pressure, V is volume.

Incompressibility or **bulk modulus** (K) is a measure of a substance's resistance to uniform compression and can be formally defined:

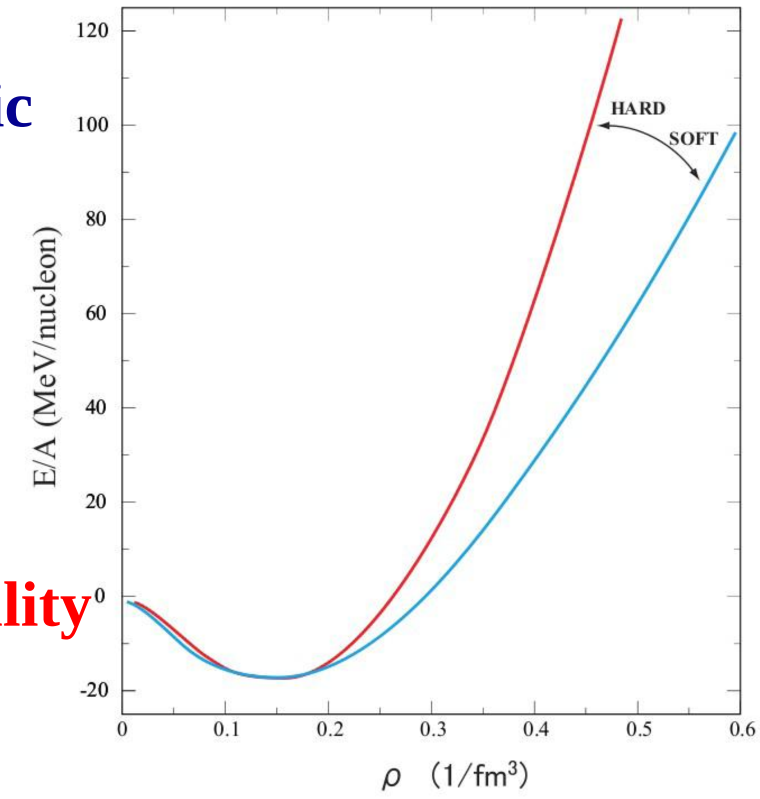
$$K = -V\frac{\partial P}{\partial V}$$

For the equation of state of symmetric nuclear matter at saturation nuclear density:

$$\left[\frac{d(E/A)}{d\rho}\right]_{\rho=\rho_0} = 0$$

and one can derive the incompressibility of nuclear matter:

$$K_{nm} = \left[9\rho^2 \frac{d^2(E/A)}{d\rho^2}\right]_{\rho=\rho_0}$$



***E/A*: binding energy per nucleon**

***ρ*   : nuclear density**

ρ₀   **: nuclear density at saturation**

J.P. Blaizot, Phys. Rep. 64 (1980) 171

Equation of state (EOS) of nuclear matter:

More complex than for infinite neutral liquids:
Neutrons and protons with different interactions
Coulomb interaction of protons

1. **Governs the collapse and explosion of giant stars (supernovae)**
2. **Governs formation of neutron stars (mass, radius, crust)**
3. **Governs collisions of heavy ions.**
4. **Important ingredient in the study of nuclear properties.**

Isoscalar Excitation Modes of Nuclei

Hydrodynamic models/Giant Resonances

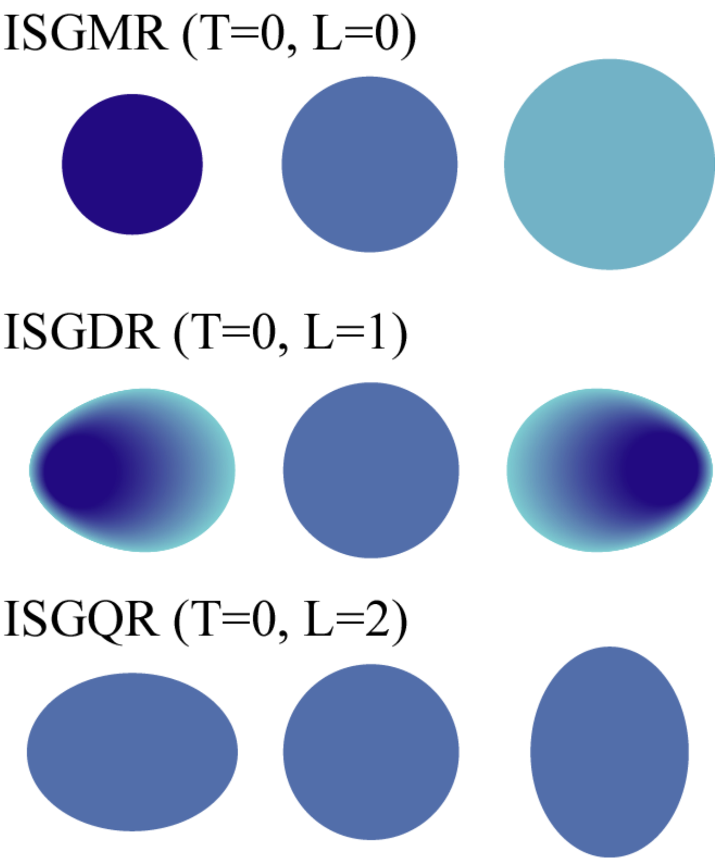
Coherent vibrations of nucleonic fluids in a nucleus.

Compression modes : ISGMR, ISGDR

In Constrained and Scaling Models:

$$E_{ISGMR} = \hbar\sqrt{\frac{K_A}{m\langle r^2\rangle}}$$

$$E_{ISGDR} = \hbar\sqrt{\frac{7}{3}\frac{K_A + \frac{27}{25}\varepsilon_F}{m\langle r^2\rangle}}$$



ε_F is the Fermi energy and the nucleus incompressibility:

$$K_A = \left[r^2(d^2(E/A)/dr^2)\right]_{r=R_0}$$

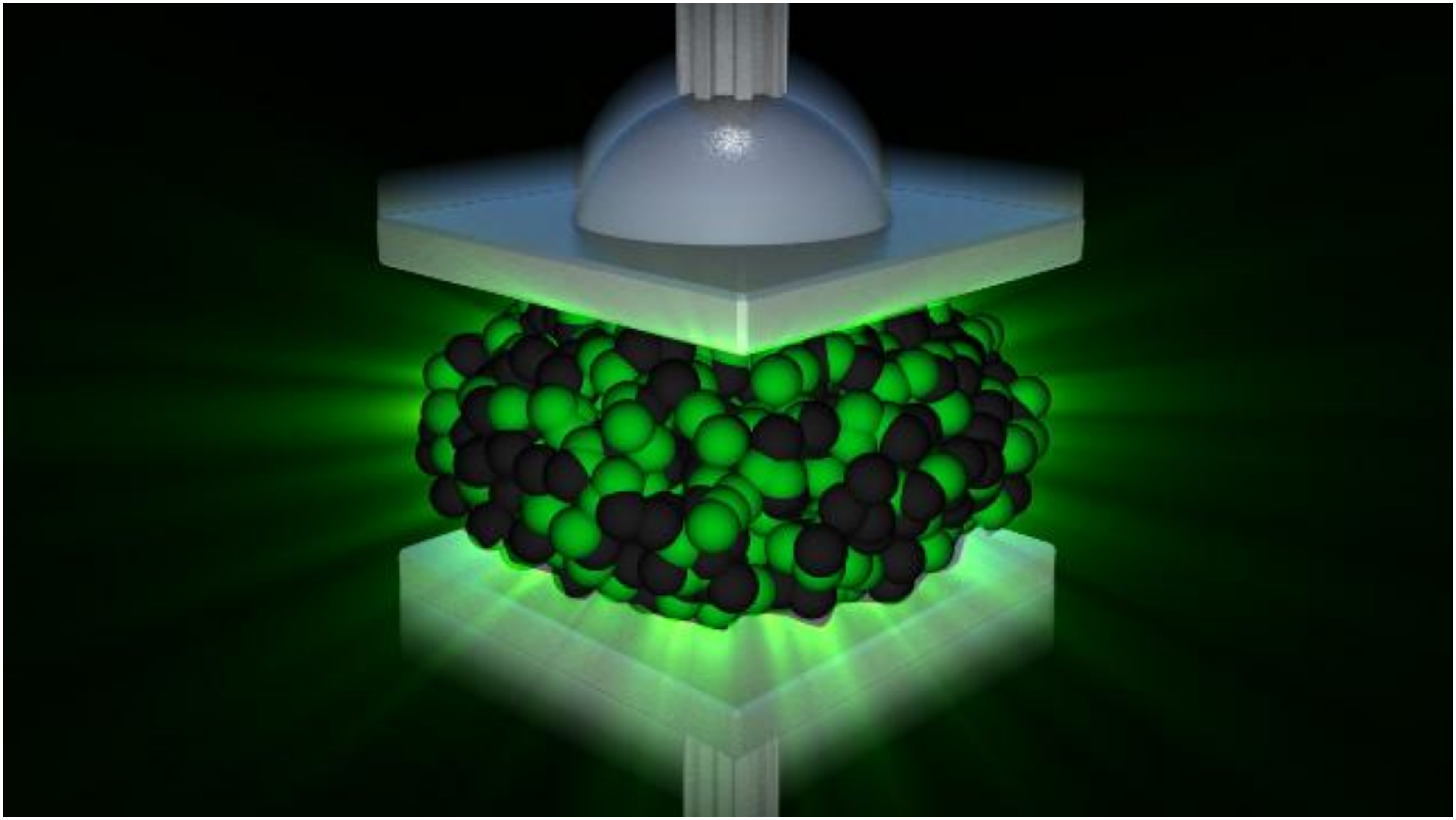
J.P. Blaizot, Phys. Rep. 64 (1980) 171

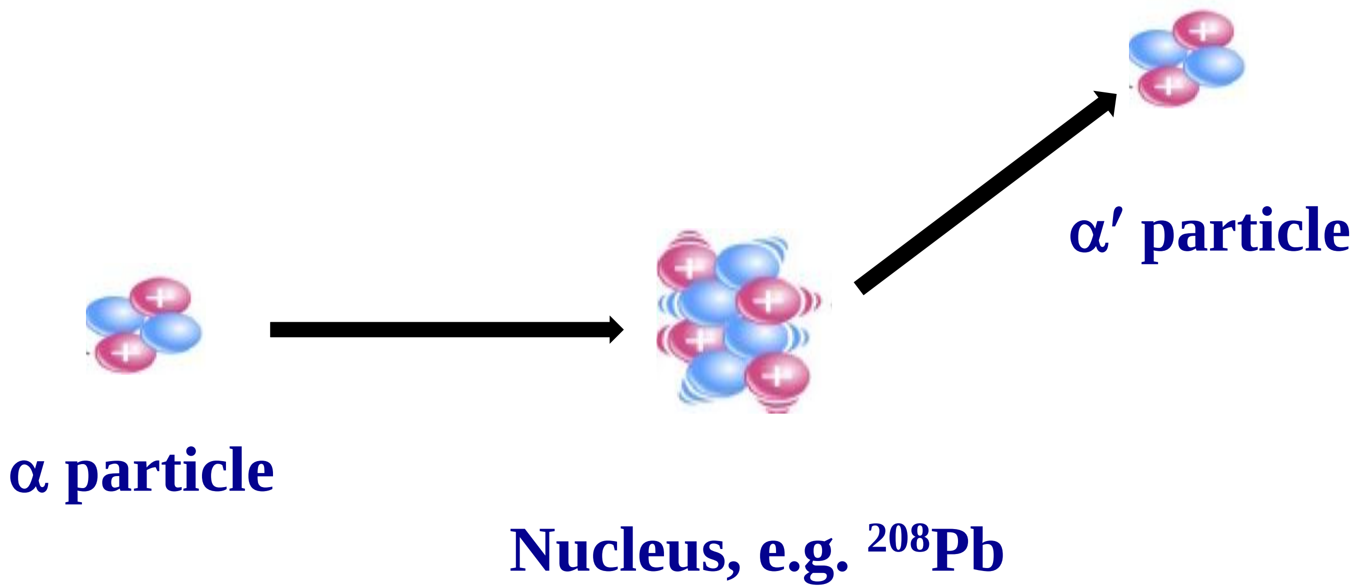
Giant resonances

- **Macroscopic properties: E_x, Γ, %EWSR**
- **Isoscalar giant resonances; compression modes**

  ISGMR, ISGDR $\Rightarrow$ Incompressibility, symmetry energy

$$K_A = K_{vol} + K_{surf}A^{-1/3} + K_{sym}((N-Z)/A)^2 + K_{Coul}Z^2A^{-4/3}$$

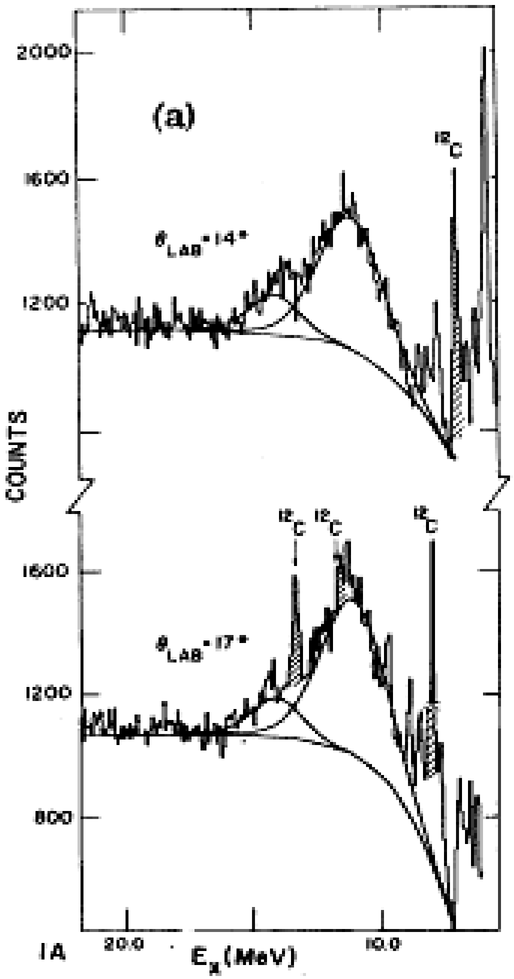




Inelastic α scattering

ISGQR, ISGMR

KVI (1977)



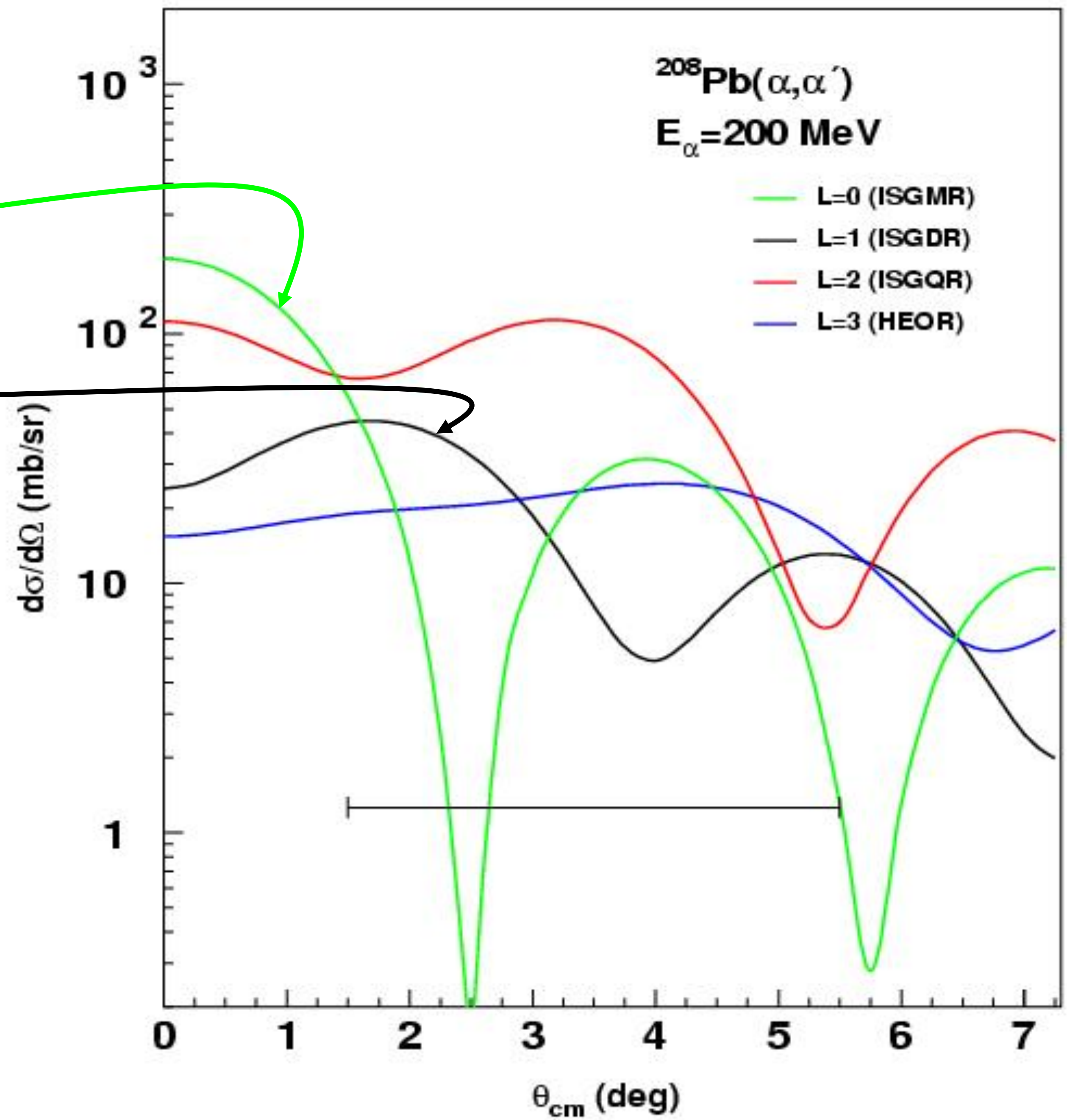
Large instrumental background!

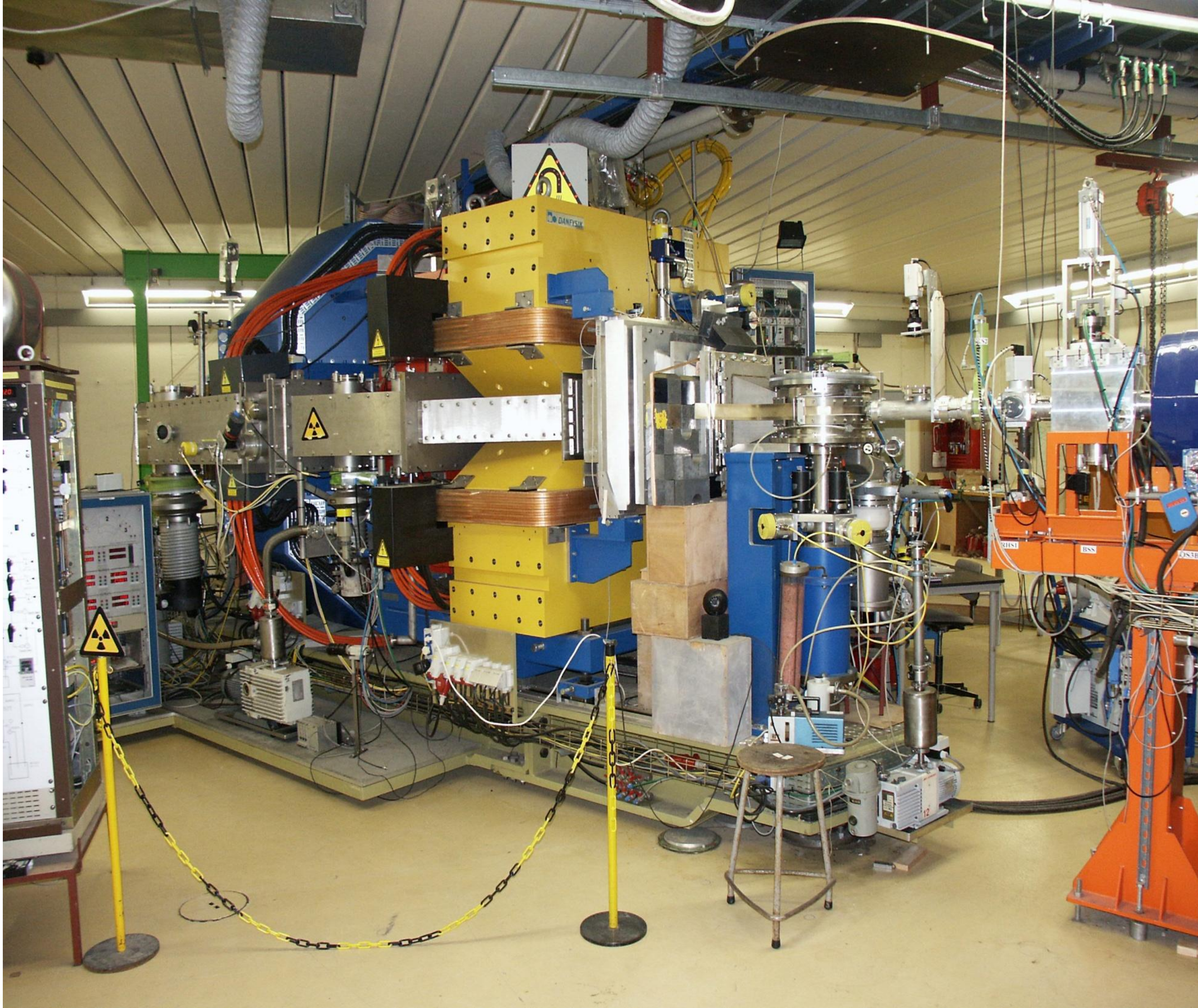
⇐ ^{208}Pb(α,α′) at E_α=120 MeV

M. N. Harakeh *et al.*, Phys. Rev. Lett. 38, 676 (1977)

ISGMR $L = 0$

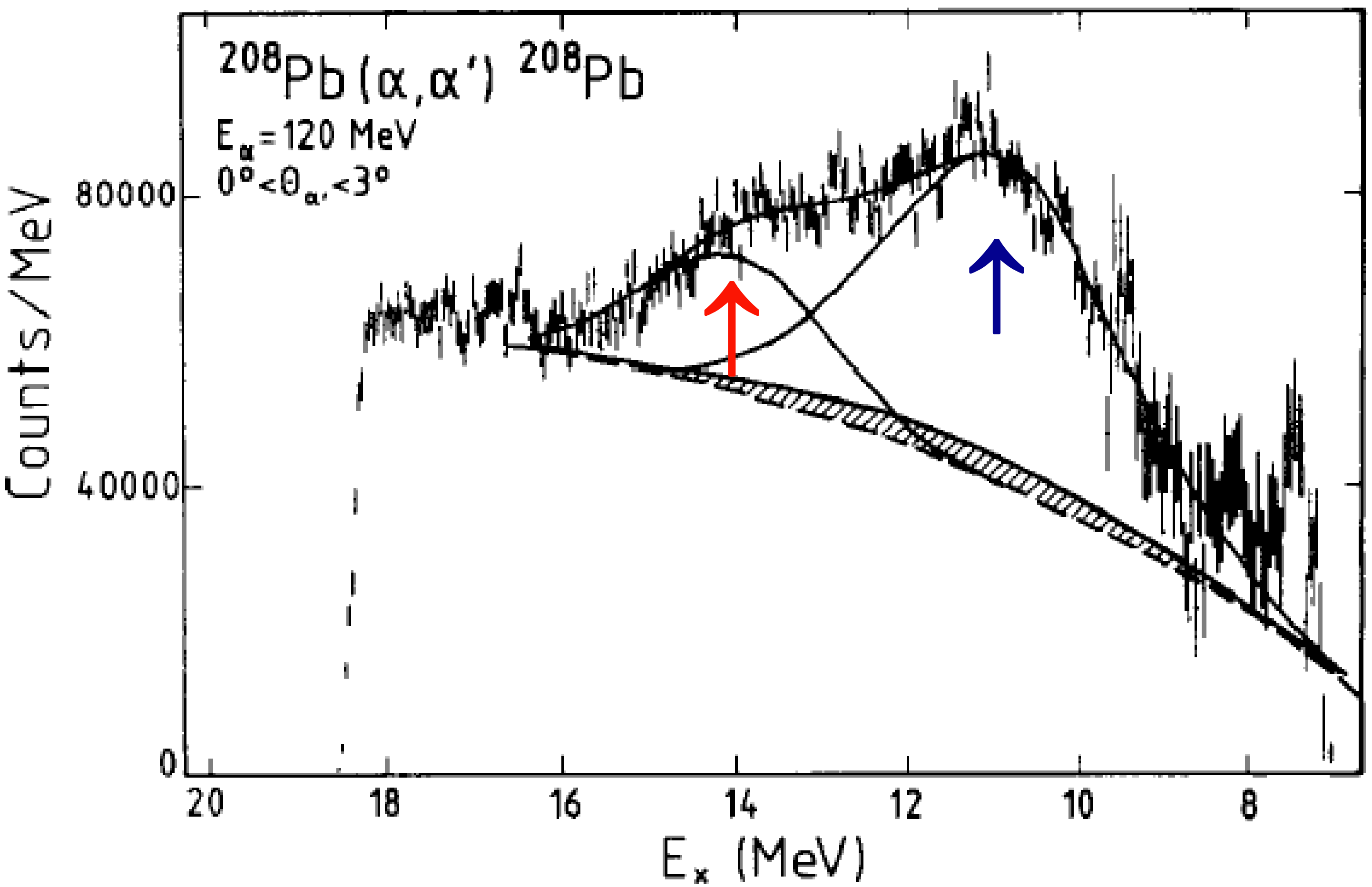
ISGDR $L = 1$




DANFYSIK

ISGQR at 10.9 MeV

ISGMR at 13.9 MeV



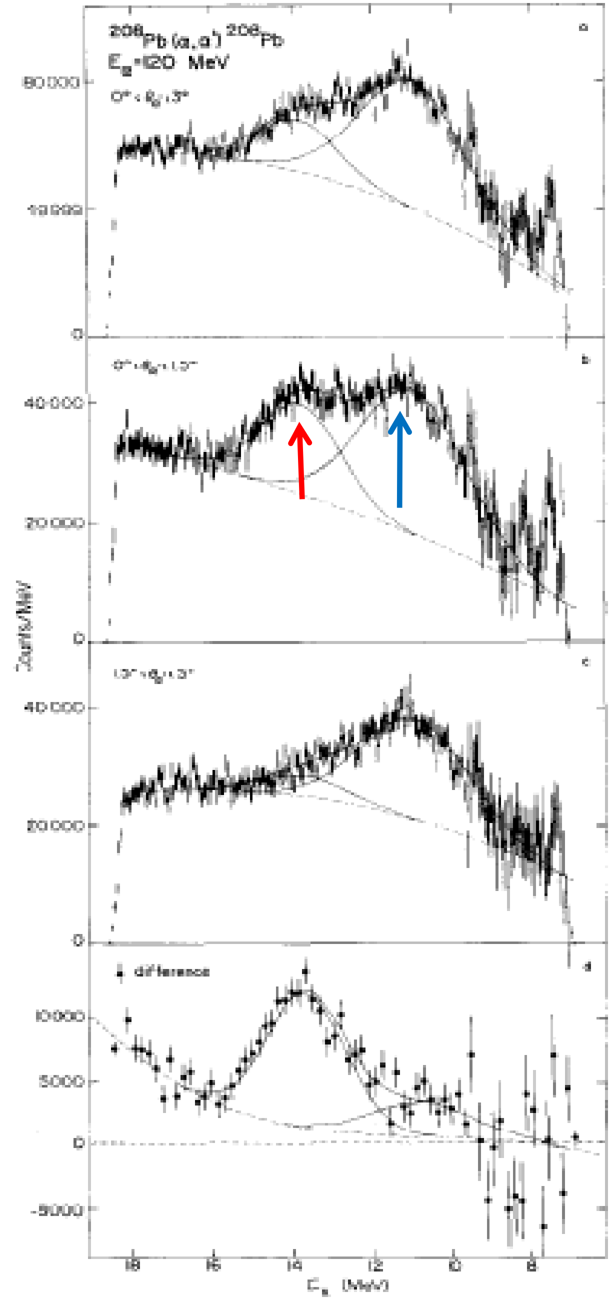
Difference of spectra

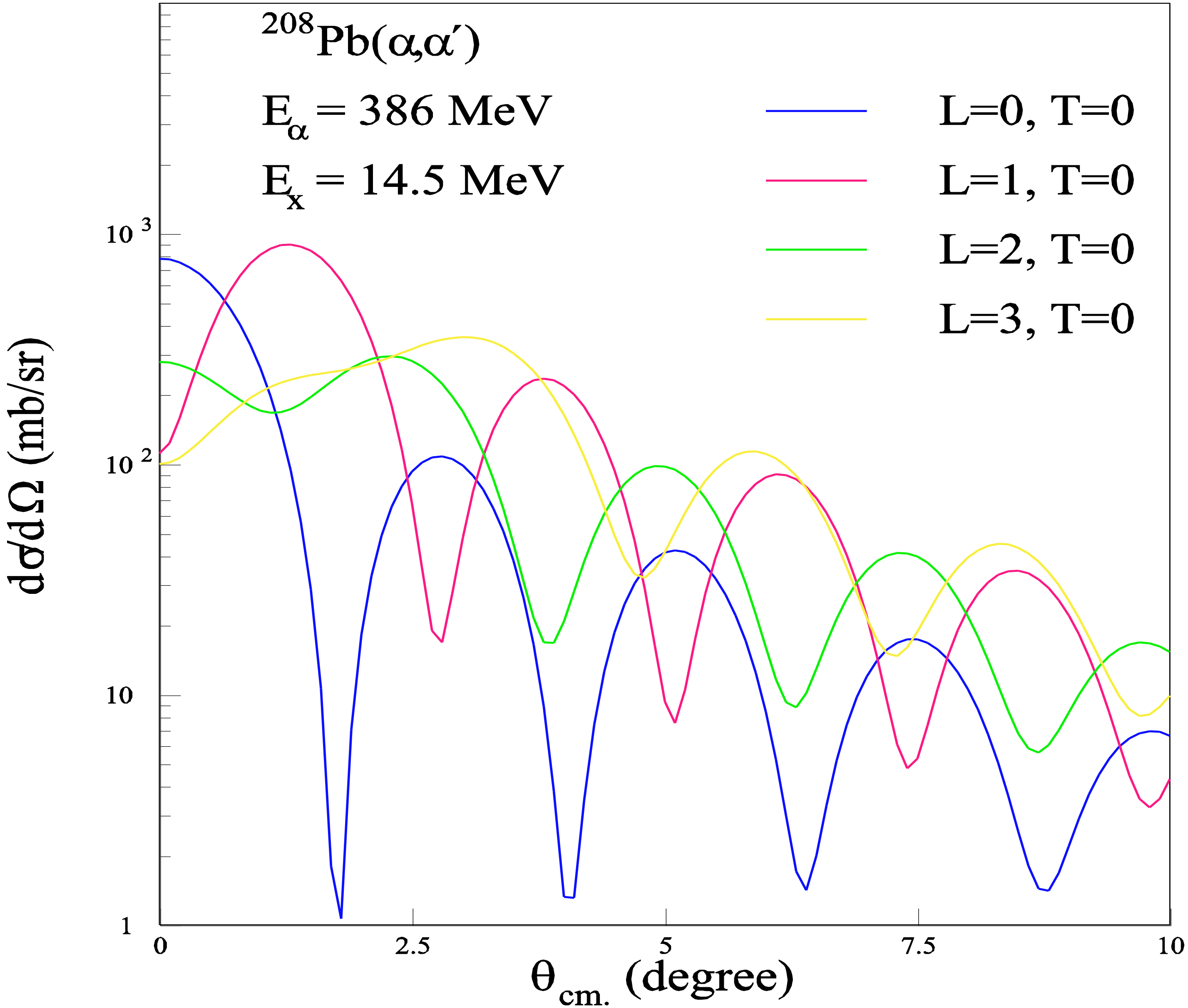
$0° < \theta_{\alpha'} < 3°$

$0° < \theta_{\alpha'} < 1.5°$

$1.5° < \theta_{\alpha'} < 3°$

Difference



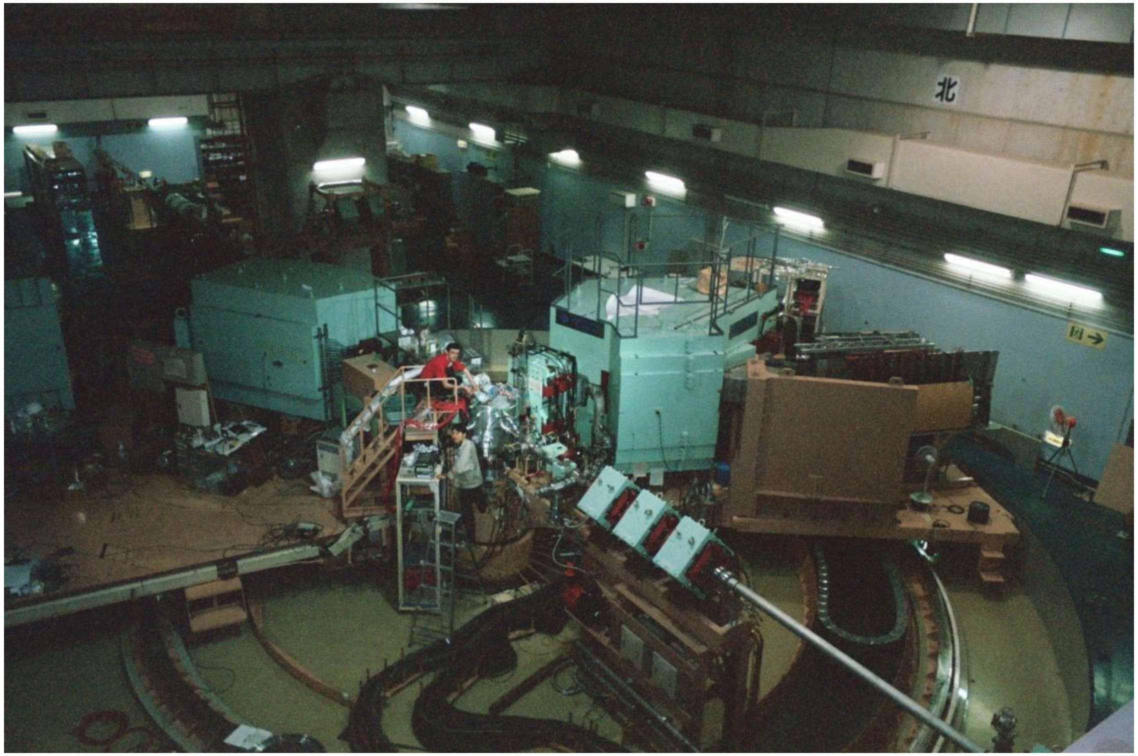


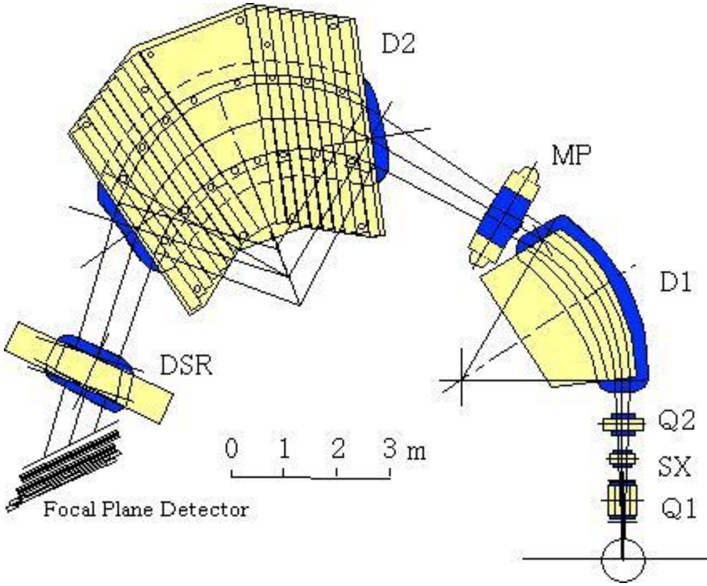
ISGMR, ISGDR

ISGQR, HEOR

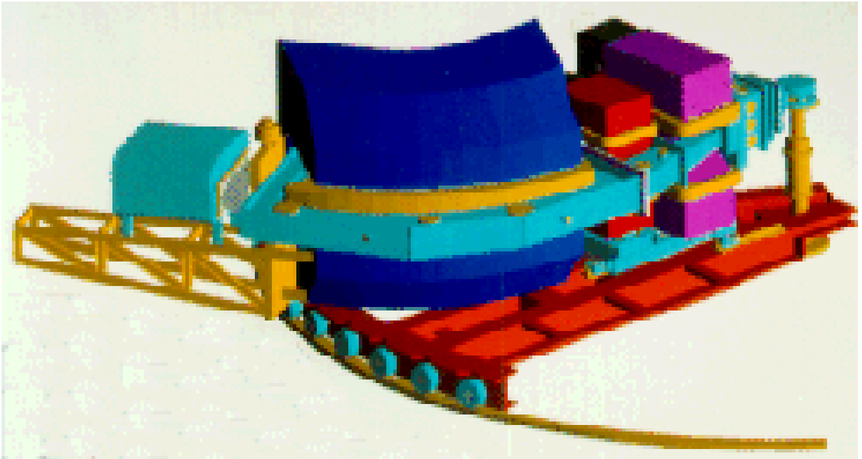
100 % EWSR

At E_x= 14.5 MeV



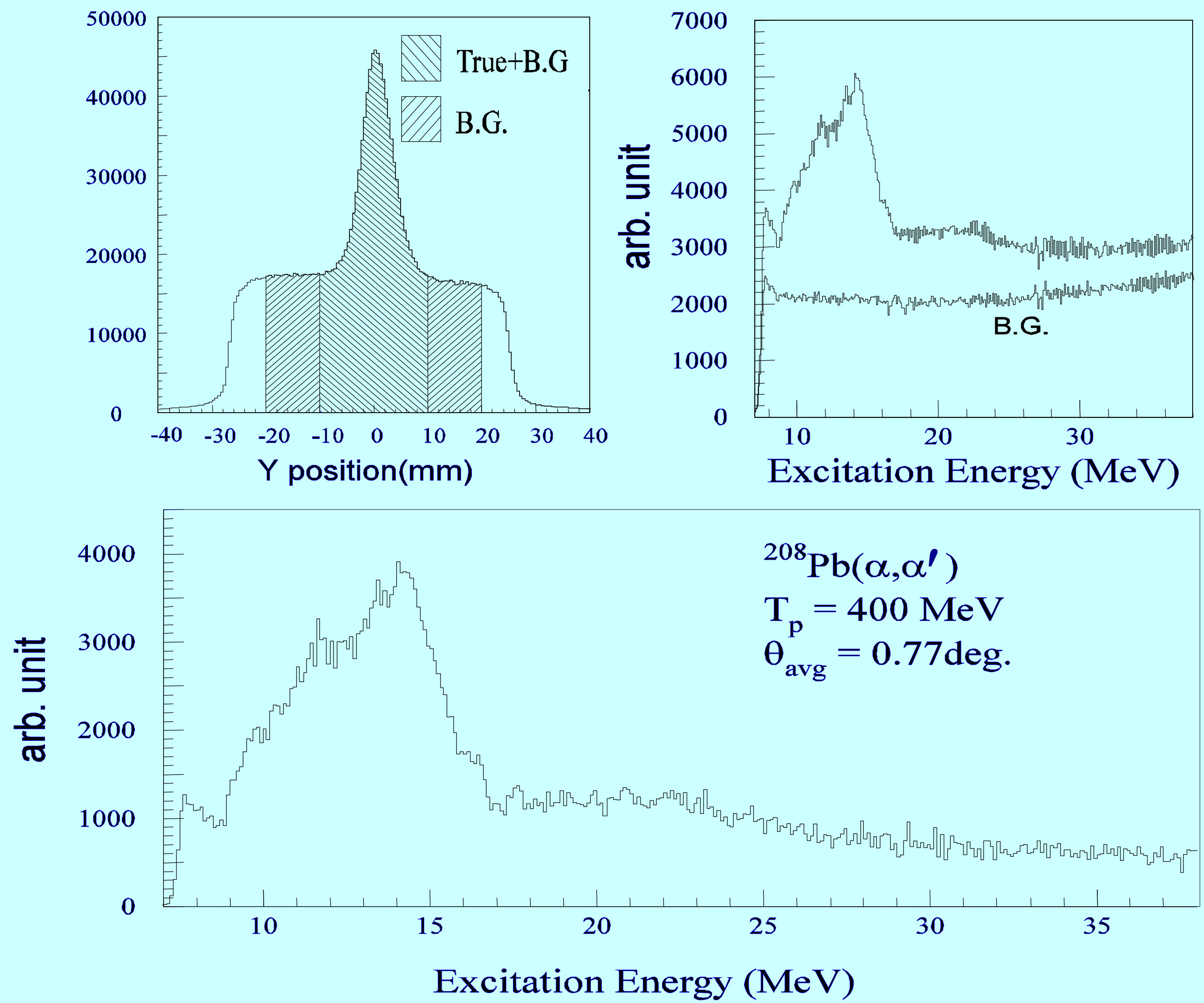


Grand Raiden@ RCNP



(α,α′) at E_{α}~ 400 & 200 MeV at RCNP & KVI, respectively

BBS@KVI


50000
40000
30000
20000
10000
0
-40 -30 -20 -10 0 10 20 30 40
Y position(mm)
True+B.G
B.G.
arb. unit
7000
6000
5000
4000
3000
2000
1000
0
10 20 30
Excitation Energy (MeV)
B.G.
arb. unit
4000
3000
2000
1000
0
10 15 20 25 30 35
Excitation Energy (MeV)
208Pb(α,α′)
Tp = 400 MeV
θavg = 0.77deg.

Multipole decomposition analysis (MDA)

$$\left(\frac{d^2\sigma}{d\Omega dE}(\vartheta_{c.m.},E)\right)^{\exp.} = \sum_L a_L(E)\left(\frac{d^2\sigma}{d\Omega dE}(\vartheta_{c.m.},E)\right)_L^{calc.}$$

$$\left(\frac{d^2\sigma}{d\Omega dE}(\vartheta_{c.m.},E)\right)^{\exp.} : \text{Experimental cross section}$$

$$\left(\frac{d^2\sigma}{d\Omega dE}(\vartheta_{c.m.},E)\right)_L^{calc.} : \text{DWBA cross section (unit cross section)}$$

$$a_L(E) : \text{EWSR fraction}$$

a. **ISGR (L<15)+ IVGDR (through Coulomb excitation)**

b. **DWBA formalism; single folding ⇒ transition potential**

$$\delta U(r,E) = \int d\vec{r}\,'\,\delta\rho_L(\vec{r}\,',E)[V(|\vec{r}-\vec{r}\,'|,\rho_0(r')) + \rho_0(r')\frac{\partial V(|\vec{r}-\vec{r}\,'|,\rho(r'))}{\partial\rho_0(r')}]$$

$$U(r) = \int d\vec{r}\,'V(|\vec{r}-\vec{r}\,'|,\rho_0(r'))\rho_0(r')$$

Transition density

- **ISGMR Satchler, Nucl. Phys. A472 (1987) 215**

$$\delta\rho_0(r,E) = -\alpha_0[3 + r\frac{d}{dr}]\rho_0(r)$$

$$\alpha_0^2 = \frac{2\pi\hbar^2}{mA < r^2 > E}$$

- **ISGDR Harakeh & Dieperink, Phys. Rev. C23 (1981) 2329**

$$\delta\rho_1(r,E) = -\frac{\beta_1}{R\sqrt{3}}[3r^2 + \frac{d}{dr} + 10r - \frac{5}{3} < r^2 > \frac{d}{dr} + \varepsilon(r\frac{d^2}{dr^2} + 4\frac{d}{dr})]\rho_0(r)$$

$$\beta_1^2 = \frac{6\pi\hbar^2}{mAE}\frac{R^2}{(11 < r^4 > -(25/3) < r^2 >^2 - 10\varepsilon < r^2 >)}$$

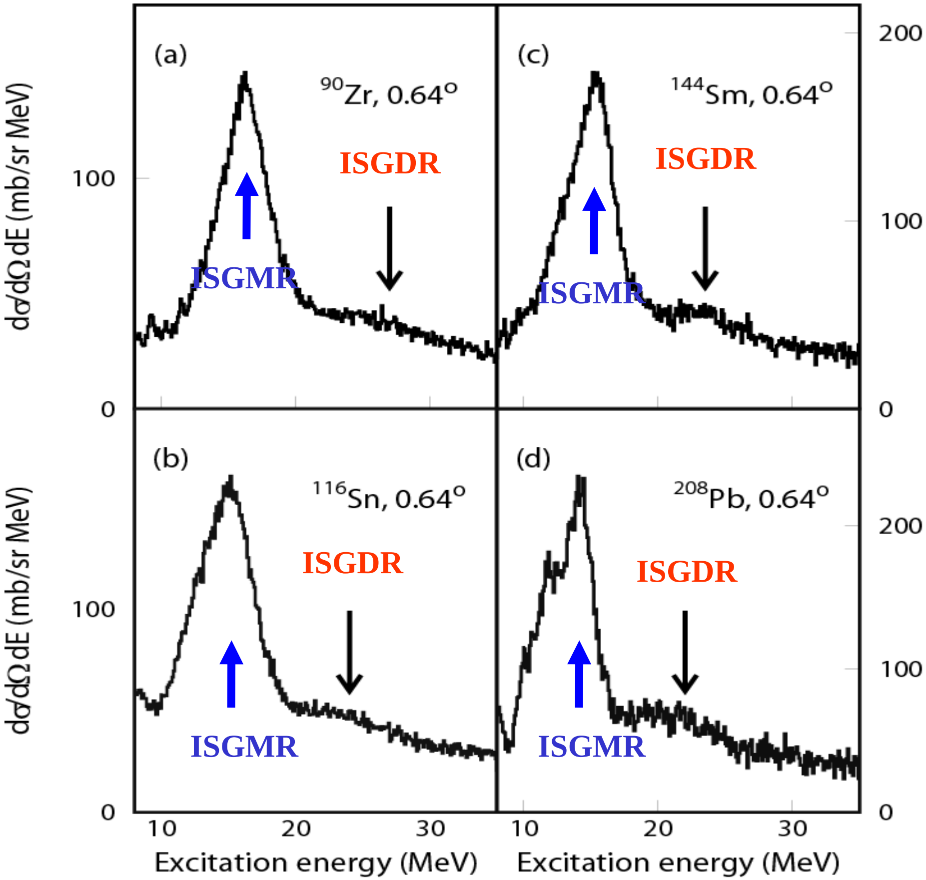
- **Other modes Bohr-Mottelson (BM) model**

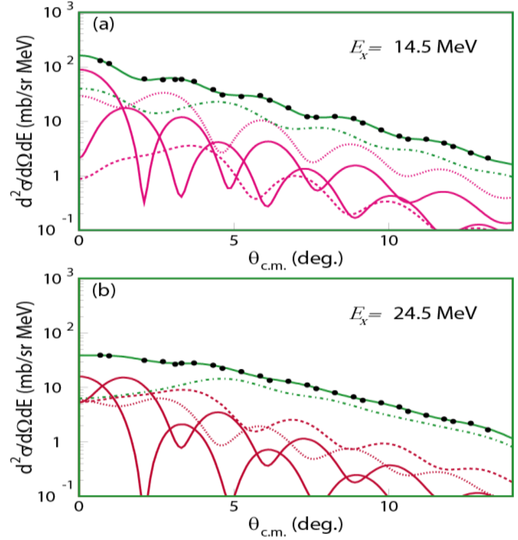
$$\delta\rho_L(r,E) = -\delta_L\frac{d}{dr}\rho_0(r)$$

$$\delta_L^2 = (\beta_L c)^2 = \frac{l(2l+1)^2}{(l+2)^2}\frac{2\pi\hbar^2}{mAE}\frac{< r^{2l-2} >}{< r^{l-2} >^2}$$

Uchida *et al.*,
Phys. Lett. B557 (2003) 12
Phys. Rev. C69 (2004) 051301

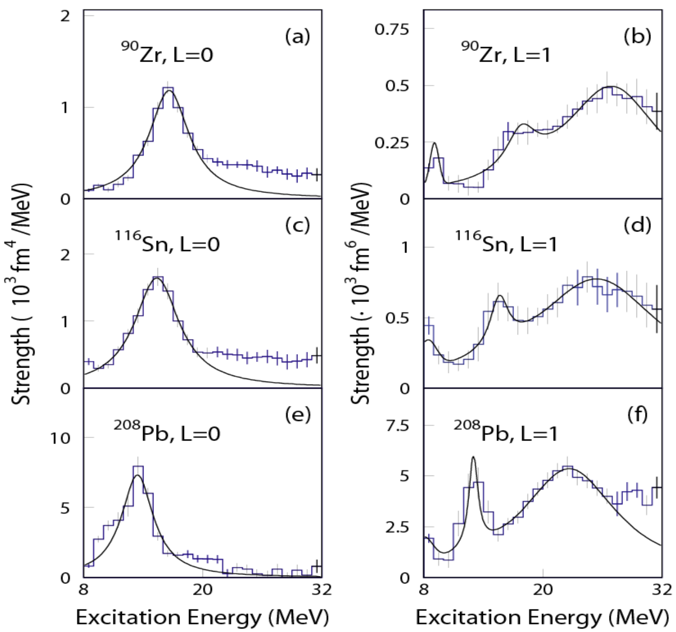
(α,α′) spectra at 386 MeV

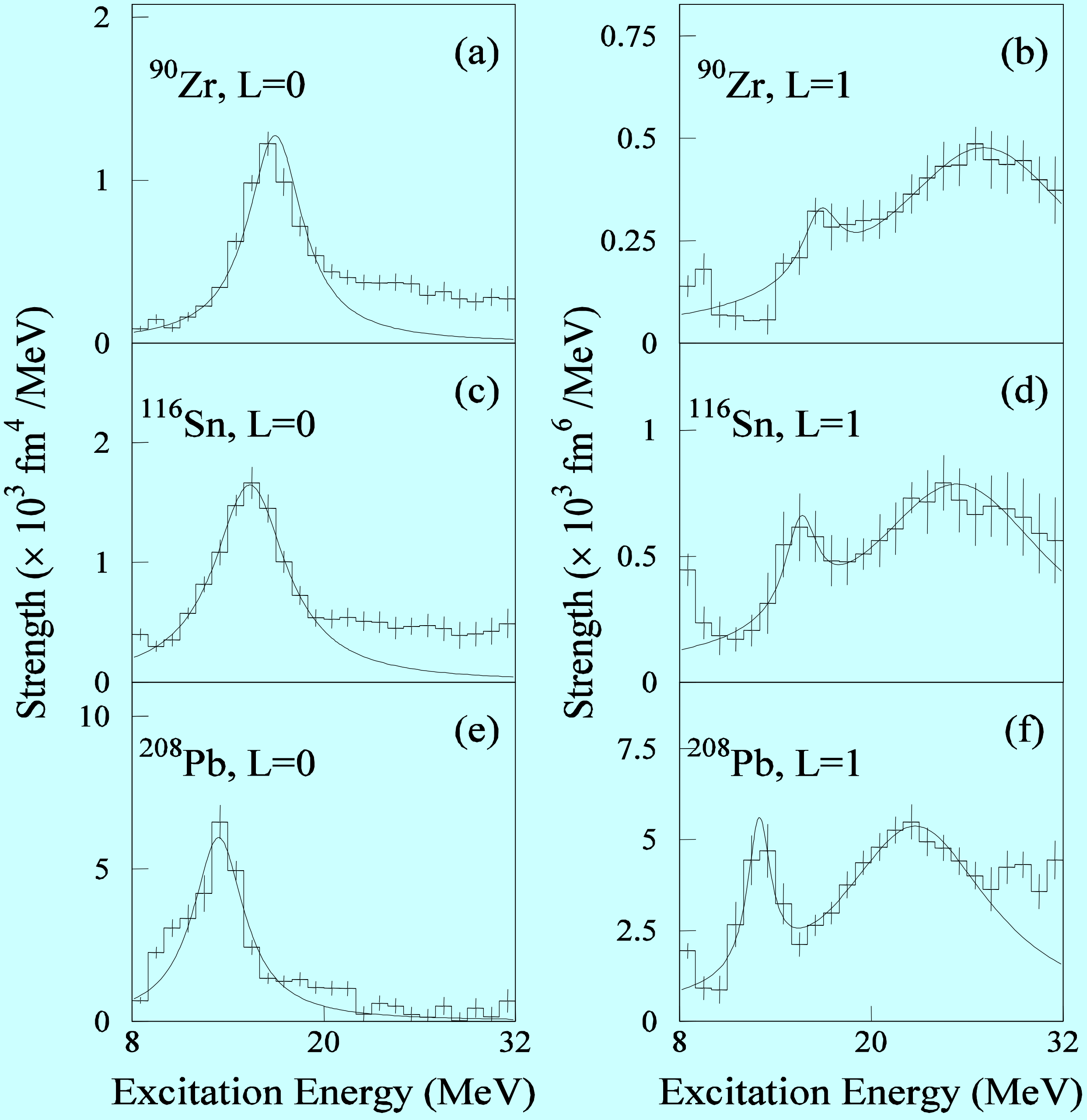




^{116}Sn

MDA results for L=0 and L=1




(a)
^{90}Zr, L=0
(b)
^{90}Zr, L=1
(c)
^{116}Sn, L=0
(d)
^{116}Sn, L=1
(e)
^{208}Pb, L=0
(f)
^{208}Pb, L=1
Strength (× 10^3 fm^4 /MeV)
Strength (× 10^3 fm^6 /MeV)
Excitation Energy (MeV)
Excitation Energy (MeV)

In HF+RPA calculations,

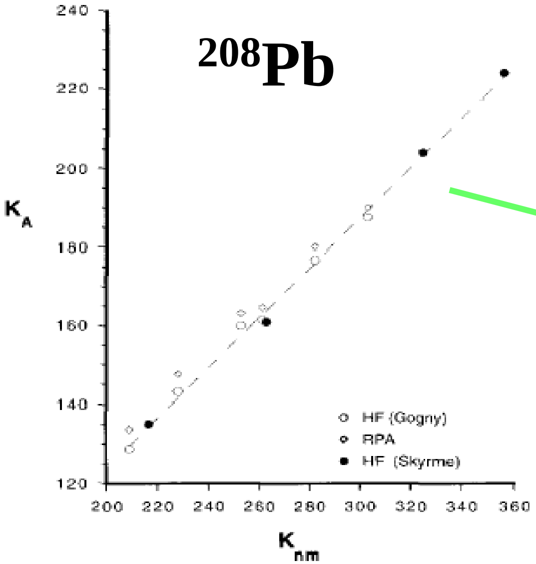
$$K_{nm} = \left[9\rho^2 \frac{d^2(E/A)}{d\rho^2} \right]_{\rho=\rho_0}$$

Nuclear matter

***E/A*: binding energy per nucleon** ***K_A*: incompressibility**

***ρ* : nuclear density**

***ρ_0* : nuclear density at saturation**



K_A is obtained from excitation energy of ISGMR & ISGDR

$K_A = 0.64 K_{nm} - 3.5$

J.P. Blaizot, NPA591 (1995) 435

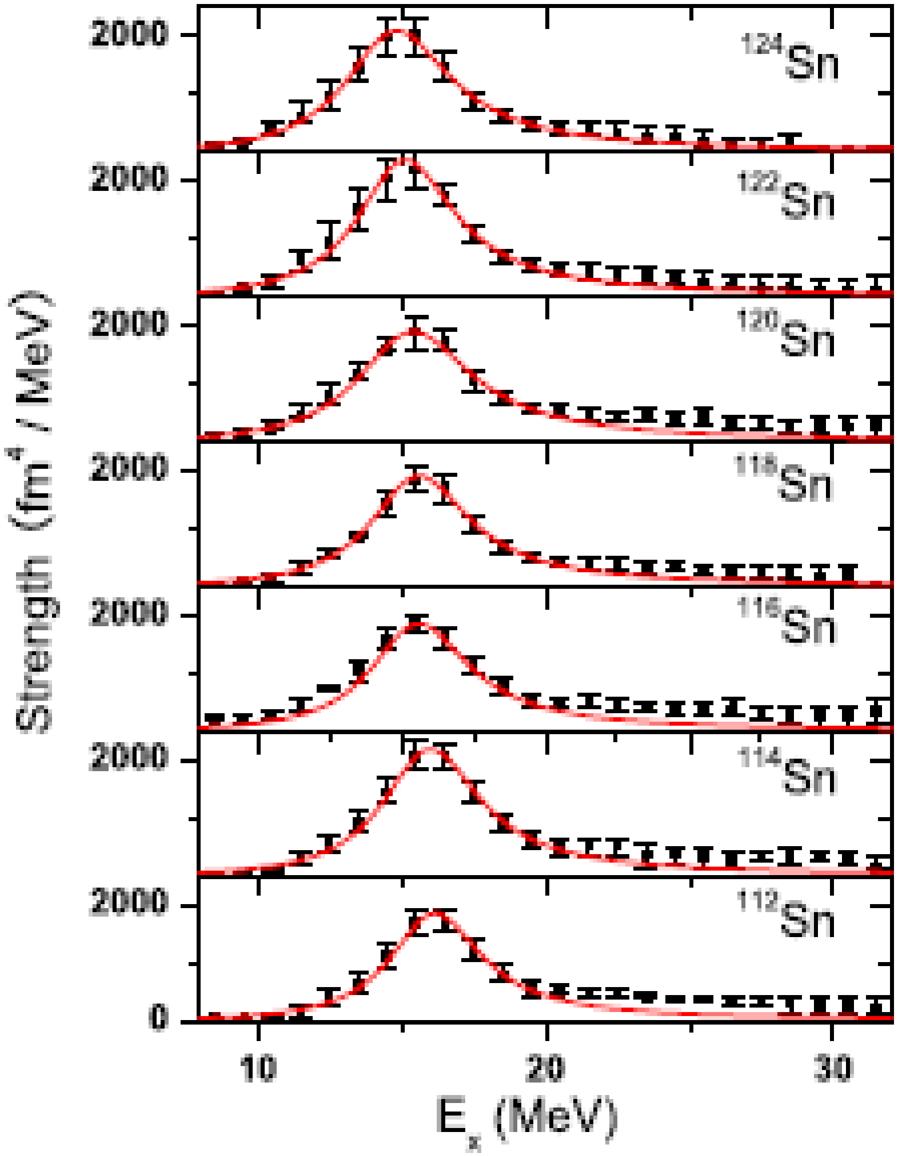
From GMR data on ^{208}Pb and ^{90}Zr,

$$K_\infty = 240 \pm 10 \text{ MeV}$$

[See, *e.g.*, G. Colò *et al.*, Phys. Rev. C 70 (2004) 024307]

This number is consistent
with both ISGMR and ISGDR Data
and
with non-relativistic and relativistic calculations

Isoscalar GMR strength distribution in Sn-isotopes obtained by Multipole Decomposition Analysis of singles spectra obtained in ^{A}Sn(α,α′) measurements at incident energy 400 MeV and angles from 0º to 9º



$$K_A \sim K_{vol}(1 + cA^{-1/3}) + K_\tau((N - Z)/A)^2 + K_{Coul} Z^2 A^{-4/3}$$

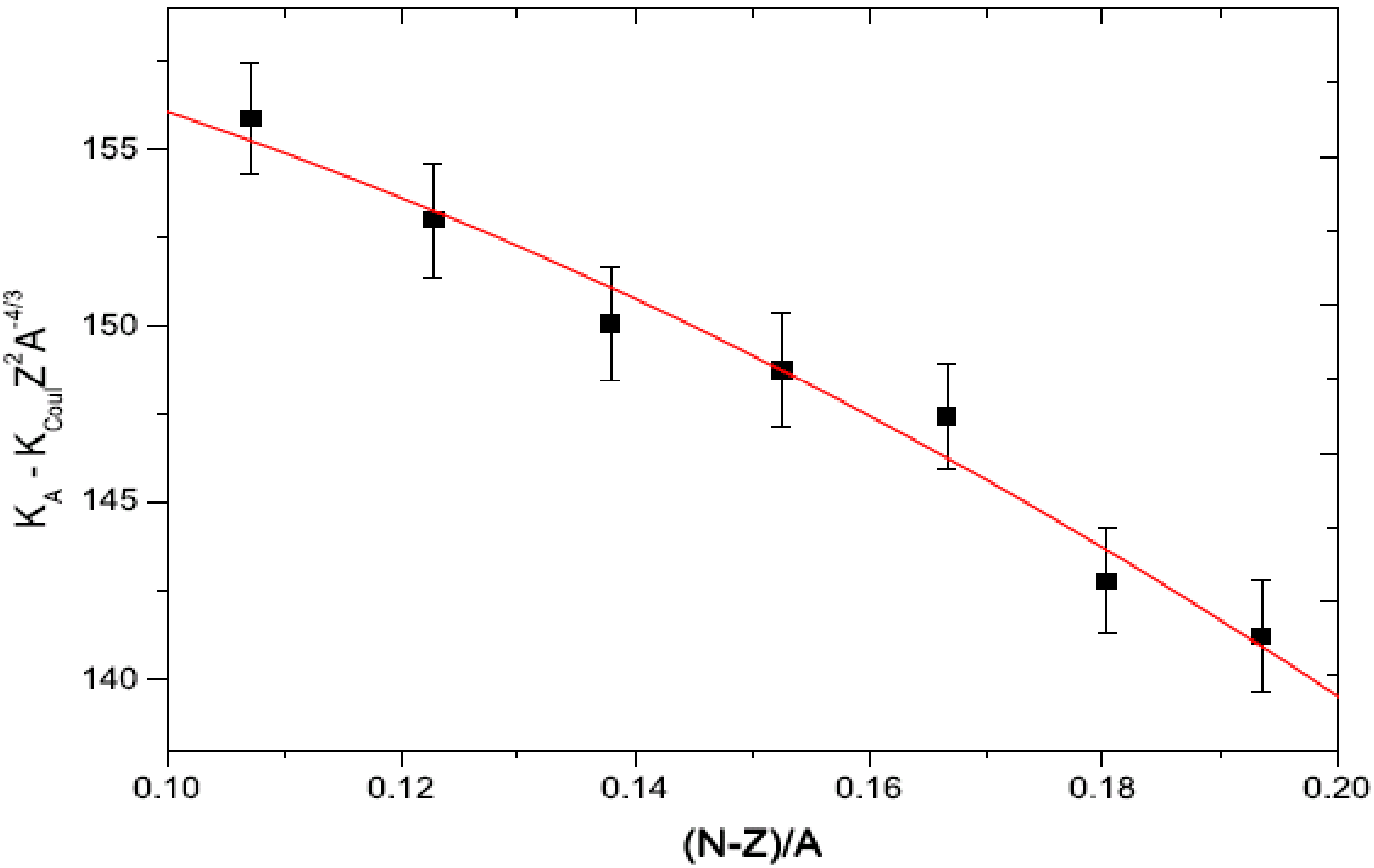
$$K_A - K_{Coul} Z^2 A^{-4/3} \sim K_{vol}(1 + cA^{-1/3}) + K_\tau((N - Z)/A)^2$$

$$\sim \text{Constant} + K_\tau((N - Z)/A)^2$$

We use K_{Coul} = - 5.2 MeV (from Sagawa)

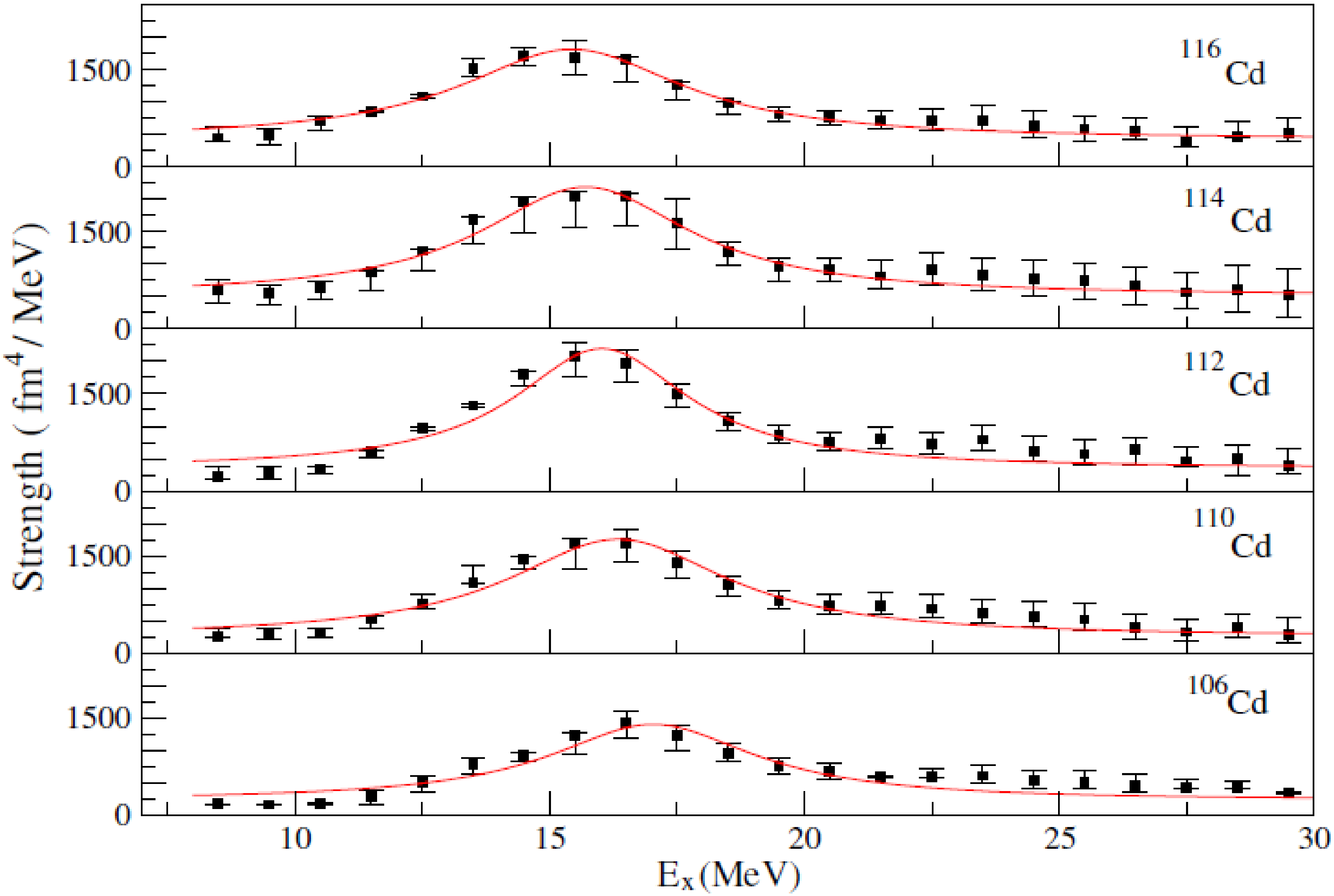
N-Z/A

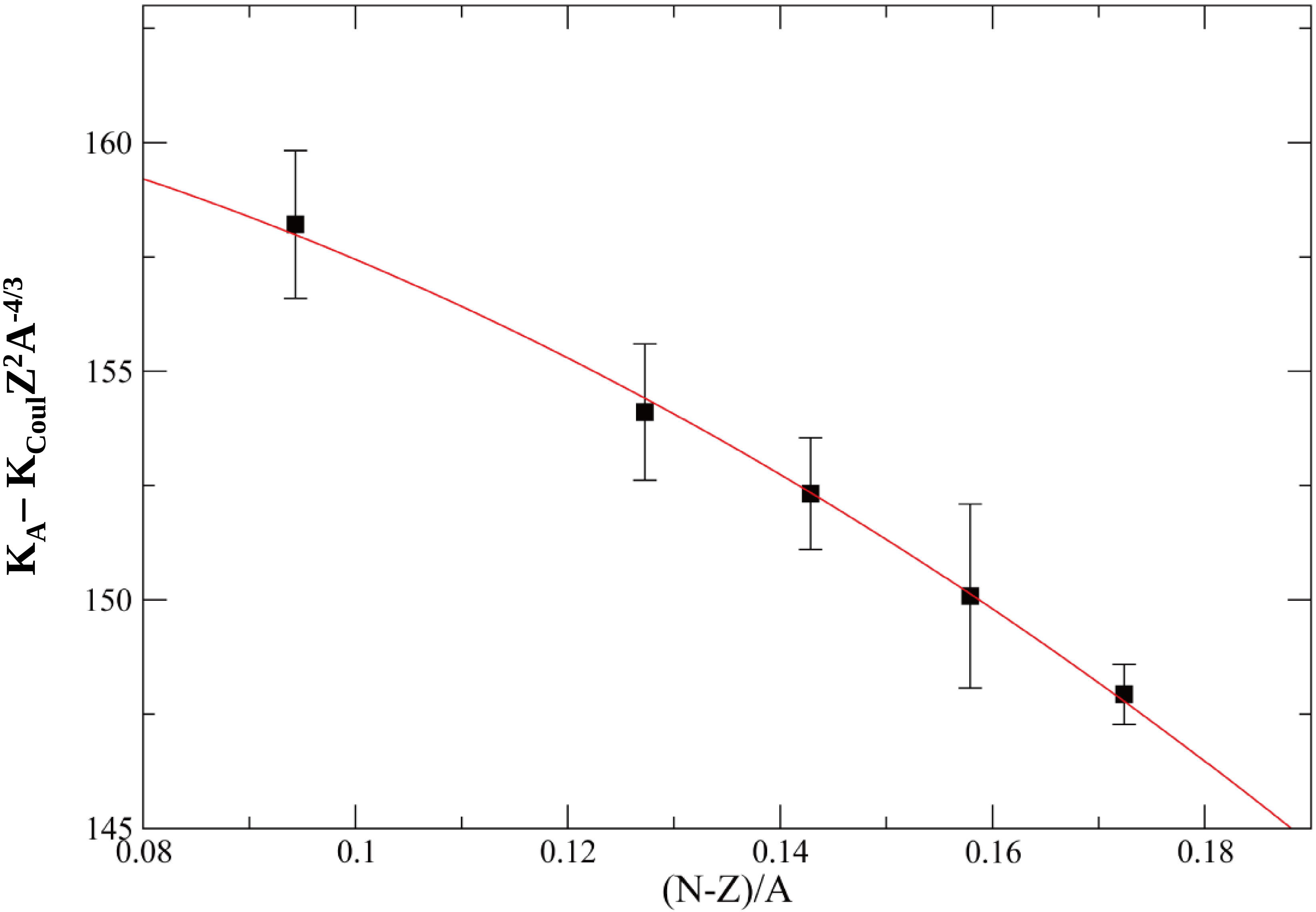
^{112}Sn – ^{124}Sn: 0.107 – 0.194



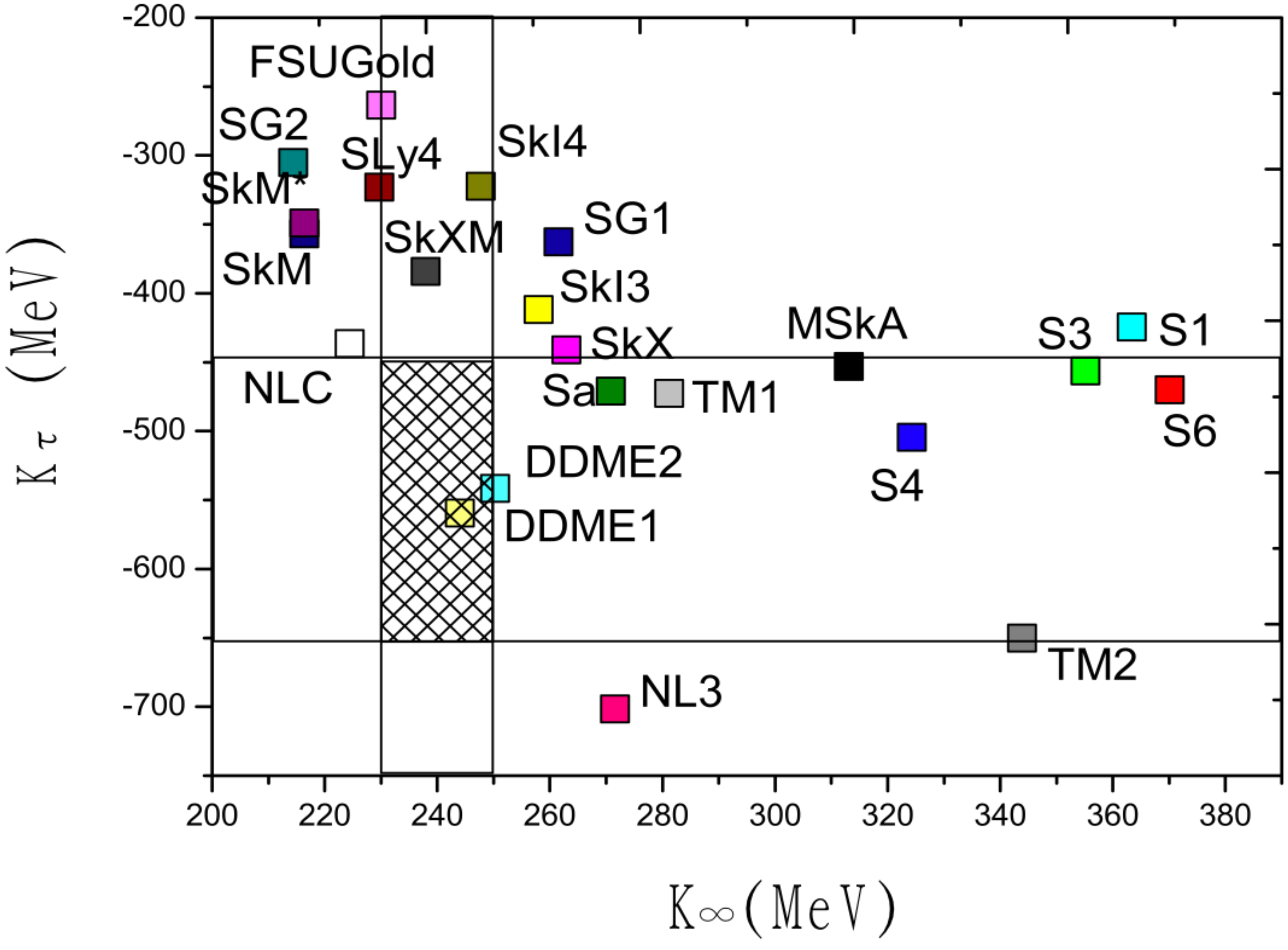
$$K_\tau = -550 \pm 100 \text{ MeV}$$

Monopole strength Distribution

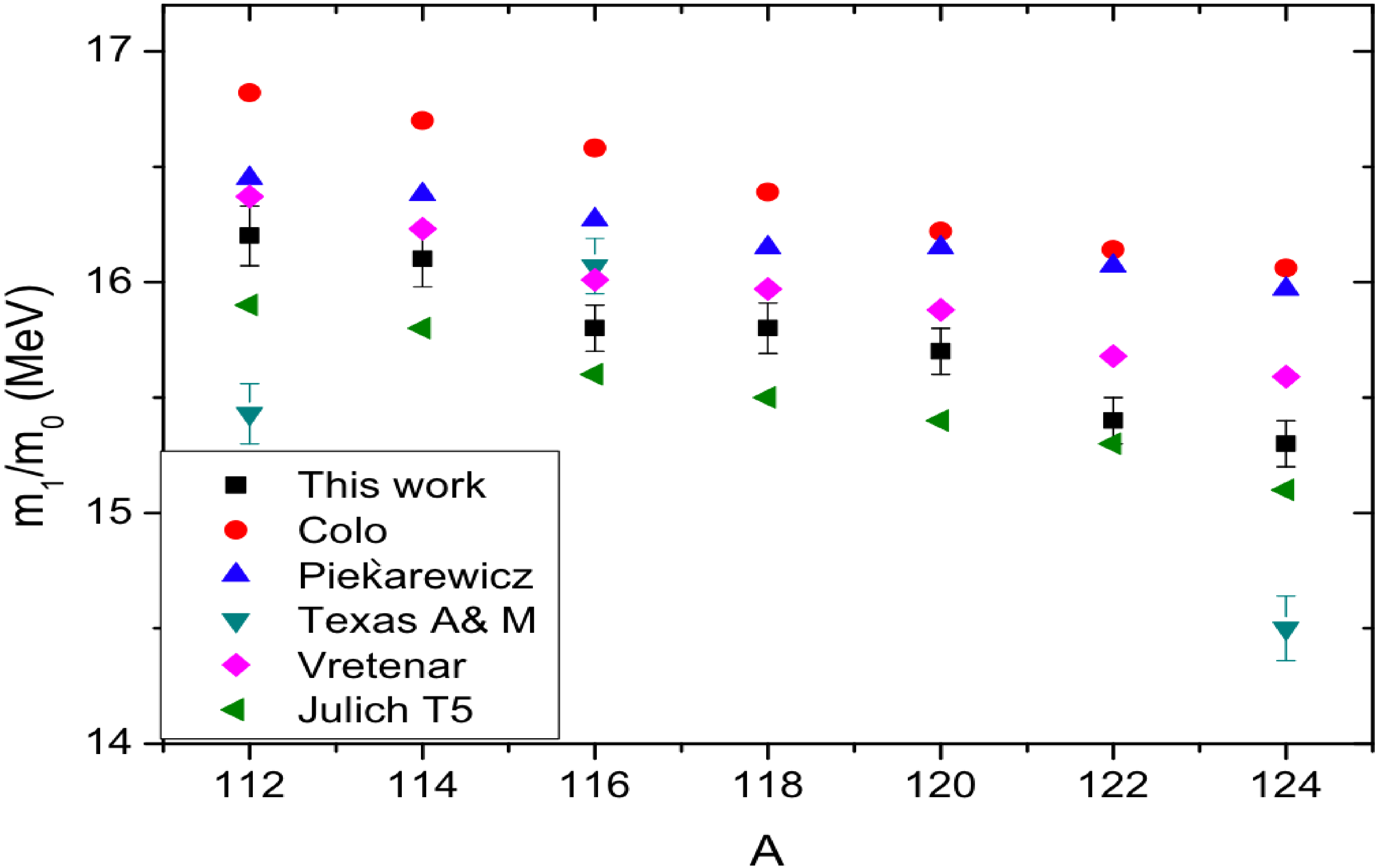




$$K_\tau = -480 \pm 100\ \mathrm{MeV}$$



Data from H. Sagawa *et al.*, Phys. Rev. C **76**, 034327 (2007)



17
16
15
14
m_1/m_0 (MeV)
112
114
116
118
120
122
124
A
This work
Colo
Piekarewicz
Texas A& M
Vretenar
Julich T5

Colò *et al*.: Non-relativistic RPA (without pairing) reproduces ISGMR in ^{208}Pb and ^{90}Zr.

Piekarewicz: Relativistic RPA (FSUGold model) reproduces g.s. observables and ISGMR in ^{208}Pb, ^{144}Sm and ^{90}Zr

Vretenar: Relativistic mean field (DD-ME2: density-dependent mean-field effective interaction).

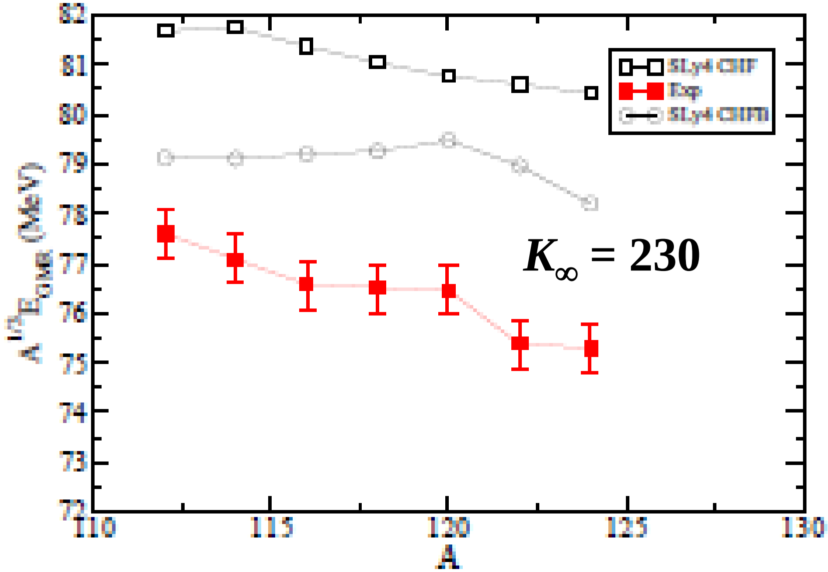
Possibly agreement is fortuitous since strength distributions are not much different from those by Colò *et al*. and Piekarewicz.

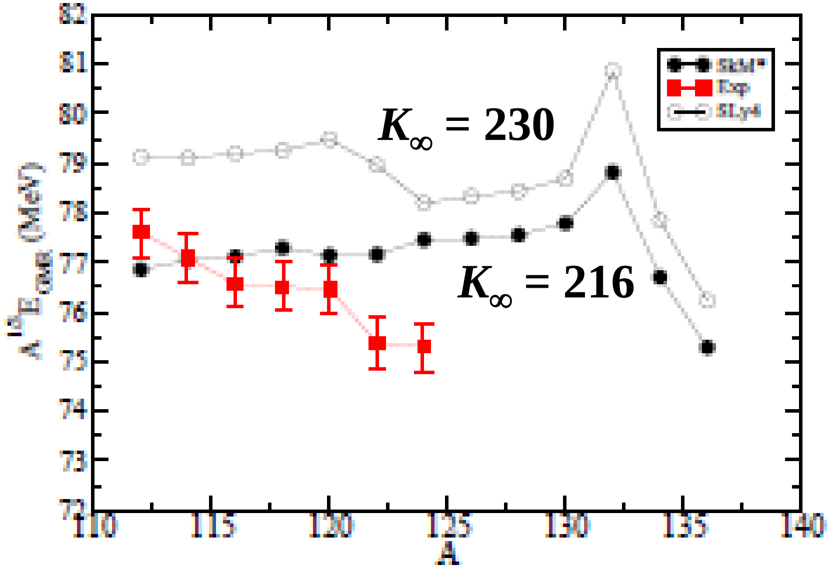
Tselyaev *et al*.: Quasi-particle time-blocking approximation (QTBA) (T5 Skyrme interaction)

K_∞ = 202 MeV?!

Softness of Sn-nuclei is still unresolved

E. Khan, PRC 80, 011307(R) (2009)





The Giant Monopole Resonances in Pb isotopes

E. Khan, Phys. Rev. C 80, 057302 (2009).

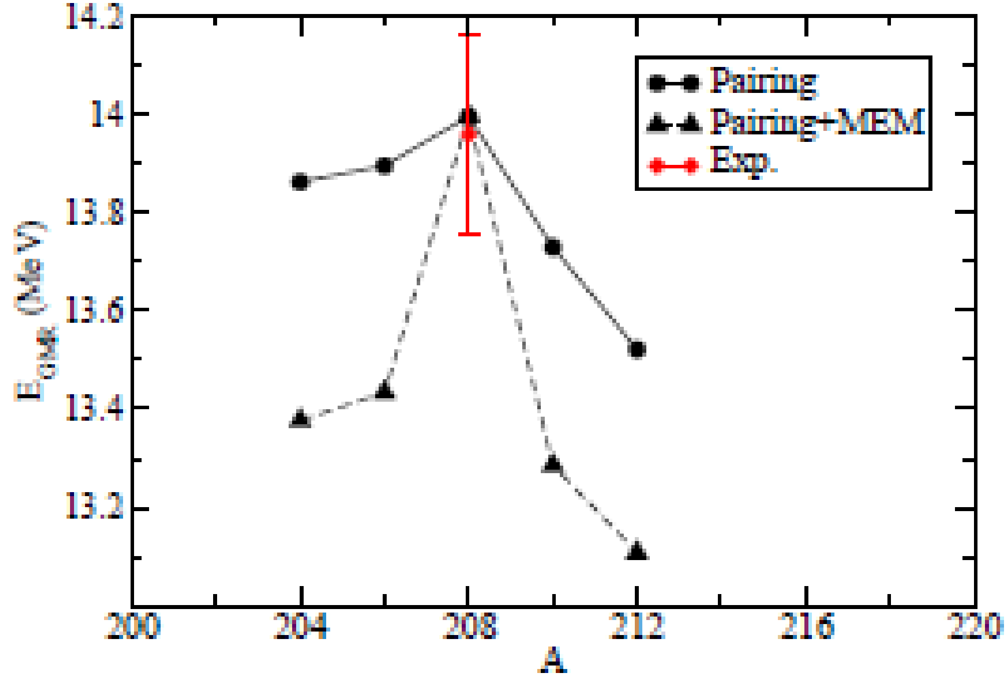
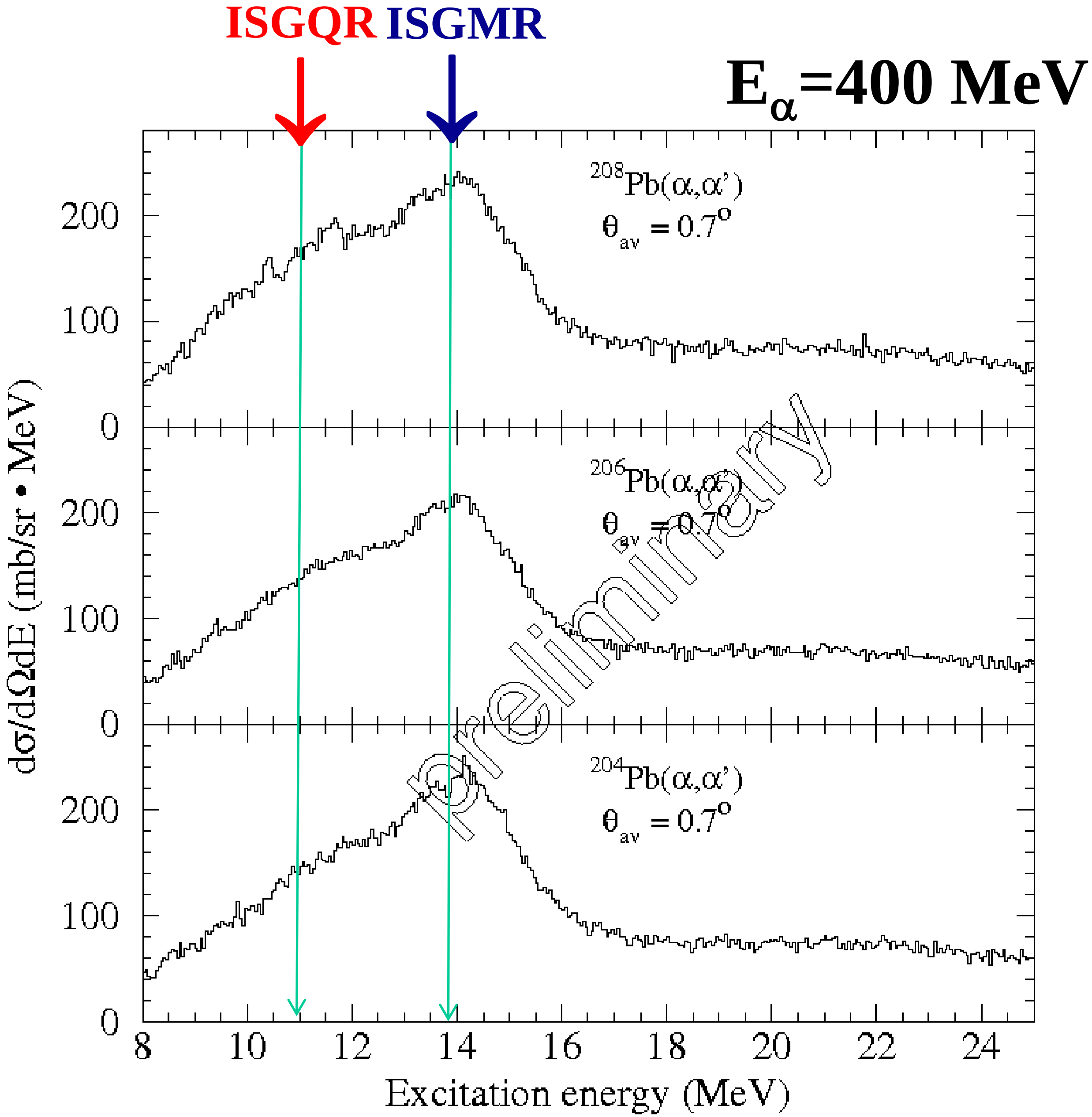


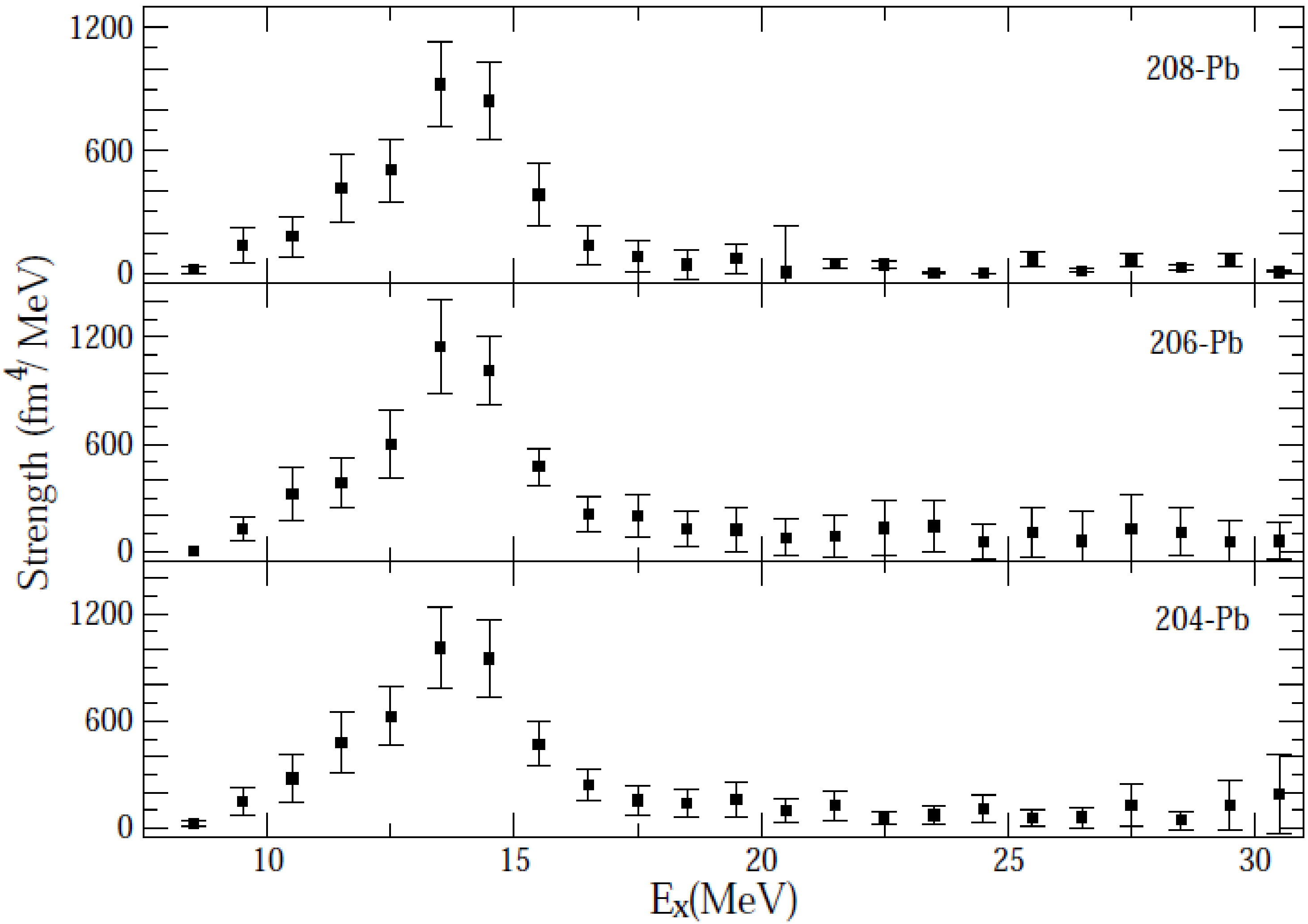
FIG. 2: Excitations energies of the GMR in $^{204-212}$Pb isotopes calculated with constrained HFB method, taking into account the MEM effect (see text). The experimental data is taken from Ref. [22]

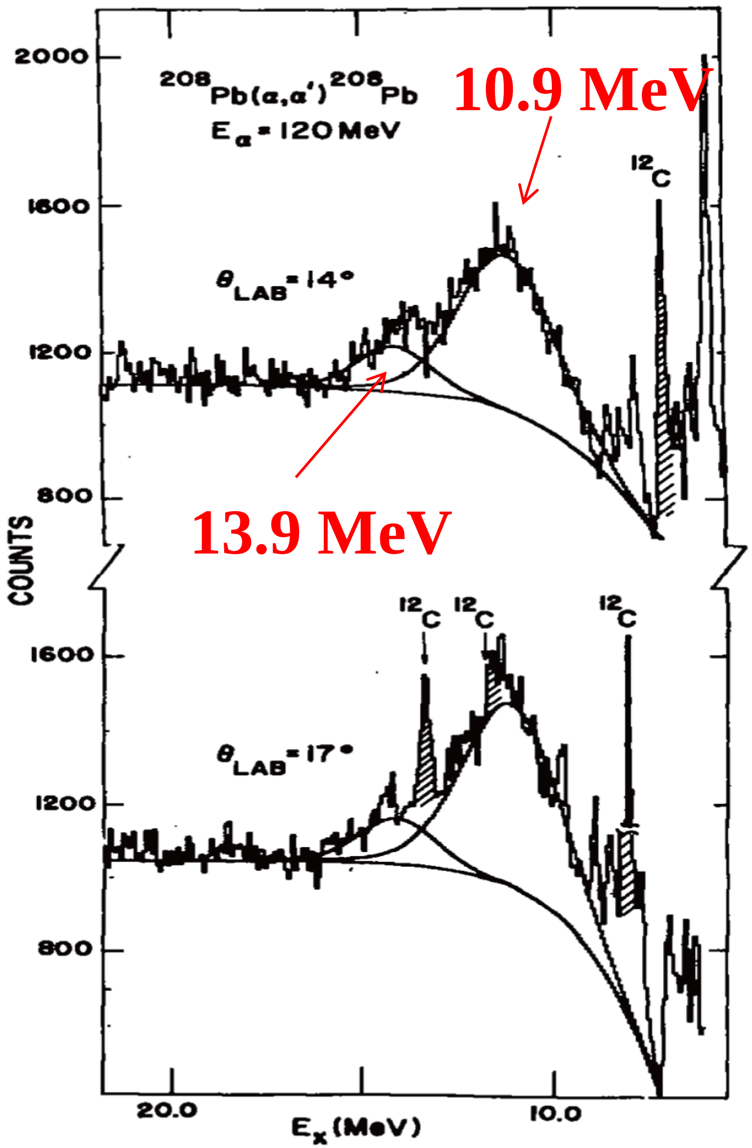
Mutually Enhanced Magicity (MEM)?

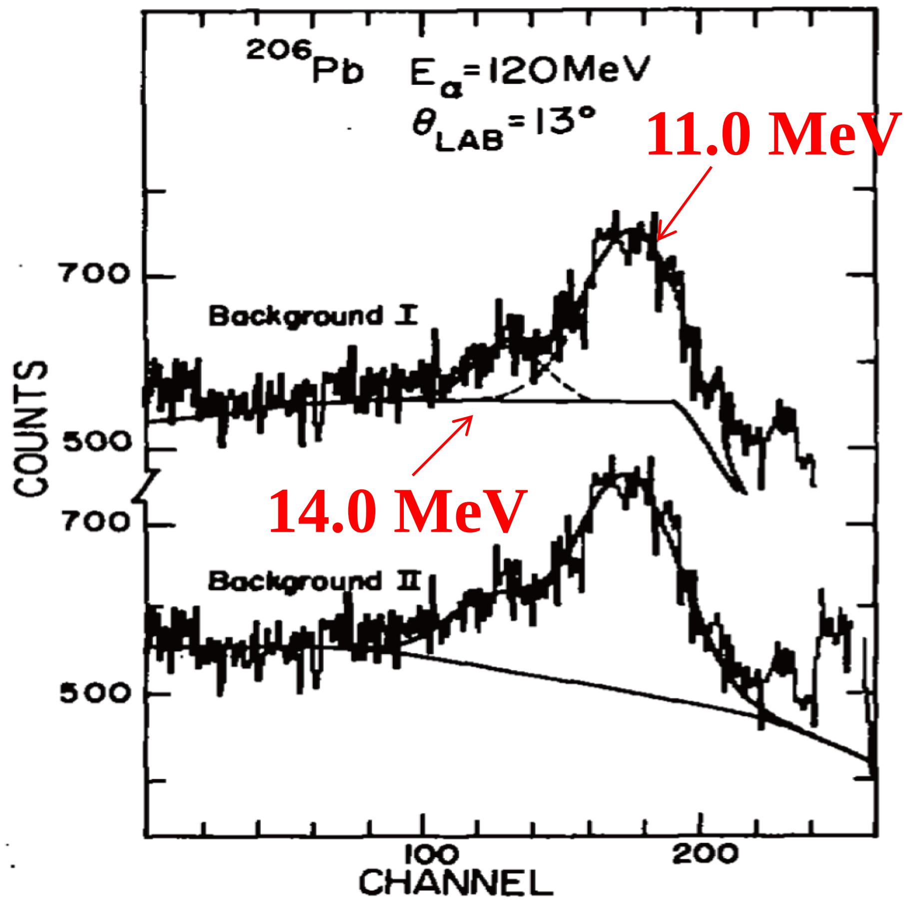


208	**13.9**	**MeV**
206	**13.94**	**MeV**
204	**13.98**	**MeV**

Monopole strength Distribution







M.N. Harakeh *et al.*, Nucl. Phys. A327, 373 (1979)

Decay of giant resonances

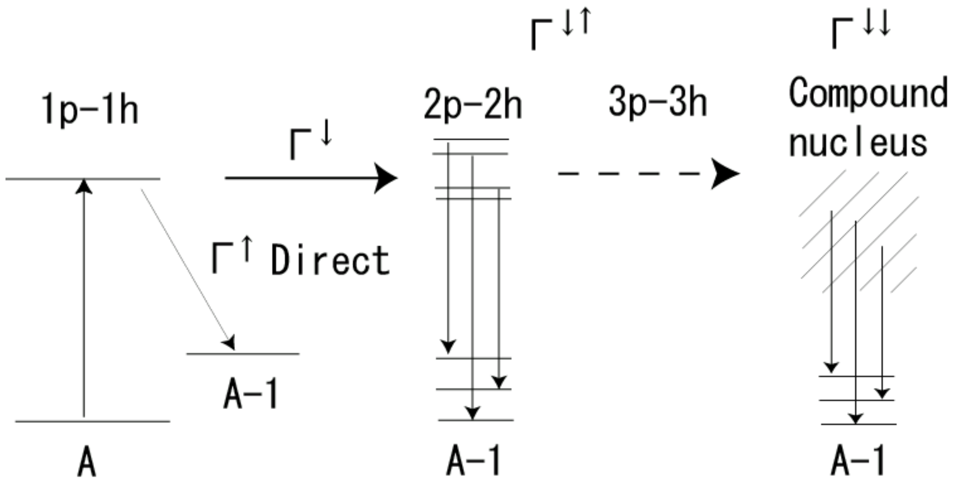
- **Width of resonance**

  $\Gamma, \Gamma^{\uparrow}, \Gamma^{\downarrow}$ $(\Gamma^{\downarrow\uparrow}, \Gamma^{\downarrow\downarrow})$
  - $\Gamma^{\uparrow}$**: direct or escape width**
  - $\Gamma^{\downarrow}$ **: spreading width**

    $\Gamma^{\downarrow\uparrow}$**: pre-equilibrium,** $\Gamma^{\downarrow\downarrow}$**: compound**
- **Decay measurements**

  $\Rightarrow$ **Direct reflection of damping processes**

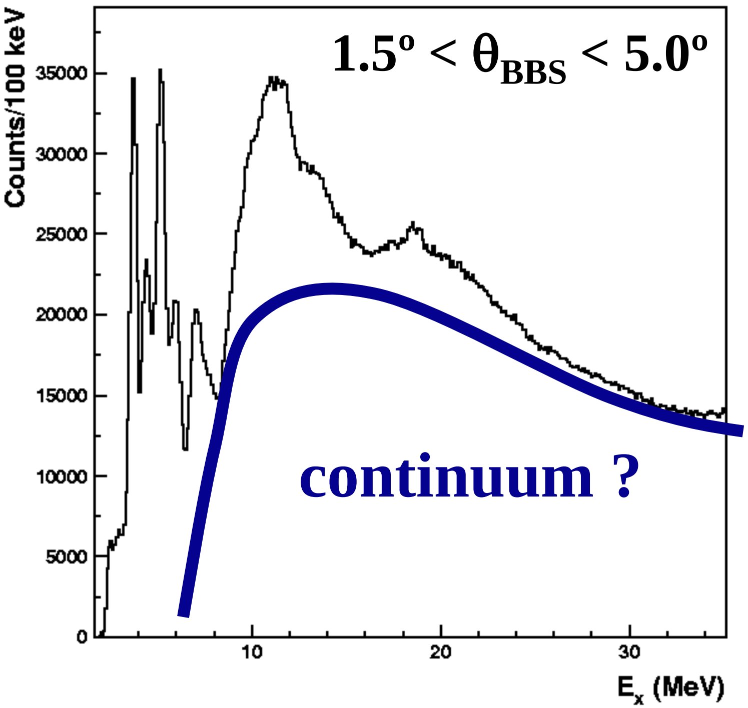
  Allows detailed comparison with theoretical calculations



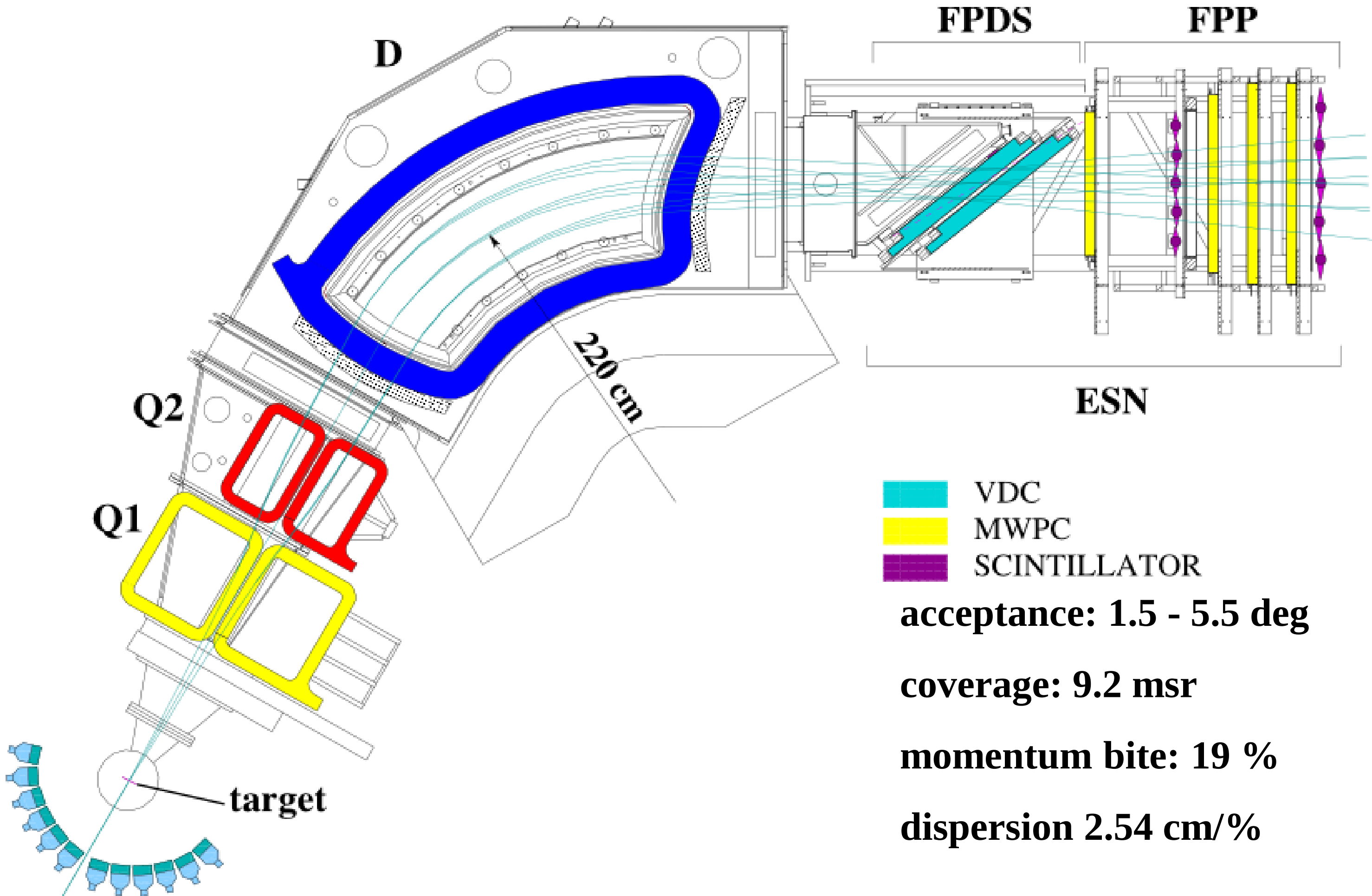
Excitation of ISGDR in ^{208}Pb

- In ^{208}Pb located around 22 MeV and width of 4 MeV
- *L=1* angular distribution peaks close to a scattering angle of 0^o
- Difficult to identify in nuclear continuum and rides on instrumental background

Singles ^{208}Pb(α,α') 200 MeV $\Rightarrow$



Si-ball
16 Si-detectors at
10 cm from the target
total solid angle: 1 sr



acceptance: 1.5 - 5.5 deg

coverage: 9.2 msr

momentum bite: 19 %

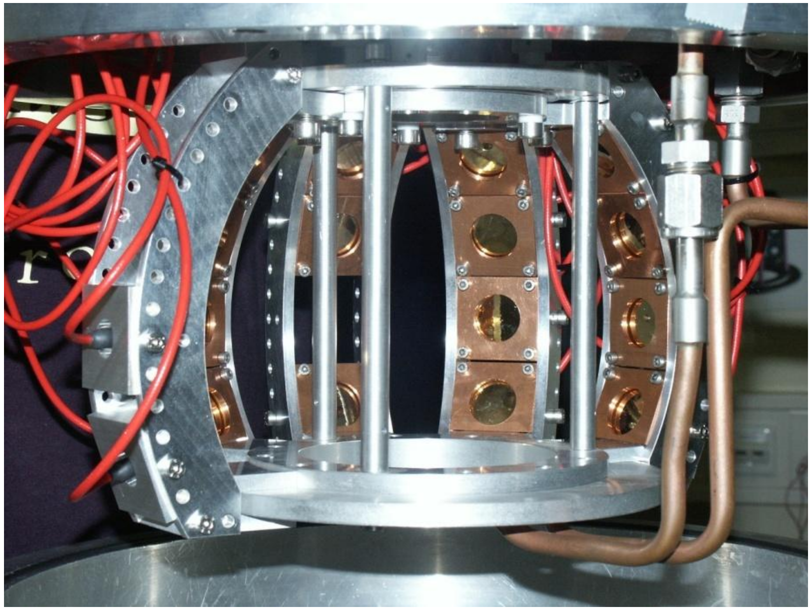
dispersion 2.54 cm/%

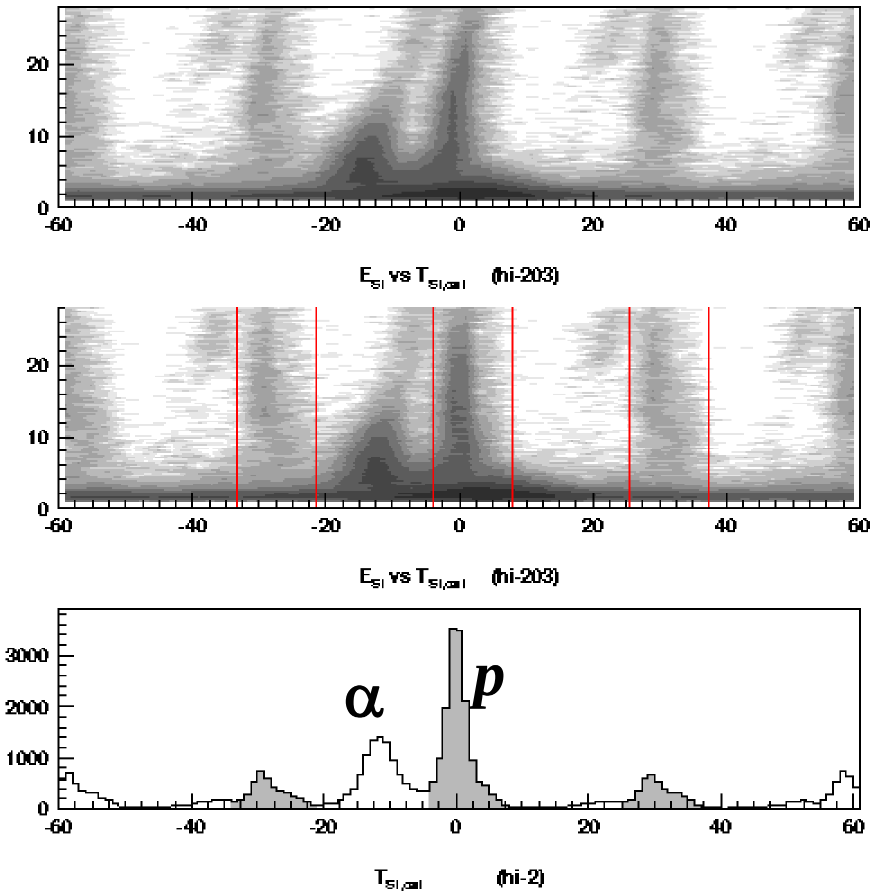
EDEN (32 detectors), NE213

total solid angle: 0.37 sr

KVI Big-Bite Spectrometer (BBS)

Proton-decay detection





α-*p* separation using rise time of signal SiLi

Microscopic structure of ISGDR

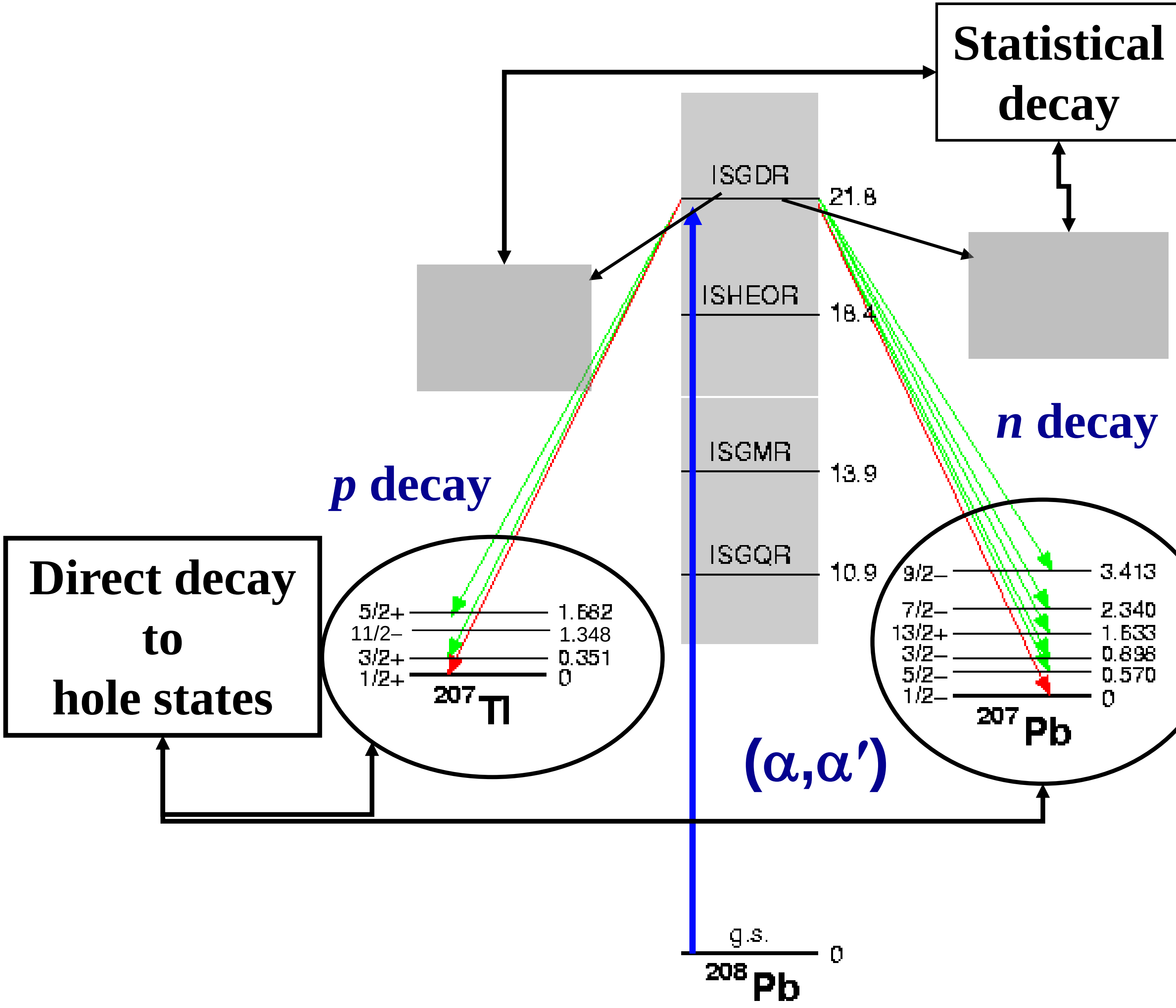
Transition operator

$$O^{L=1} = \sum_i r_i^1 Y_0^1 + \frac{1}{2}\sum_i r_i^3 Y_0^1 + \ ...$$

Spurious center of mass motion

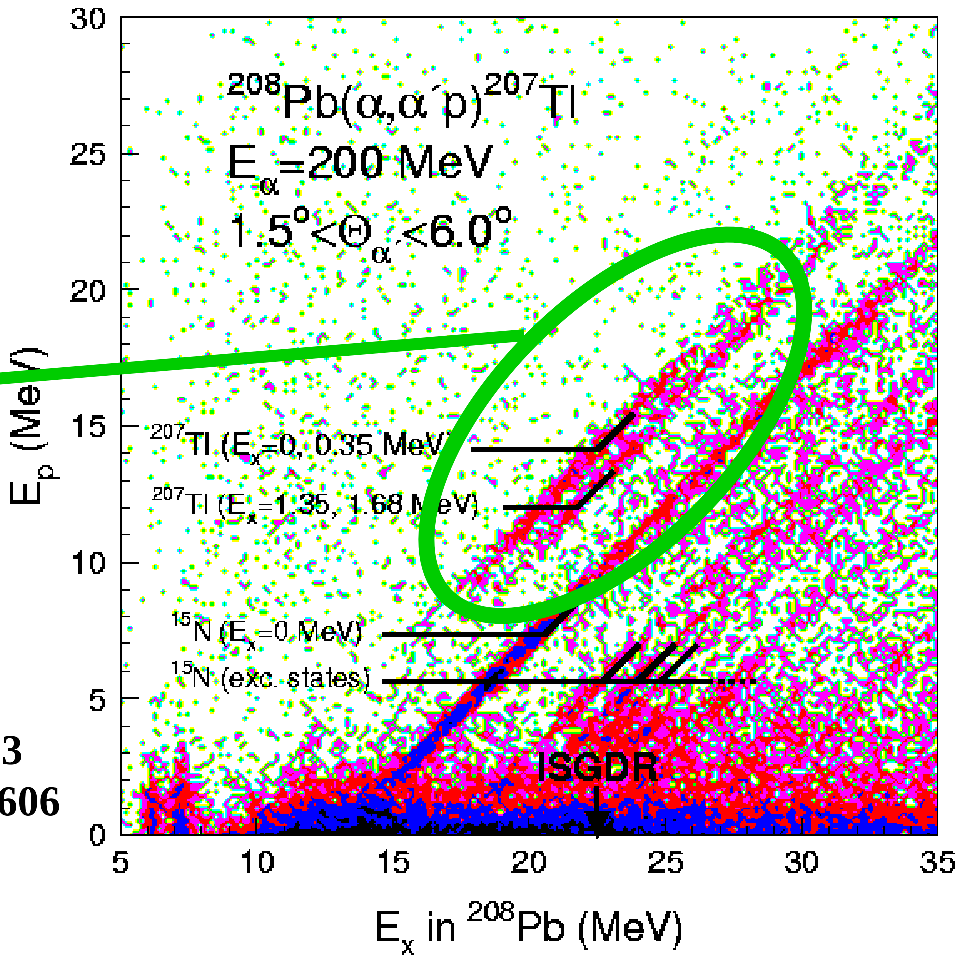
Overtone

3ħω excitation (overtone of c.o.m. motion)


Statistical decay
ISGDR
21.8
ISHEOR
18.4
ISGMR
13.9
ISGQR
10.9
n decay
p decay
Direct decay to hole states
5/2+ 1.682
11/2– 1.348
3/2+ 0.351
1/2+ 0
207Tl
9/2– 3.413
7/2– 2.340
13/2+ 1.633
3/2– 0.898
5/2– 0.570
1/2– 0
207Pb
(α,α′)
g.s. 0
208Pb

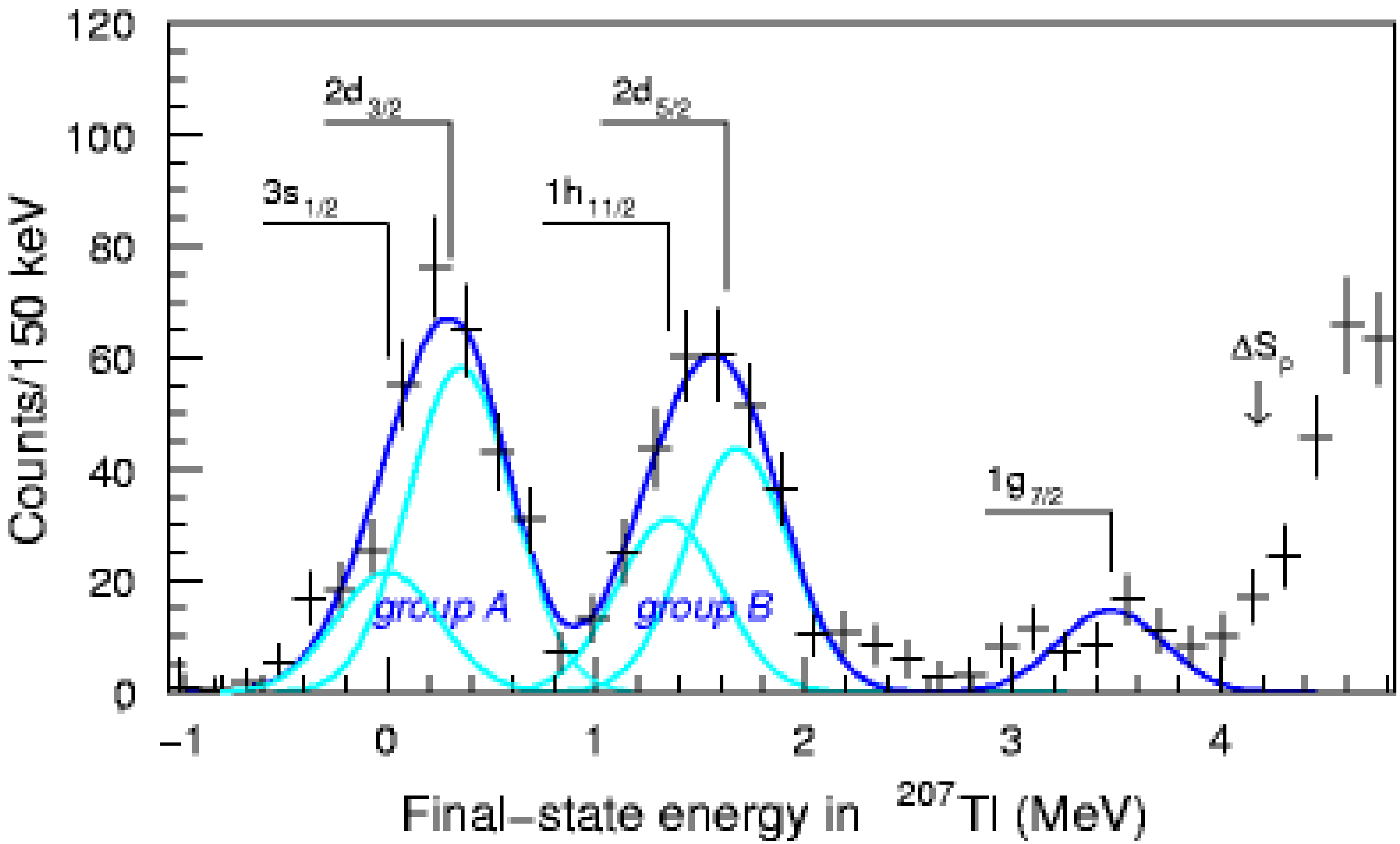
^{208}Pb(α,α′) followed by *p* decay

Decay to hole states in ^{207}Tl; branching ratios predicted by Gorelik *et al.*



M. Hunyadi *et al*., Phys. Lett. B576 (2003) 253
M. Hunyadi *et al*., Phys. Rev. C75 (2007) 014606

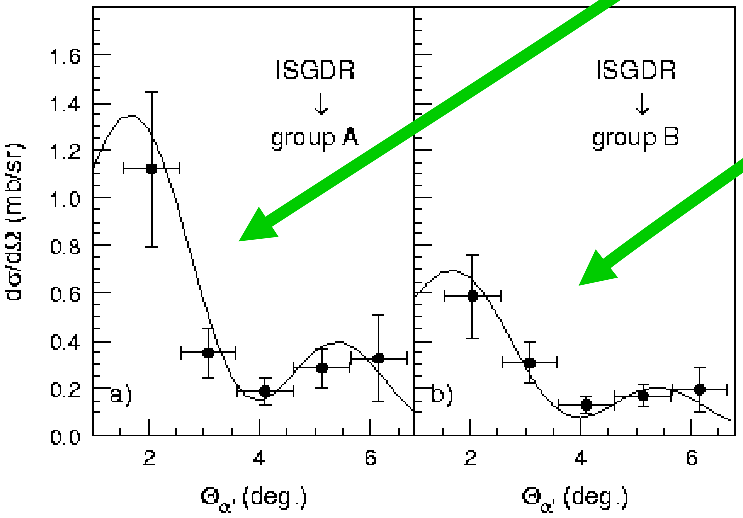
Branching ratios for decay

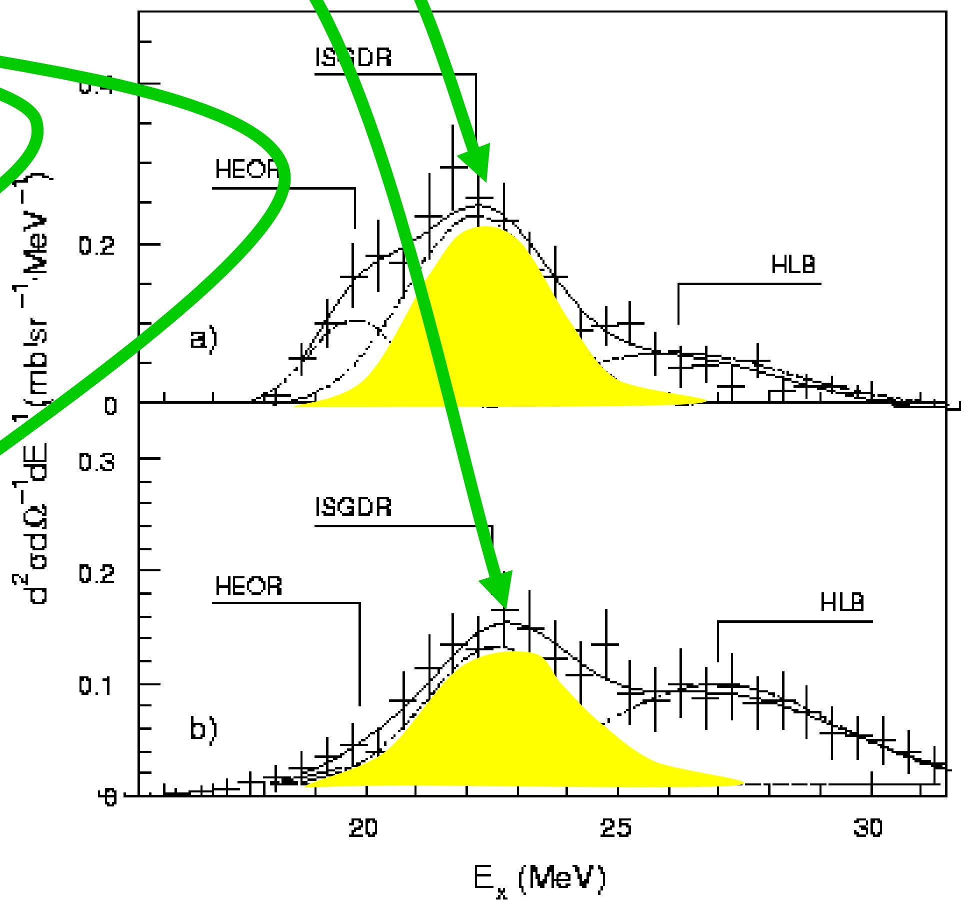


ISGDR in ^{208}Pb in p decay

E_x = 22.1 ± 0.3 MeV
***L = 1* transition**

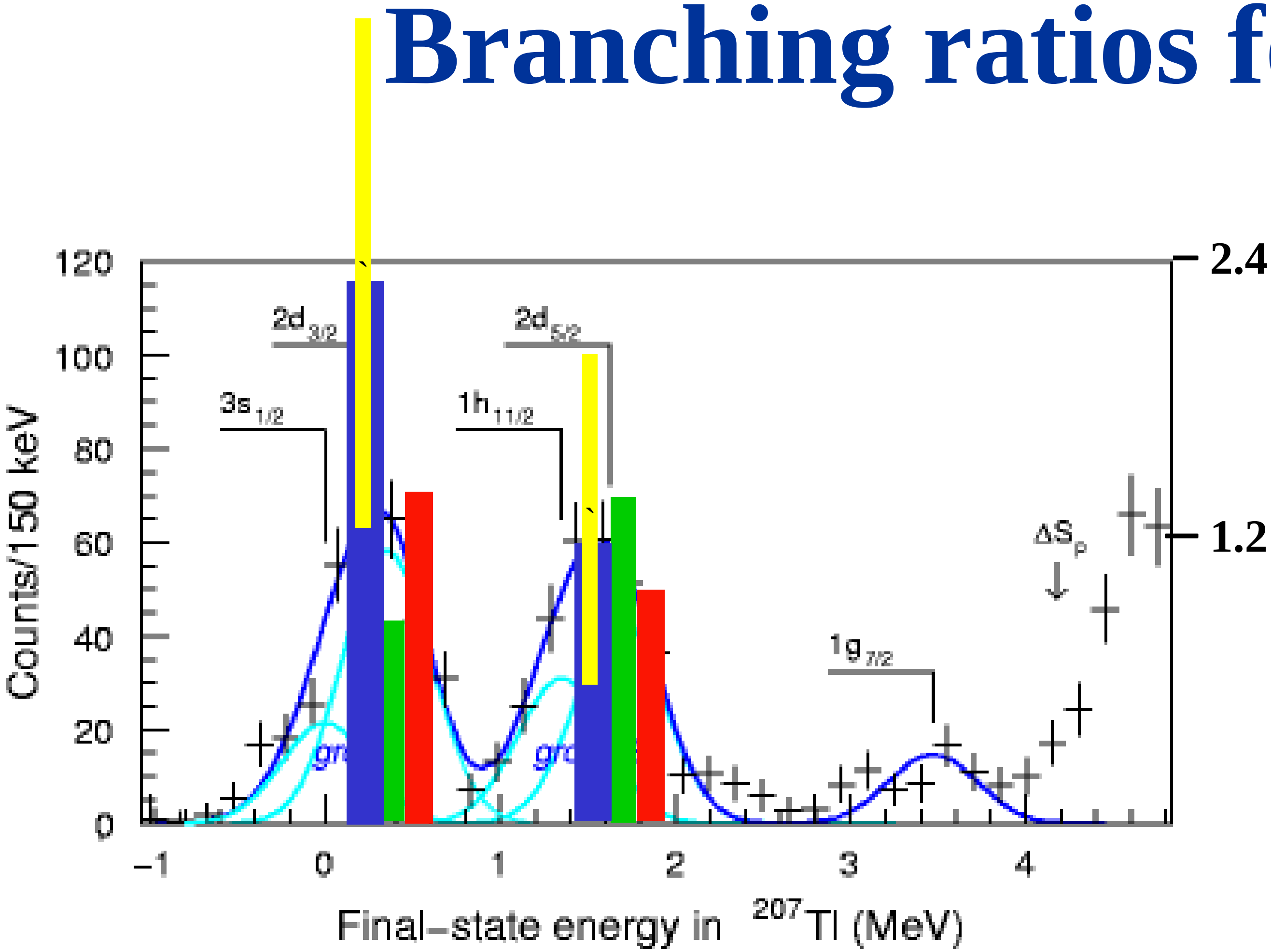
gated on hole states





KVI GSI university of groningen

Branching ratios for decay



This work

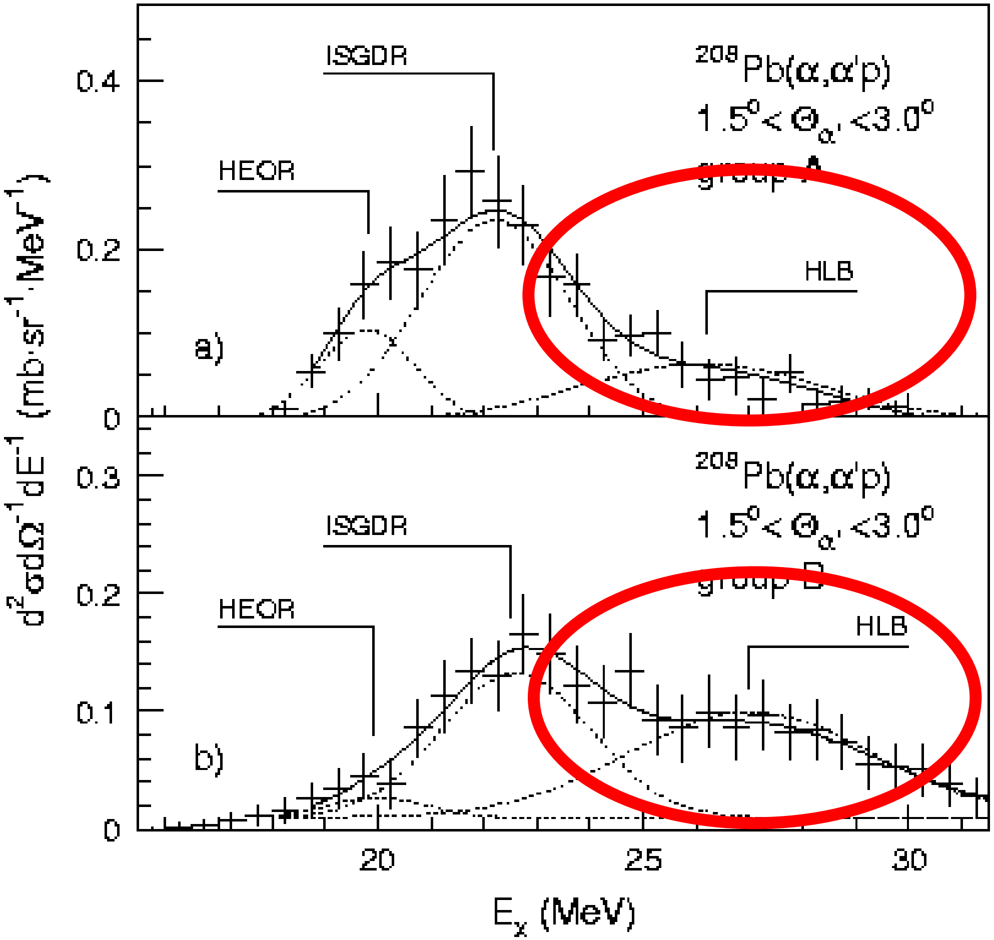
Gorelik *et al.*, PRC 62 (2000) 047301;

Continuum RPA;
Landau-Migdal
Parameters: f^{ex}, f';
Smearing parameter
Δ energy-dependent

Gorelik *et al.*, PRC 69 (2004) 054322

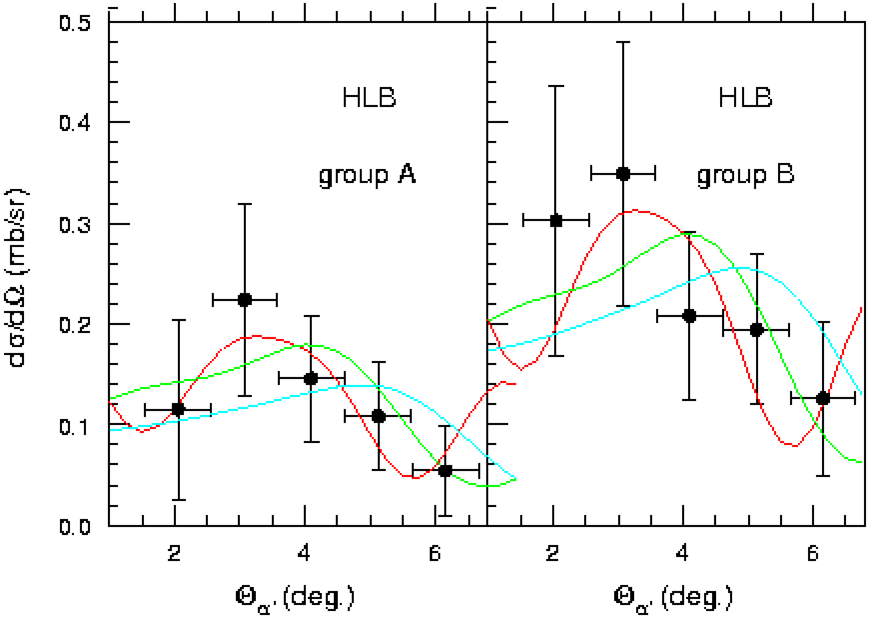
$1/2^+ + 3/2^+$	2.6%	(S~0.56) 1.45%	2.3 ± 1.1
$11/2^- + 5/2^+$	1.9%	(S~0.56) 1.04%	1.2 ± 0.7

Overtone of the ISGQR? [r^4Y_2]

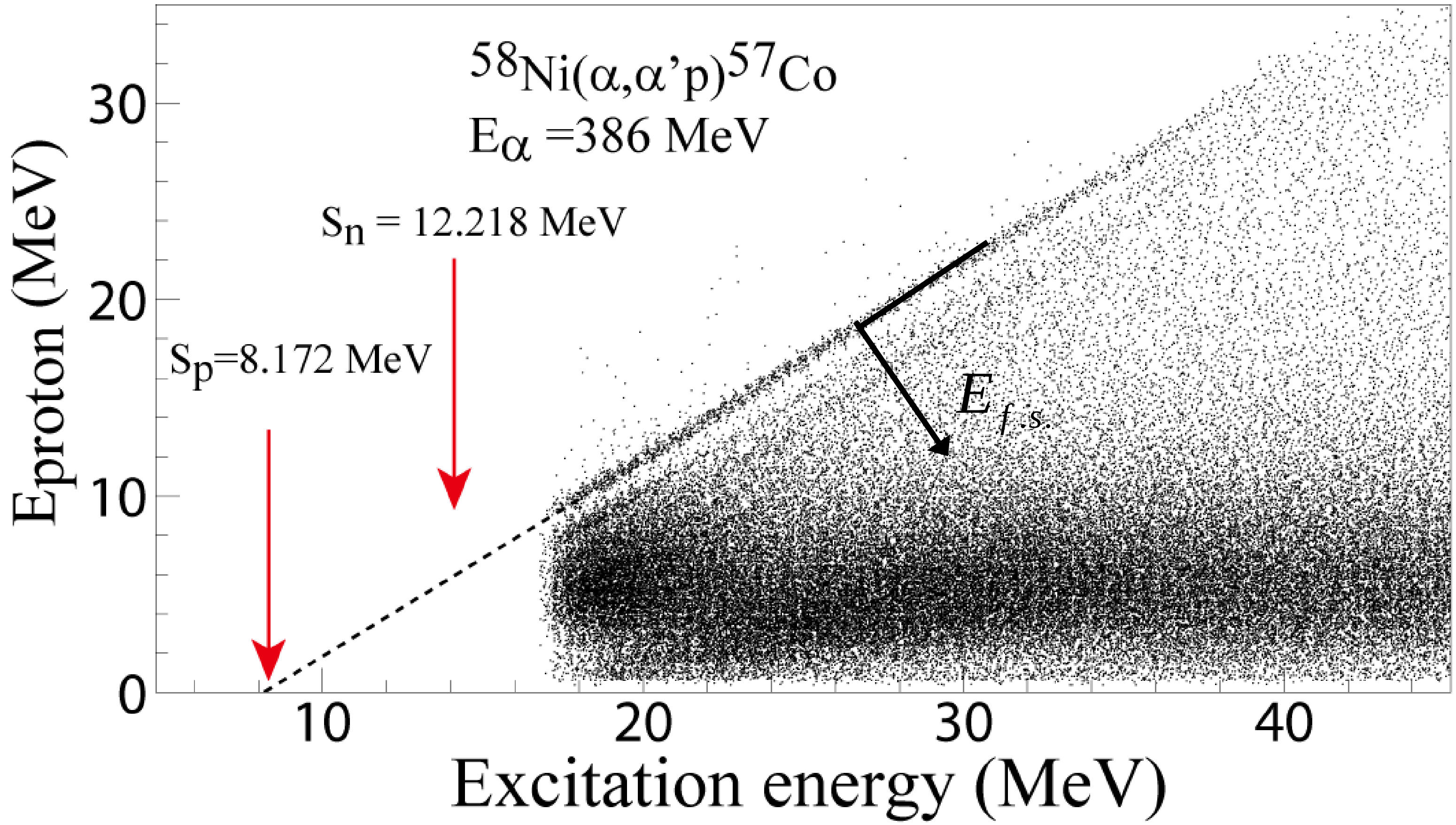


E_x = 26.9 ± 0.7 MeV

Muraviev and Urin
Bull. Acad. Sci. USSR
Phys. Ser. 52 (1988) 123
E_x = 28.3 MeV



Data analysis: Proton decay



Final-state energy

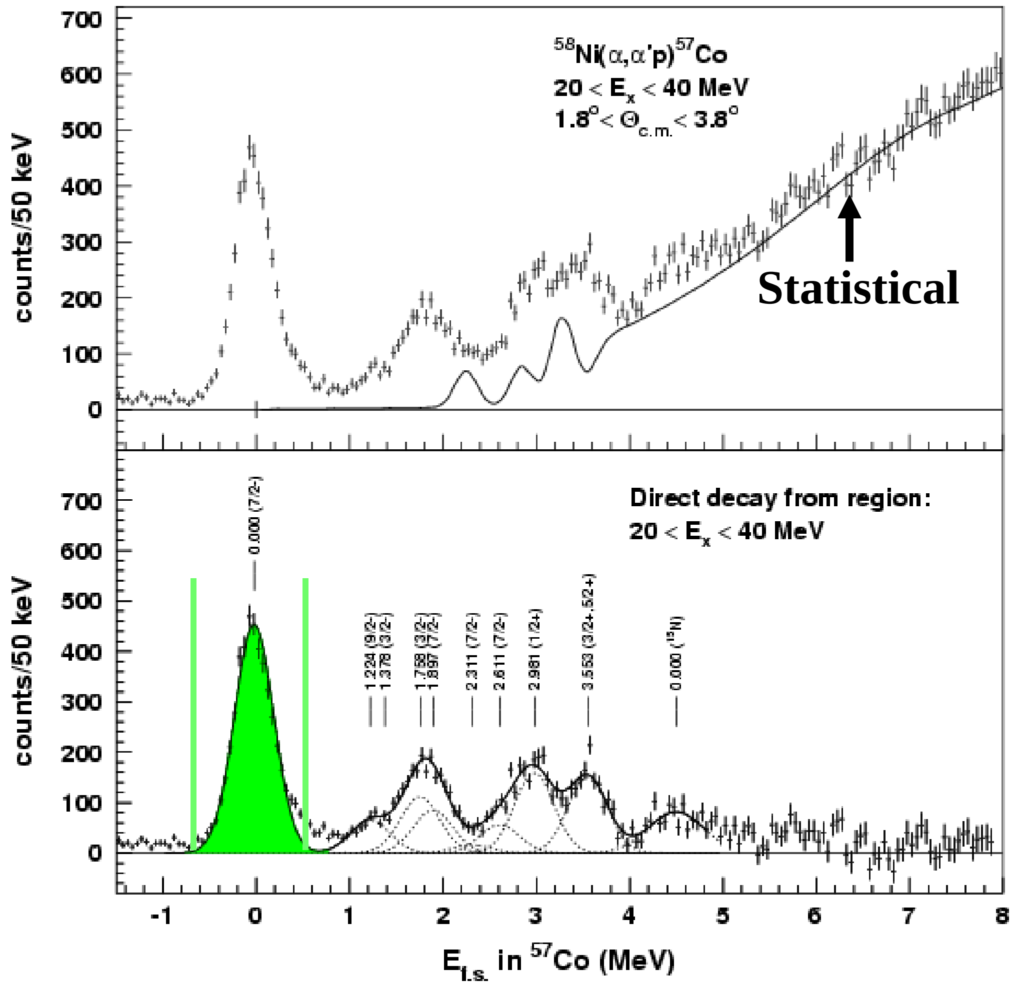
$$E_{f.s.} = E_X - E_p - S_p$$

Experimental results

M. Hunyadi *et al.*, Phys. Rev. C80 (2009) 044317

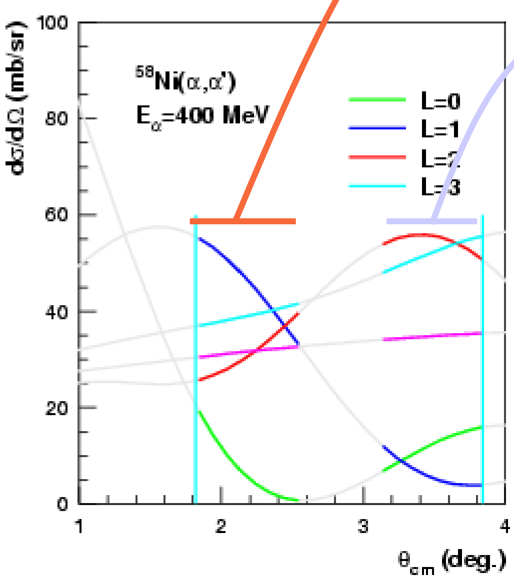
Final-state energy spectra

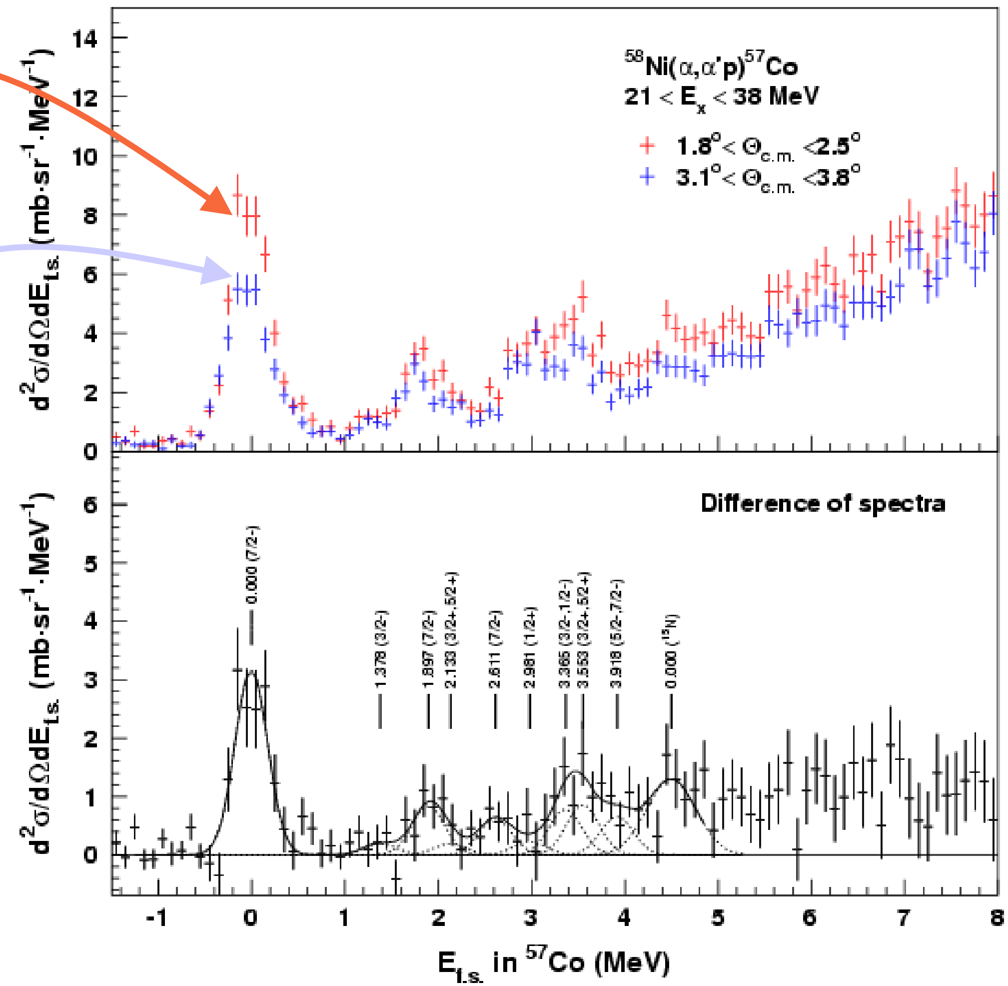
Final-state energy spectra after subtracting statistical contribution



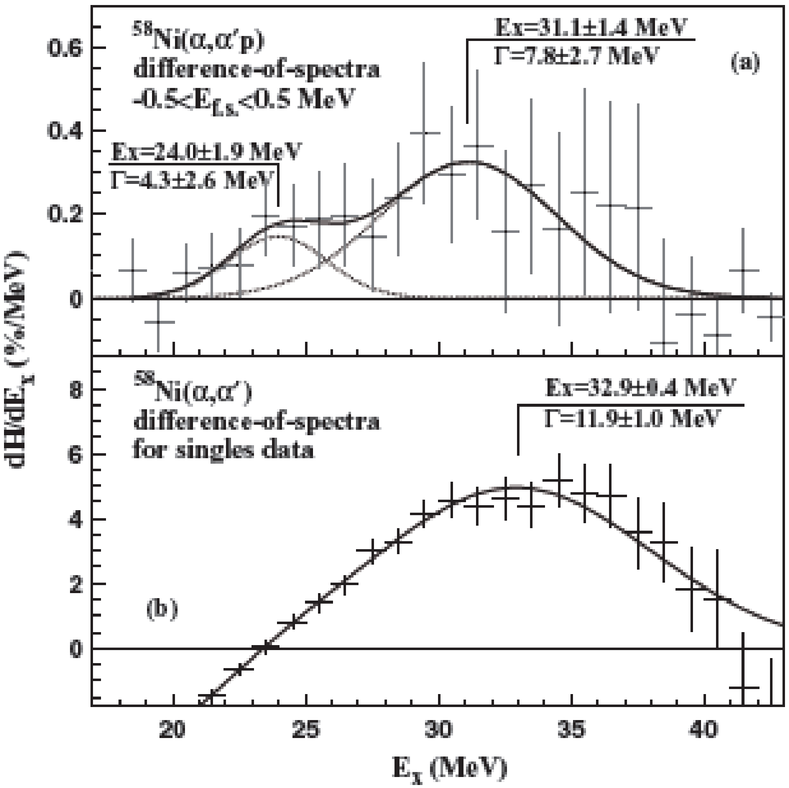
Experimental results

Difference of $E_{f.s.}$-spectra

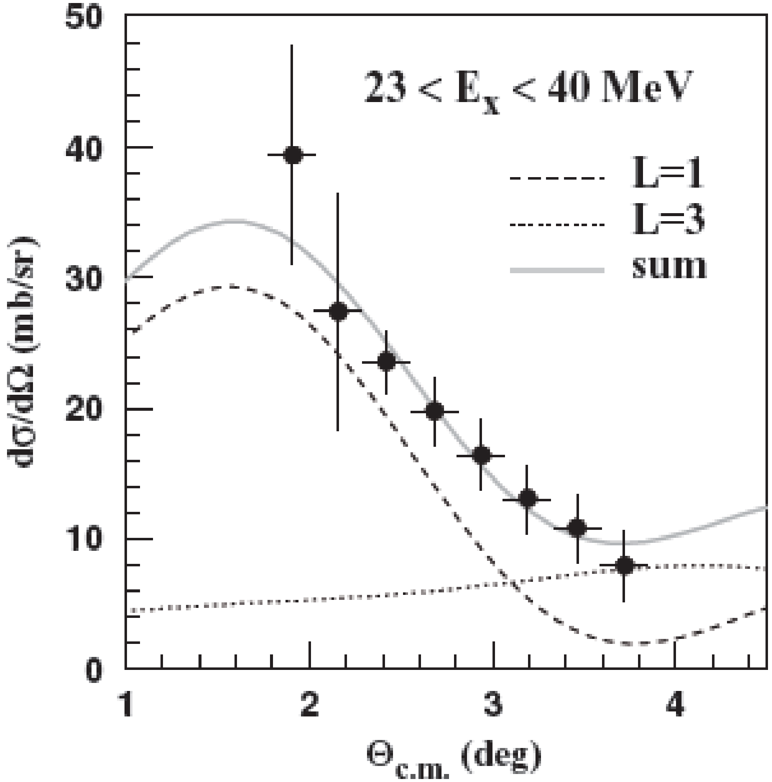




Strength distribution of ISGDR in ^{58}Ni



Spectra of L = 1 strengths obtained with DOS method in percentage of isoscalar EWSR; a) coincidence data gated on g.s. decay and b) singles data.



Differential cross section of resonance structure fitted with $L = 1$ and $L = 3$ DWBA calculations.

Proton-decay branching ratios
Normalized to 100%

	Exp. (%) (24-38 MeV)	Cal. (%) (15-40 MeV)
$7/2^-$	61.3 (with $5/2^-$)	47
$3/2^-$	7.9	3.1
$3/2^-$, $1/2^-$	9.9	2.2 (only for 1/2-)
$5/2^-$	3.2 ± 3.4	-
$1/2^+$	2.0 ± 4.2	13.4
$3/2^+$, $5/2^+$	15.9	34.3
Σ	100 %	100%

Calculations: M.L. Goerlik, I.V. Safonov, and M.H. Urin, Phys. Rev. C69 (2004) 054322

Conclusions!

- **There has been much progress in understanding ISGMR & ISGDR macroscopic properties**

  Systematics: E_x, Γ, %EWSR

  $\Rightarrow K_{nm} \approx 240$ MeV

  $\Rightarrow K_{\tau} \approx -500$ MeV

- **Sn nuclei are softer than ^{208}Pb and ^{90}Zr.**
- **Recently, Microscopic Structure for a few nuclei**

  CRPA has some success in ^{208}Pb & ^{58}Ni but fails badly in ^{116}Sn & ^{90}Zr.

- **Possible observation quadrupole compression mode, i.e. overtone of ISGQR**

Outlook

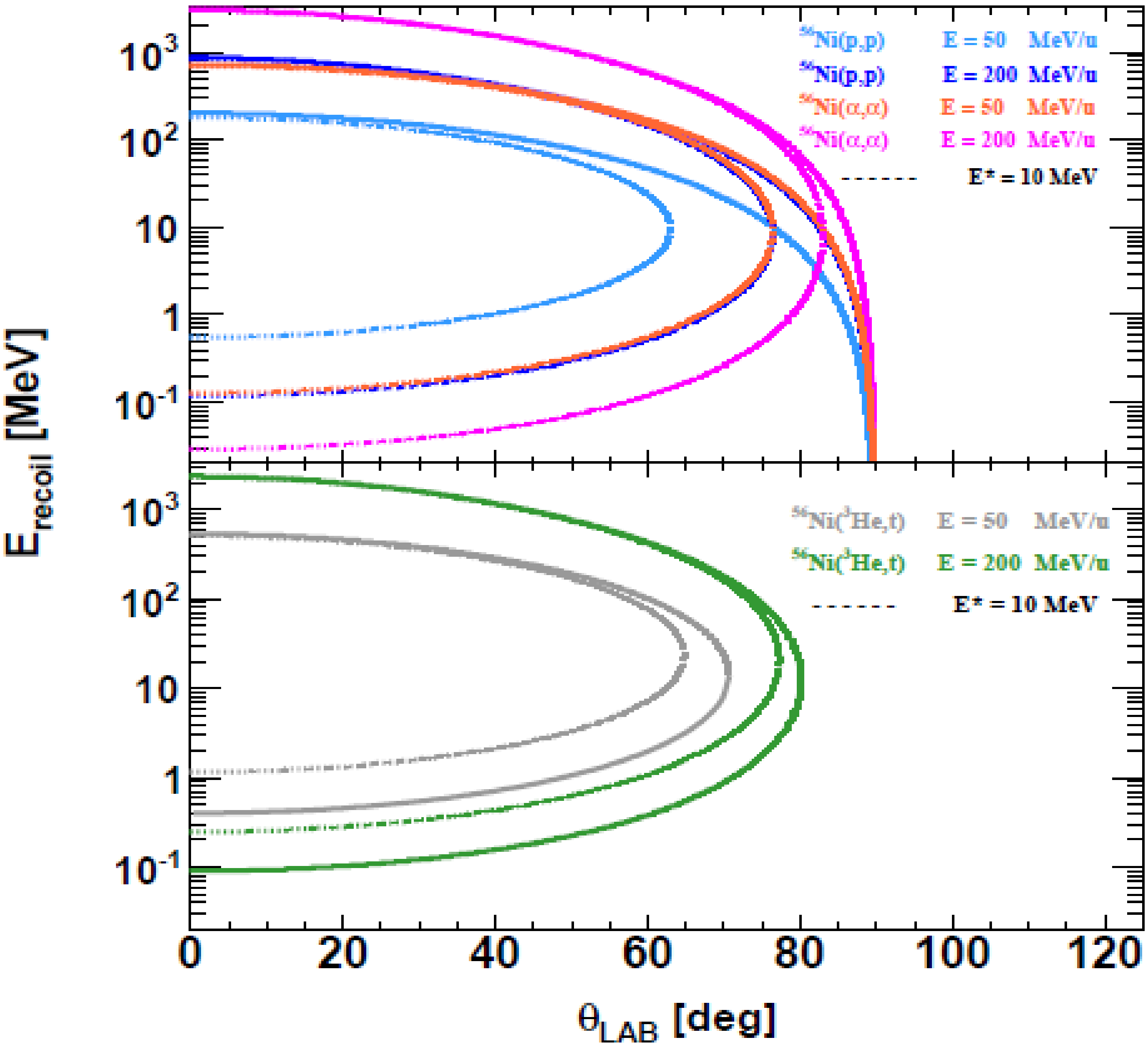
Radioactive ion beams will be available at energies where it will be possible to study excitation of ISGMR and ISGDR

RIKEN, FAIR, SPIRAL2, NSCL, EURISOL

Determine ISGMR and ISGDR in unstable Sn nuclei.

A = 106 to 134 possible

⇒A more precise determination of K_τ



^{56}Ni(p,p) E = 50 MeV/u
^{56}Ni(p,p) E = 200 MeV/u
^{56}Ni(α,α) E = 50 MeV/u
^{56}Ni(α,α) E = 200 MeV/u
E* = 10 MeV
^{56}Ni(^{3}He,t) E = 50 MeV/u
^{56}Ni(^{3}He,t) E = 200 MeV/u
E* = 10 MeV
E_{recoil} [MeV]
θ_{LAB} [deg]

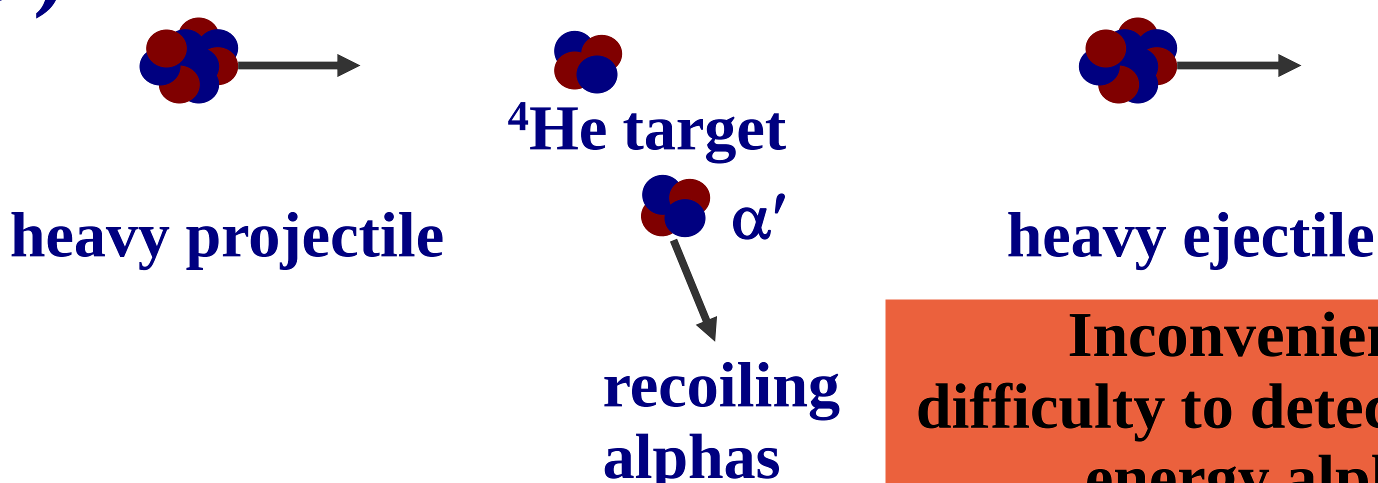
Nuclear structure studies with reactions in inverse kinematics

- Possible at FAIR and RIKEN

(beam energies of 50-100 MeV/u are needed!)

Approach (at FAIR): measure the recoiling alphas

(α,α')



Inconvenience: difficulty to detect the low-energy alphas

Detection system @ FAIR

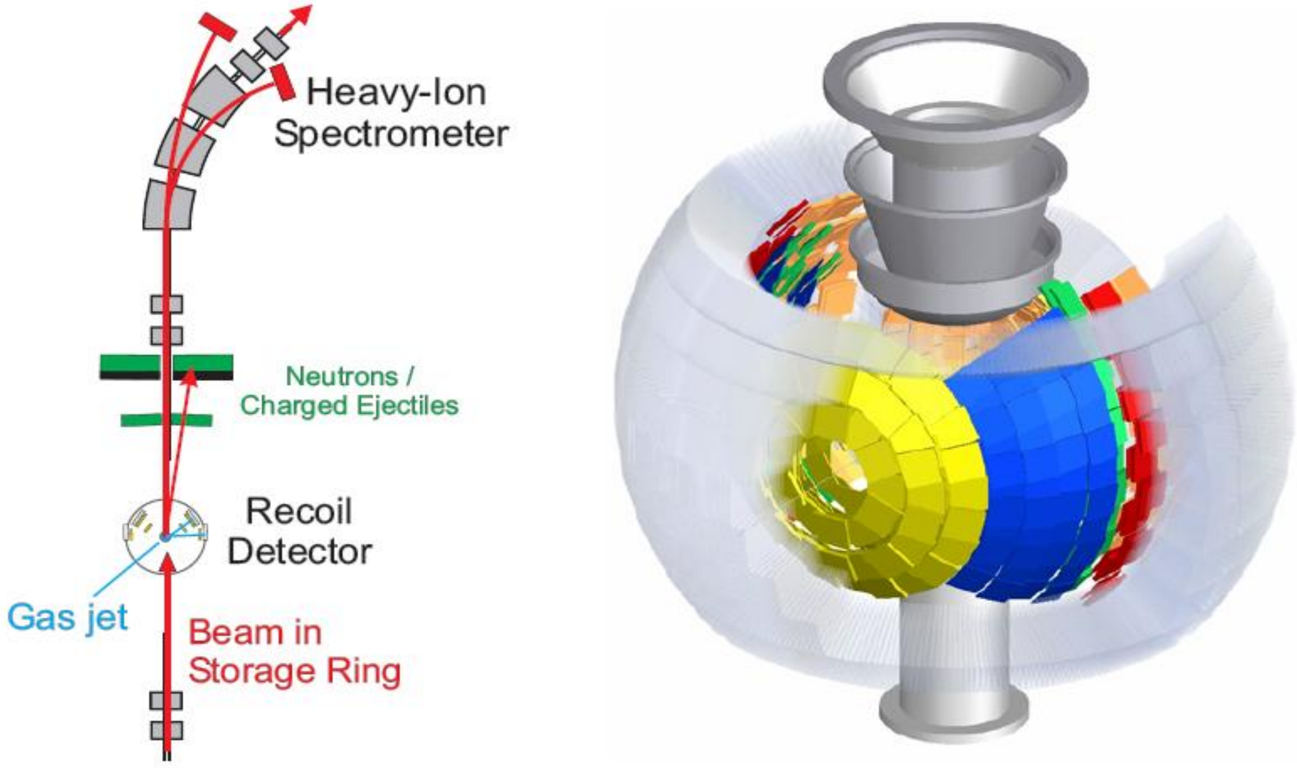
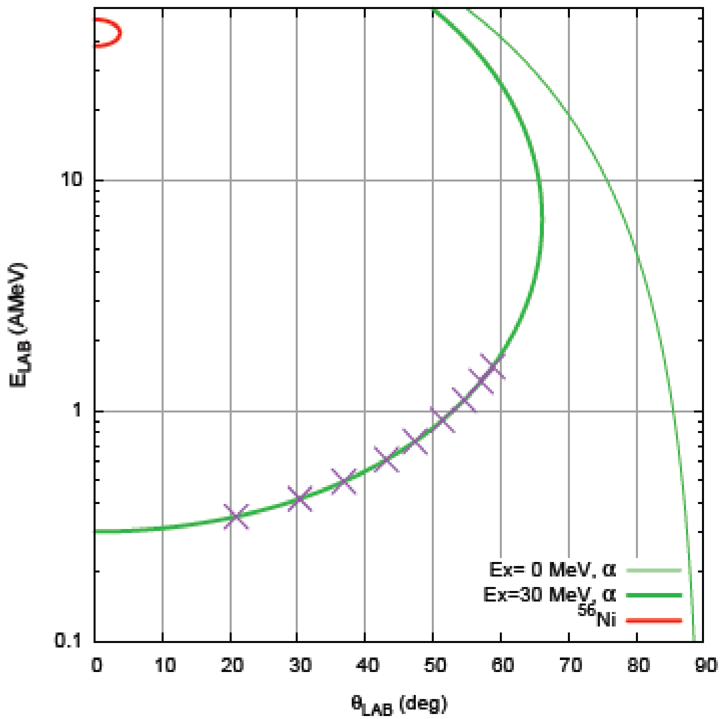
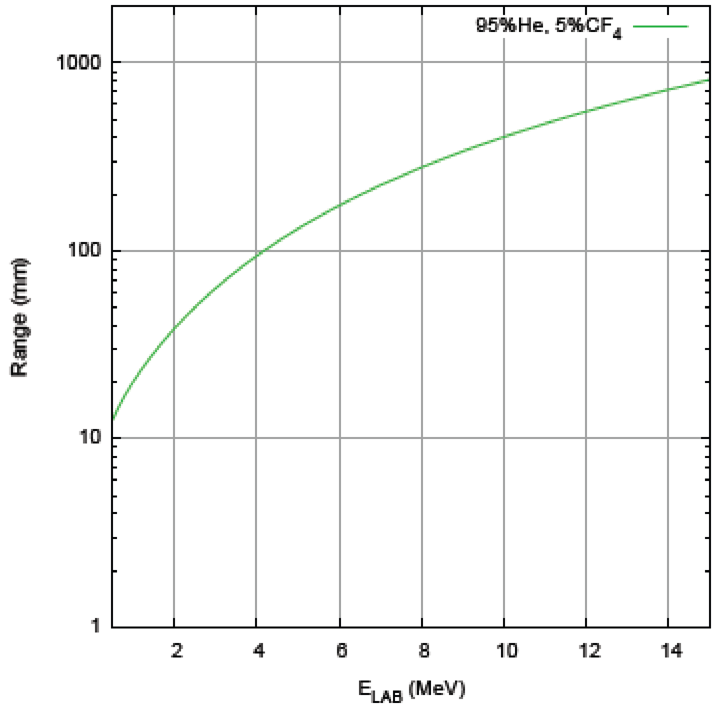


Figure 1: *Schematic view of the EXL detection systems. Left: Set-up built into the NESR storage ring. Right: Target-recoil detector surrounding the gas-jet target.*

Use of EXL recoil detector is under evaluation

Kinematics: energies





ISGDR in ^{56}Ni: 2nd ISGDR in Ni after ^{58}Ni

Conditions

- ^{56}Ni(α,α') and ^{56}Ni(α,α'p)
- LISE beam at the highest energy possible ($\geq$ 50*A* MeV)

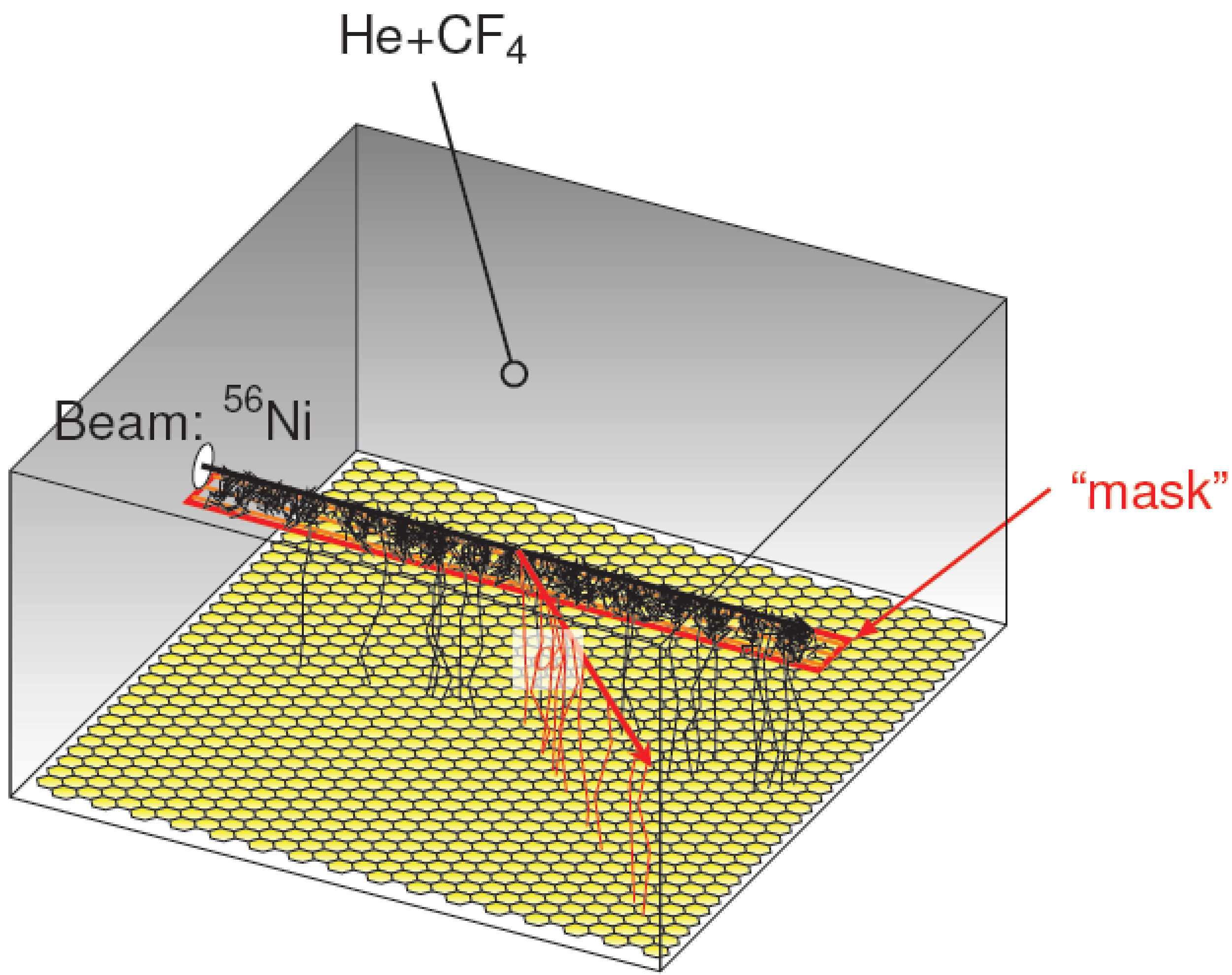
Setup

- MAYA in D6, filled w/ He (>95%), CF_4 mixture.
- All ancillary detectors (Si & CsI + S1)
- "Active" masking of the beam

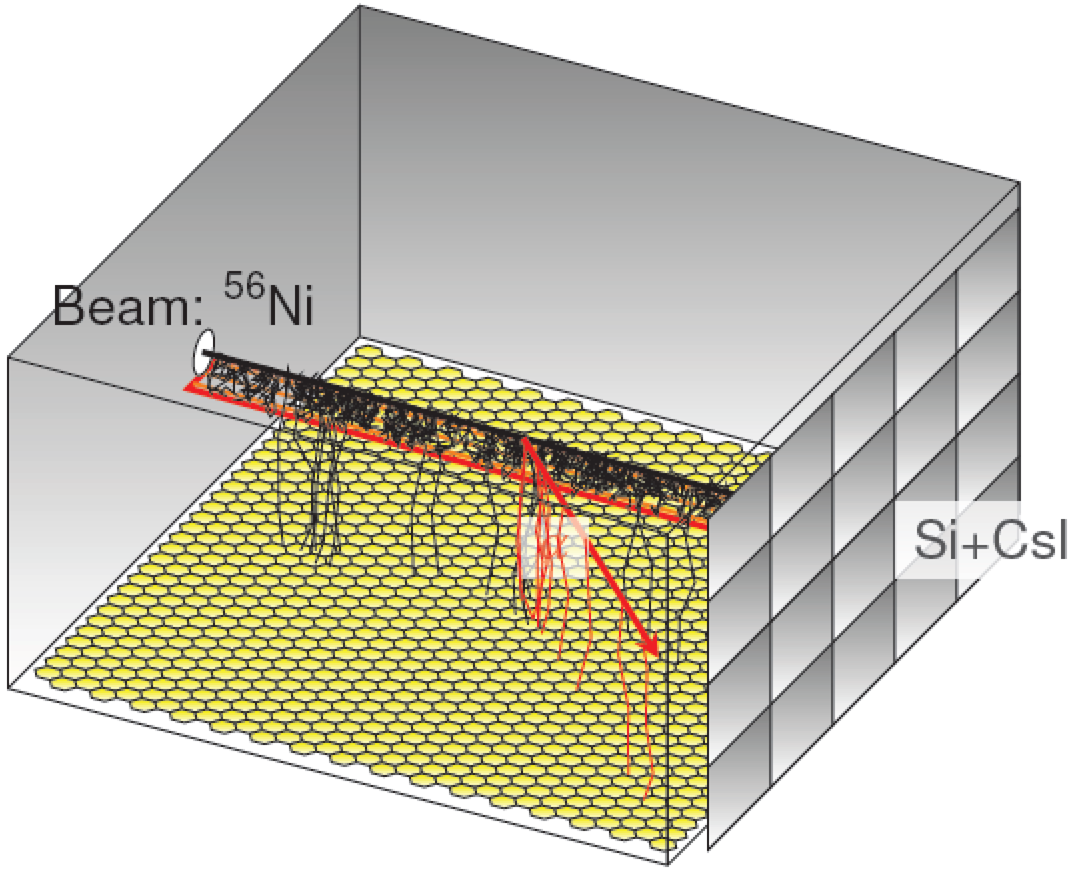
Predictions

- Detect alpha down to $\approx$ 1 MeV, between [2°,10°] in C.M.
- Remeasure ISGMR in ^{56}Ni
- Detect protons coming from ISGDR decay

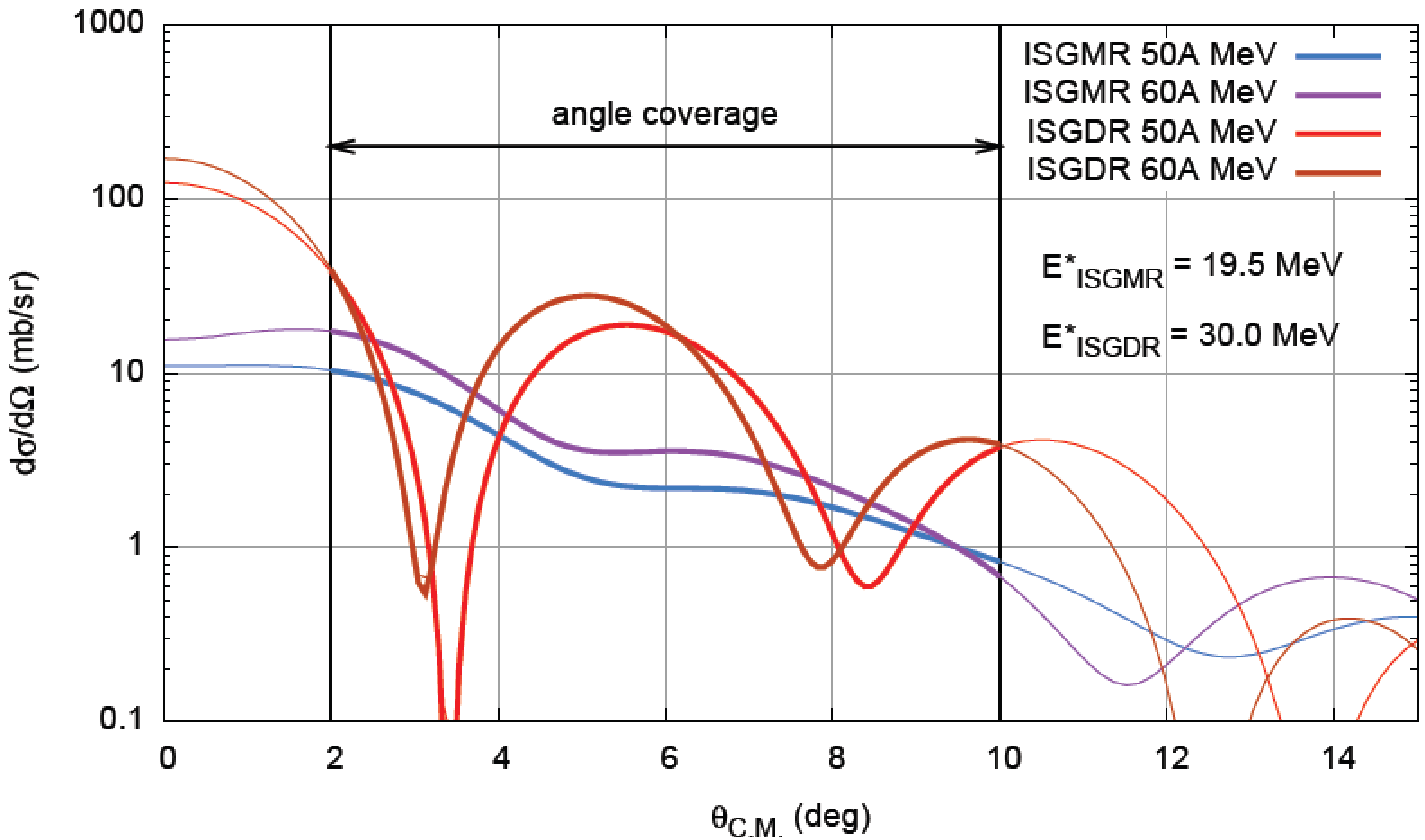
Setup: MAYA (Active target)



Setup: MAYA (Active target)



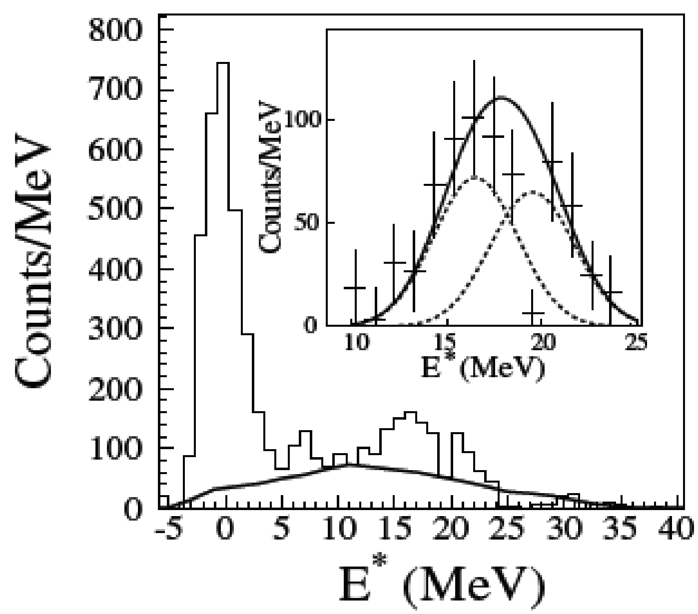
Predicted angular distribution

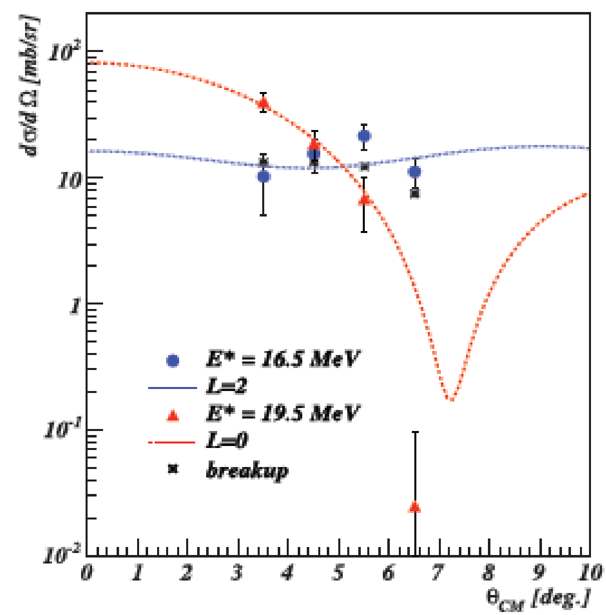


ISGMR in unstable nuclei

^{56}Ni: active target MAYA filled with deuterium gas
bombarding energy of 50*A* MeV

C. Monrozeau *et al.* @ GANIL





Excitation of the isovector GDR by inelastic α-scattering as a measure of the neutron skin of nuclei

A. Krasznahorkay *et al*.,

Nucl. Phys. A567 (1994) 521

Measuring the cross section of the isovector giant dipole resonance (IVGDR) for

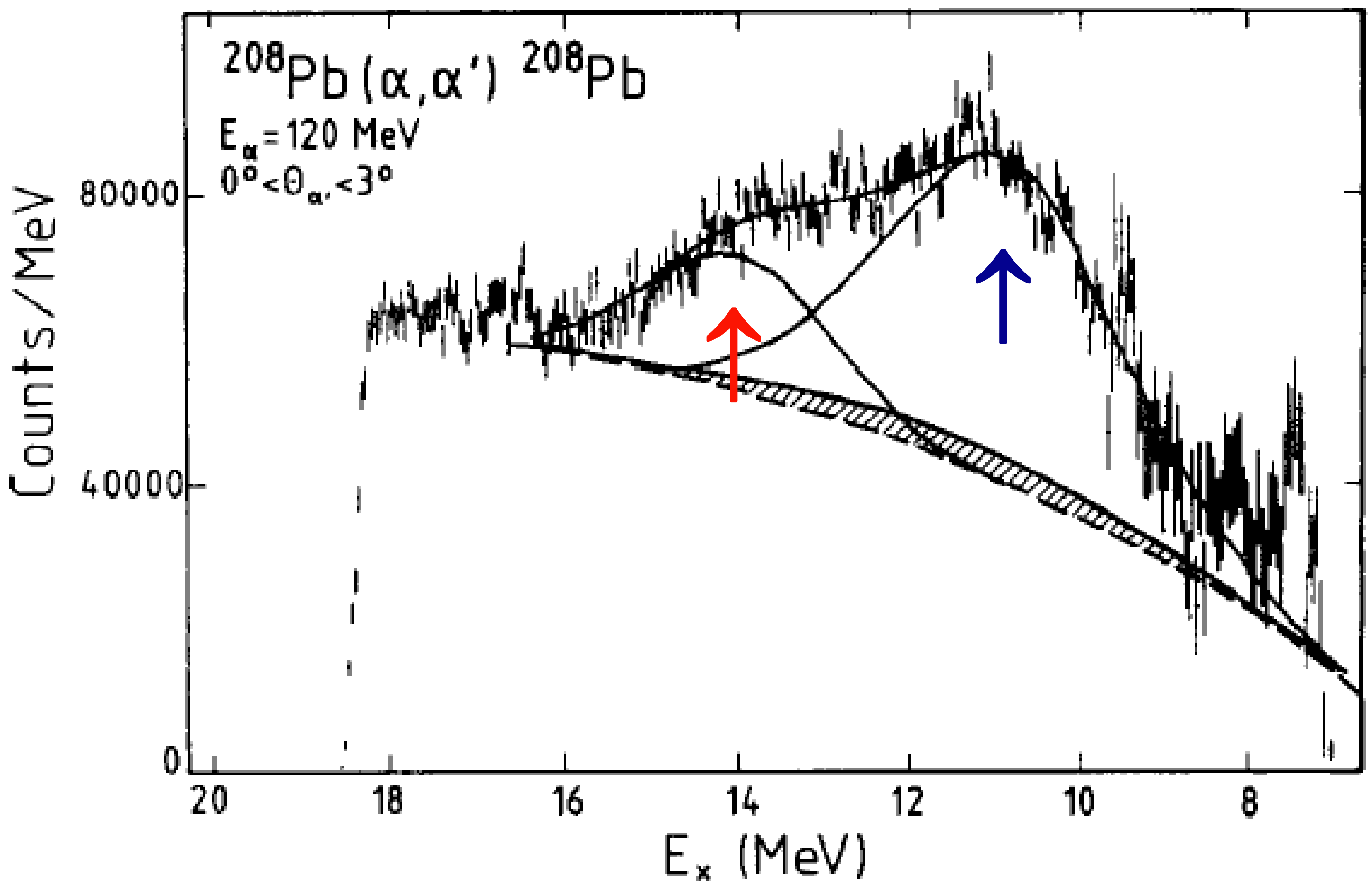
^{116,124}Sn($\alpha,\alpha'\gamma_0$), ^{150}Nd($\alpha,\alpha'\gamma_0$) and ^{208}Pb($\alpha,\alpha'\gamma_0$) reactions at E_α = 120 MeV and $0^o \leq \theta_{\alpha'} \leq 3^o$

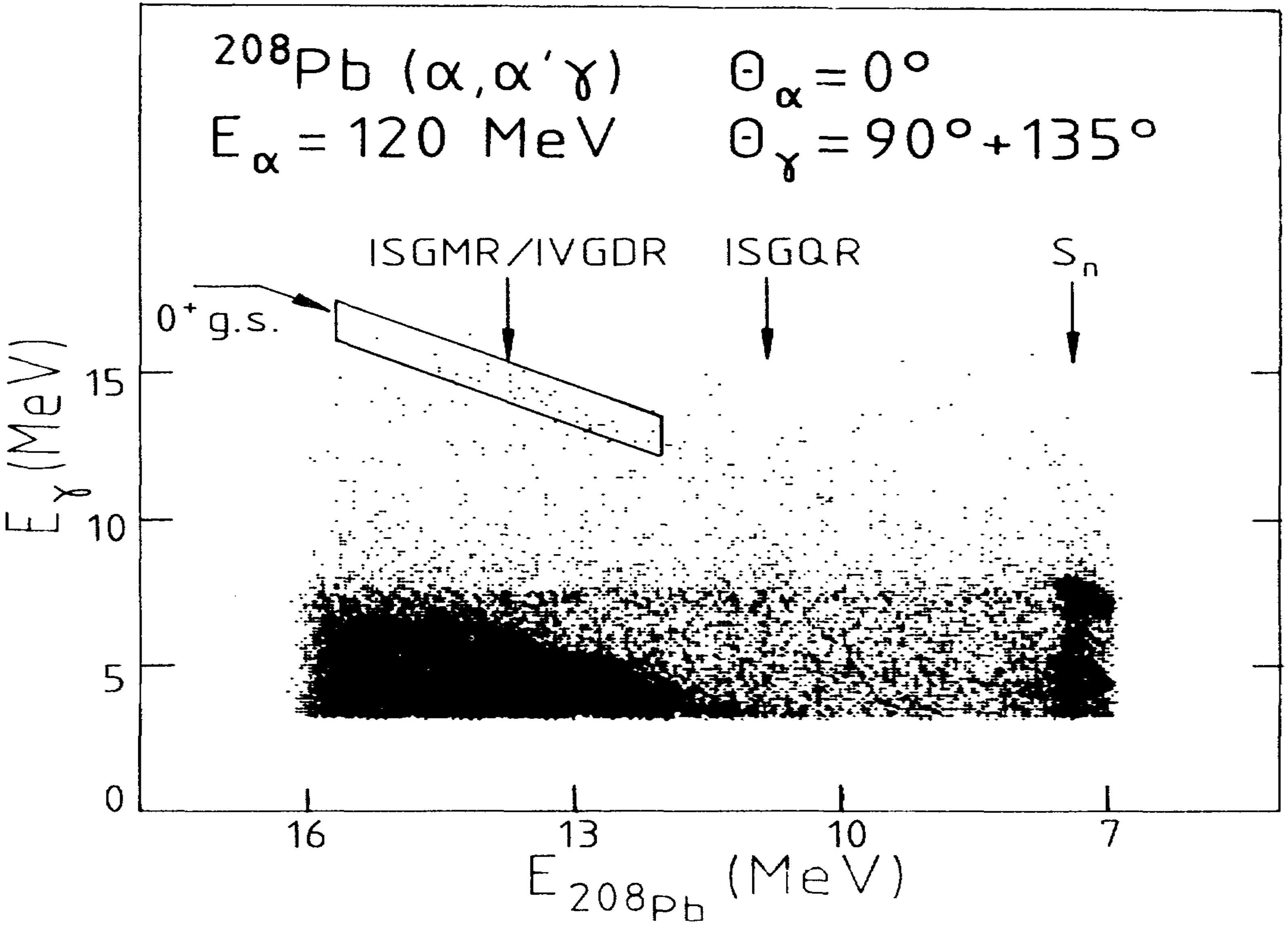
⇒ Determine neutron-skin thickness

ISGQR at 10.9 MeV

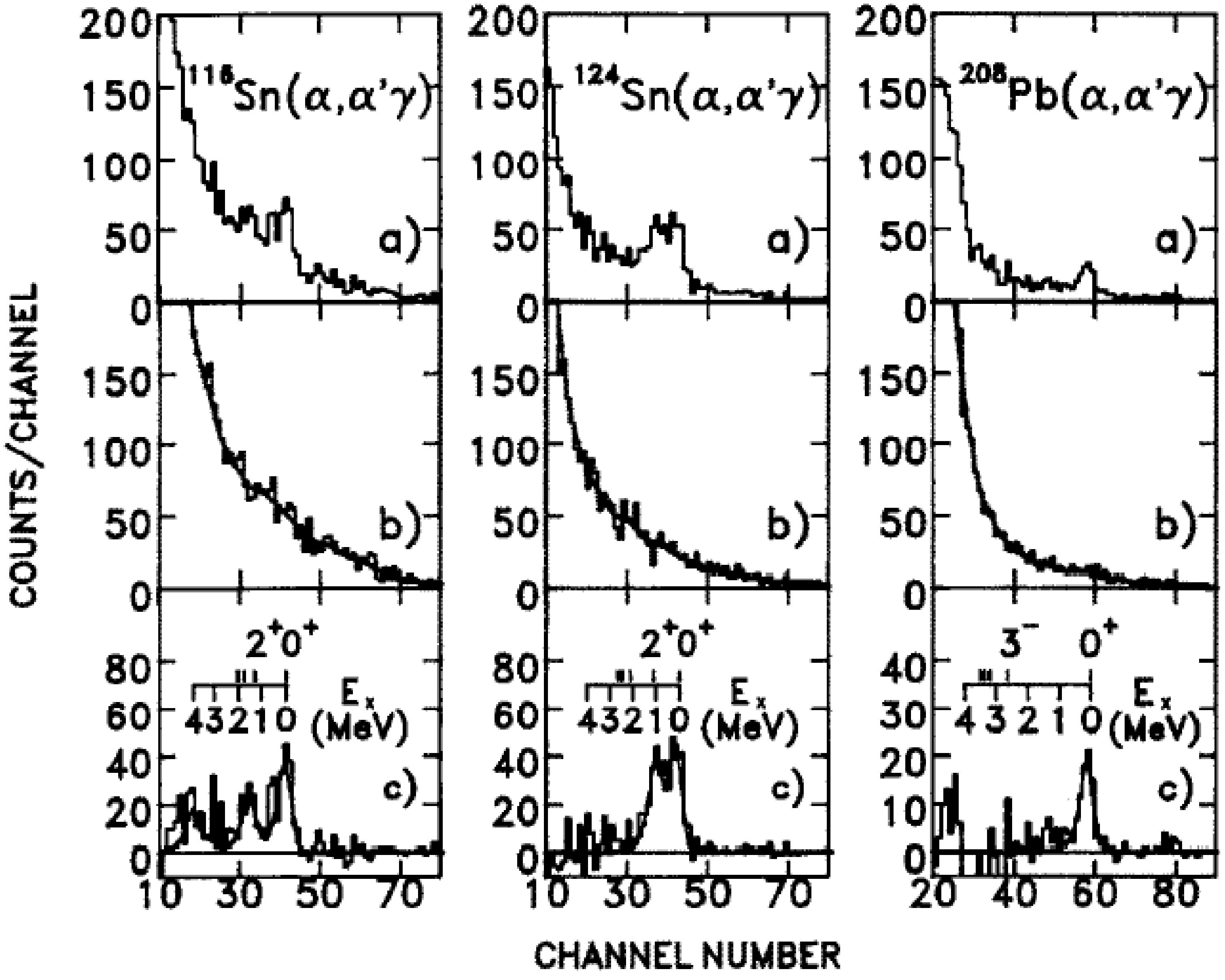
ISGMR at 13.9 MeV

Hatched area ⇒ IVGDR contribution (Coulomb + nuclear)





Two-dimensional spectrum showing inelastic α-scattering in coincidence with γ-decay



True
+
Random

Random

True

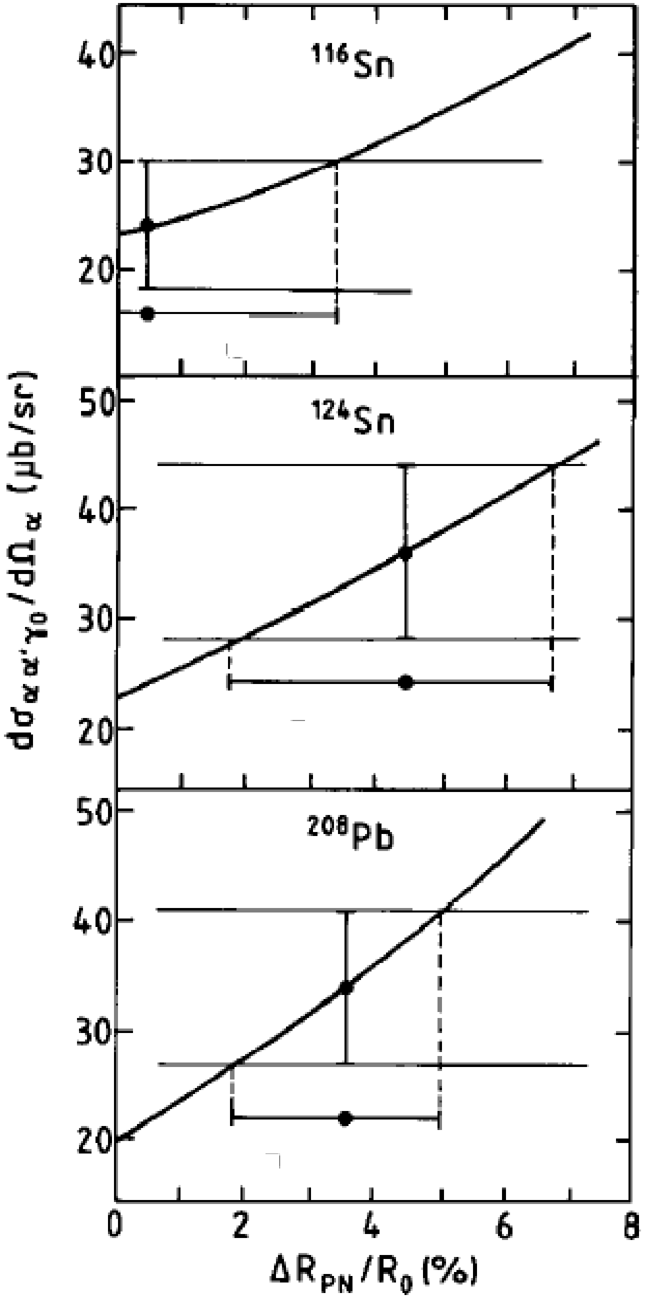
Isoscalar transition density in the Goldhaber-Teller model for excitation of IVGDR in inelastic α-scattering.

$$g_1^{10}(r) = g_1^{n}(r) - g_1^{p}(r) = \alpha_1 \gamma \left(\frac{N-Z}{A} \right) \left(\frac{d\rho(r)}{dr} + \frac{1}{3} c \frac{d^2\rho(r)}{dr^2} \right).$$

$$\frac{\Delta R_{PN}}{R_0} = \frac{R_n - R_p}{\frac{1}{2}(R_n + R_p)} = \gamma \frac{2(N-Z)}{3A}.$$

Therefore, DWBA cross sections can be calculated as function of $\Delta R_{pn}/R_0$ for the Goldhaber-Teller model and similarly for the Steinwedel-Jensen model.

γ-decay branching ratios are known from photoabsorption experiments.



Full line is the calculated $\alpha\gamma_0$ coincidence cross section, averaged over the solid angle of the α-particle and integrated over the full γ-ray solid angle (4π) and over the ΔE energy range as function of $\Delta R_{pn}/R$.

The experimental $\alpha\gamma_0$ cross sections for the IVGDR are shown as full circles with vertical error bars. The deduced values for $\Delta R_{pn}/R$ with the associated uncertainty (full circles with horizontal error bars) are also indicated.

Isotope	Present work $\Delta R_{PN}/R_0$ (%)	Present work ΔR_{PN} (fm)	Batty et al. [2] ΔR_{PN} (fm)	Angeli et al. [6] ΔR_{PN} (fm)
^{116}Sn	0.5 ± 2.7	0.02 ± 0.12	0.15 ± 0.05	0.13
^{124}Sn	4.4 ± 2.4	0.21 ± 0.11	0.25 ± 0.05	0.22
^{208}Pb	$3.5^{+1.5}_{-1.6}$	0.19 ± 0.09	0.14 ± 0.04	0.22

Participants

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E605: ISGDR in ^{58}Ni

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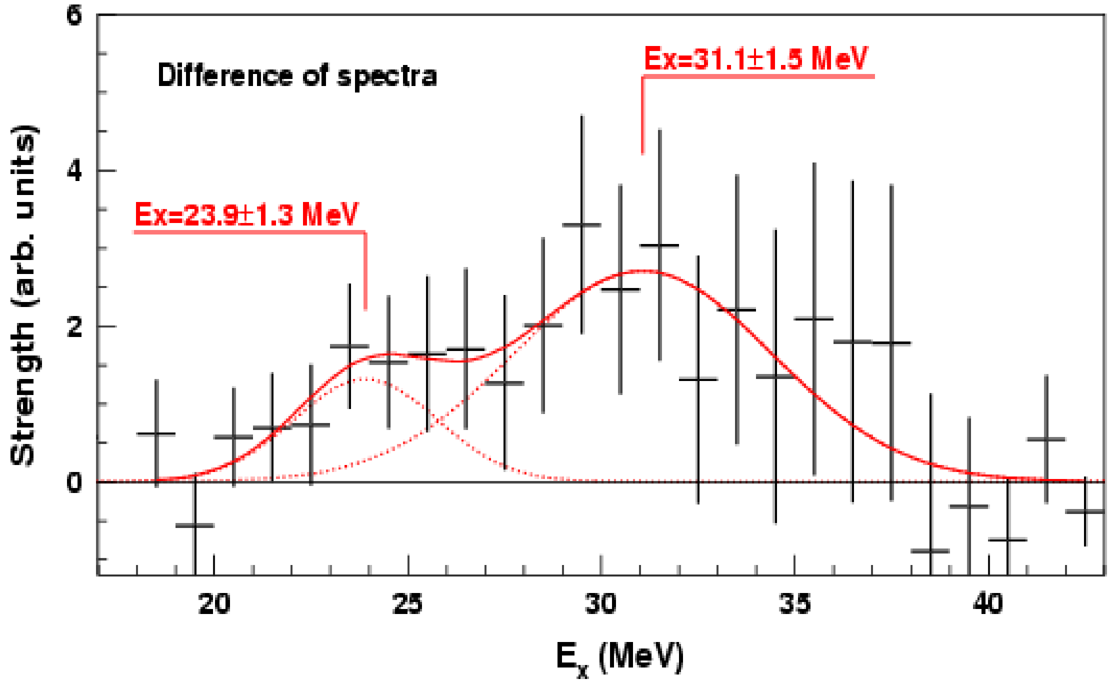
Thomas Roger	Sara Sambi	Hervé Savajols
Christelle Stodel	Laura Suen	Jean Charles Thomas
Marine Vandebrouck	Jarno van de Walle	

Strength distribution of ISGDR in ^{58}Ni

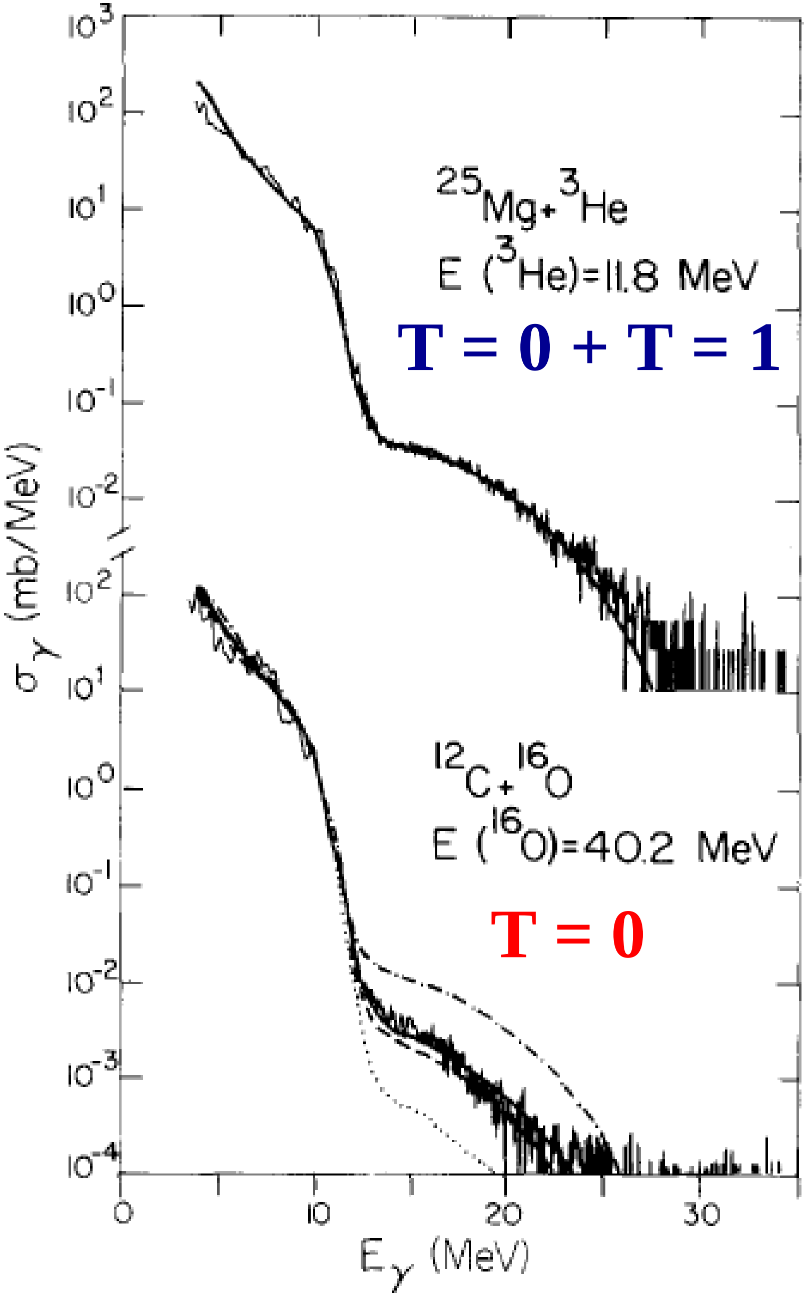
Double differential cross section of DOS divided by double differential cross section calculated by DWBA in 1.8 ~2.5 degree (dipole favoring angle region)



Centroid energies:

<E>= 31.1 ± 1.5 MeV higher-lying peak

<E>= 29.7 ± 0.8 MeV for 20<E_x<40 MeV



Role of isospin in the statistical decay of the giant dipole resonance built on excited states

M.N. Harakeh, D.H. Dowell, G. Feldman, E.F. Garman, R. Loveman, J.L. Osborne and K.A. Snover

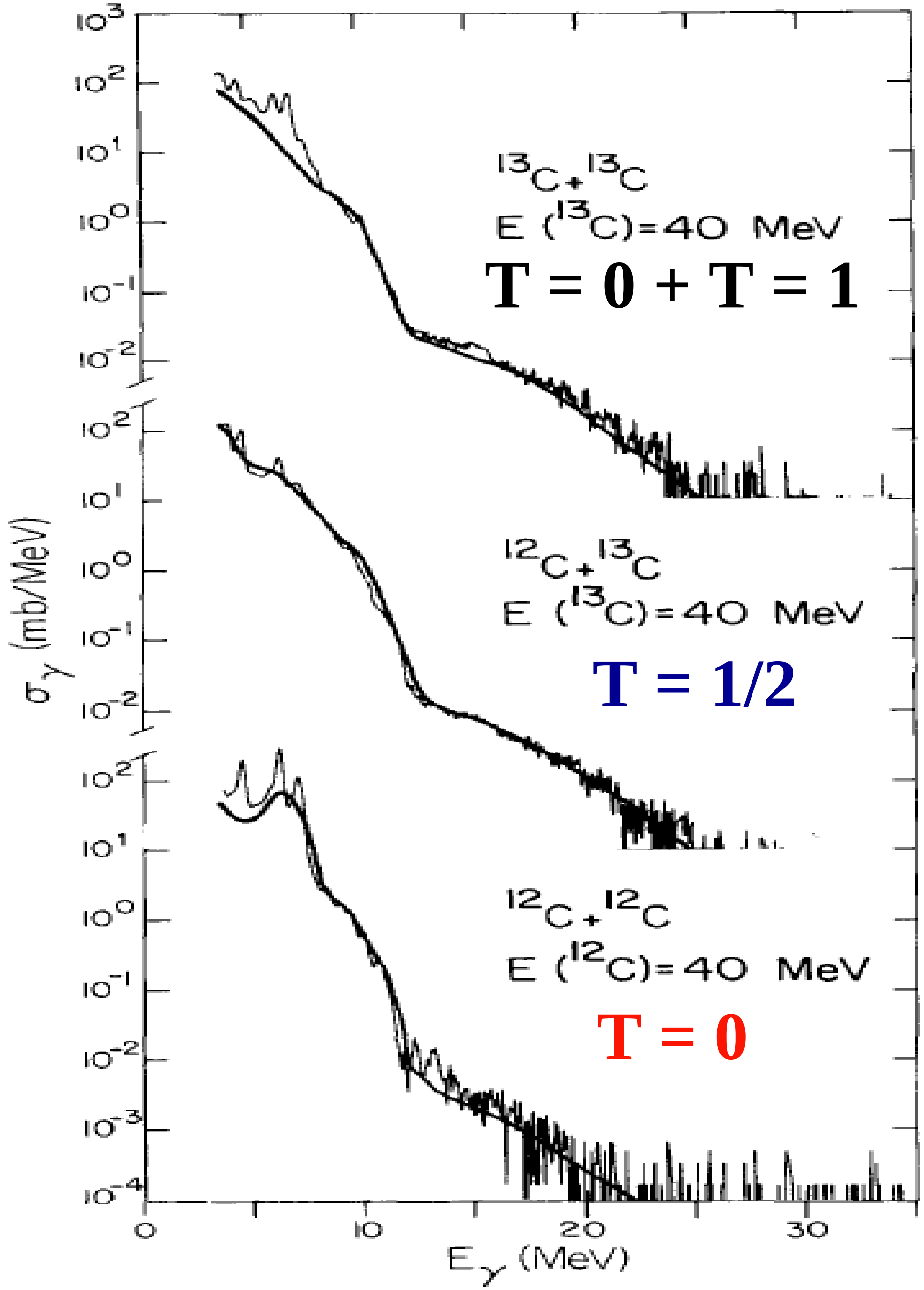
Phys. Lett. B176 (1986) 297

Dotted: pure ISGQR

Dashed: pure isospin

Dash-dotted: complete isospin mixing

Solid: isospin mixing (~ 5%)


$^{13}C+^{13}C$
$E(^{13}C)=40$ MeV
T = 0 + T = 1
$^{12}C+^{13}C$
$E(^{13}C)=40$ MeV
T = 1/2
$^{12}C+^{12}C$
$E(^{12}C)=40$ MeV
T = 0
σ_γ (mb/MeV)
E_γ (MeV)